# Models that link and suggest data about elementary particles, dark matter, and the cosmos 

Thomas J. Buckholtz

Ronin Institute for Independent Scholarship, Montclair, New Jersey 07043, USA


#### Abstract

We suggest progress regarding the following six physics opportunities. List all elementary particles. Describe dark matter. Explain ratios of dark matter to ordinary matter. Explain eras in the history of the universe. Link properties of objects. Interrelate physics models. We use models based on Diophantine equations.


Keywords: Beyond the Standard Model, Dark matter, Galaxy formation, Neutrino masses, Evolution of the universe

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## 1. Introduction

This essay pursues the following two challenges. Describe new elementary particles, dark matter, and dark energy. Use descriptions of elementary particles, dark matter, and dark energy to explain astrophysics data and cosmology data.

Our work has roots in the known elementary particles and in concordance cosmology. Our work features a hypothesis that nature includes six isomers of known elementary particles - of which only one isomer associates with ordinary matter and five isomers associate with most dark matter. Our work features new modeling based on Diophantine equations.

We suggest new elementary particles, insight regarding modeling regarding gravity, a specification for dark matter, and insight regarding galaxy formation.

We suggest explanations for known data for which - seemingly - other modeling does not offer explanations. For example, we show bases for at least a half-dozen seemingly significant ranges of observed ratios of dark matter effects to ordinary matter effects.

We suggest insight about mechanisms that shape eras in the history of the universe - including the two known eras, the possible inflationary epoch, and two possible earlier eras. For example, we suggest a resolution regarding tensions between cosmology data and concordance cosmology.

We suggest other data that people might be able to verify or refute. For example, we suggest masses for the neutrinos, more accurate masses for some other known elementary particles, and masses for elementary particles that we suggest.

Figure 1 shows flow - from roots through results - regarding our research.

Our explanations regarding large-scale data might help validate our descriptions of possible new elementary particles, our description of dark matter, and our modeling regarding dark energy.

## 2. Methods

This unit develops and deploys modeling that does the following. Associate properties of objects with modeling regarding electromagnetism, gravity, and possible other long-range interactions. Point to all known elementary particles. Suggest new elementary particles. Suggest specifications for dark matter.

We anticipate the following notions.

- We consider two sets of elementary particles. One set associates with the term LRI, as in long-range interaction. This set includes the photon, possibly a graviton, possibly a spinthree relative of the photon, and possibly a
spin-four relative of the photon. One set associates with the word simple. The three-word term simple elementary particles denotes all elementary particles other than LRI elementary particles.
- We consider solutions to the equation $\Sigma=$ $\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right| . \quad \Sigma$ is a nonnegative integer. Each one of the one or more $l_{o}$ is a positive integer. (Perhaps, see equation (2).)
- Solutions for which $\Sigma=0$ point to all known simple elementary particles and to all other possible simple elementary particles that our work suggests.
- Solutions for which $1 \leq \Sigma \leq 4$ point to the known LRI elementary particle (or, to the photon) and to all other possible LRI elementary particles that our work suggests. For example, solutions for which $\Sigma=1$ associate with electromagnetism and the photon. Solutions for which $\Sigma=2$ associate with gravity and the possible graviton.
- Solutions for which $1 \leq \Sigma \leq 4$ and $\Sigma \in \Gamma$ point to nominal properties of objects. Such nominal properties include charge, magnetic moment, mass, moments of inertia, a function of flavour (for elementary fermions), spin (for simple elementary particles), and other properties that people can measure or infer.
- Solutions for which $1 \leq \Sigma \leq 4$ and $\Sigma \notin \Gamma$ point to anomalous properties. Such anomalous properties include anomalous magnetic moment.
- Nature includes six isomers (or, similar - but not identical - copies) of the set of simple elementary particles. One isomer provides a basis for "somewhat all" ordinary matter and some dark matter. (Here, our uses of the twoword term "somewhat all" associate with notions that these sentences do not fully account for LRI aspects, such as photons.) The other five isomers provide bases for most dark matter (or, for "somewhat all" dark matter other than the aforementioned some dark matter).

Table 1 summarizes aspects to which solutions to $\Sigma=\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right|$ point.

### 2.1. Charge, mass, and other properties

This unit develops and deploys modeling that interrelates long-range forces (such as electromagnetism and gravity) and properties of objects.

Research - elementary particles, dark matter, and the cosmos


Figure 1: Research flow - from roots to results.

Table 1: Aspects to which solutions to $\Sigma=\left|\sum_{l_{0} \in \Gamma}\left( \pm l_{o}\right)\right|$ point. Here, $\Sigma$ is a nonnegative integer and each one of the one or more $l_{o}$ is a positive integer. (See equation 2l.) The three-word phrase simple elementary particles denotes all elementary particles except LRI elementary particles. LRI (or, long-range interaction) elementary particles include the photon, the possible graviton, and perhaps some other possible elementary particles. Possible elementary particles for which the isomer zero copies would measure as dark matter include three zero-charge analogs to quarks and three heavy neutrinos. Visibly observed dark matter galaxies include some ordinary matter stars. This essay designates the six isomers of simple elementary particles via the two-word terms isomer zero, isomer one, ..., and isomer five.

| Solutions | Solutions point to ... | Examples |
| :--- | :--- | :--- |
| $\Sigma=0$ | Simple elementary particles | Higgs boson, quarks, electron, W boson, <br> possible inflaton |
| $1 \leq \Sigma \leq 4$ | LRI elementary particles | Photon, possible graviton <br> Charge, nominal magnetic moment, <br> $1 \leq \Sigma \leq 4$ and $\Sigma \in \Gamma$ |
| Nominal properties of objects | mass, moments of inertia |  |
| $1 \leq \Sigma \leq 4$ and $\Sigma \notin \Gamma$ | Anomalous properties | Anomalous magnetic moment |
| $\Sigma=0-$ isomer zero | Ordinary matter (except LRI <br> aspects) and some dark matter | Visible aspects of the universe |
| $\Sigma=0$ - isomers one | Most dark matter | Most stuff within a dark matter galaxy |
| through five |  |  |

### 2.1.1. We explore notions that point toward aspects of some modeling that our work uses.

We imagine two non-moving objects - object A and object B - that are located a nonzero distance $r$ from each other.

We imagine hypothetical effects that associate with hypothetical interactions by object B with hypothetical fields produced by object A. (Notions of such fields might parallel notions of an electric field or notions of a gravitational field.)

We discuss aspects of hypothetical particles that might associate with interactions between object A and object $B$. We use the two-item term object $C$ to denote such a hypothetical particle. We imagine that objects C traverse straight-line trajectories from object A to object B. We use the word axis to associate with the straight line.

We imagine objects C that have some similarities to and some differences from a planar solar system that includes one star and one or a few planets. Regarding an object C, the star moves along the axis between object A and object B. All of the planetary orbits exist in a plane that is perpendicular to the motion of object C. For this example, the only property - of object C - for which we have interest is the (net) total orbital angular momentum, with the total being the sum of the orbital angular momenta of the planets. We imagine that, with respect to the axis that runs from object A to object $B$, each orbit associates with a unique magnitude $l_{o} \hbar$ of orbital angular momentum. Here, $l_{o}$ is a positive integer. Up to one planet can occupy an orbit. The integer $l_{\text {max }}$ denotes the maximum value of $l_{o}$ that associates with an occupied orbit. Relative to the axis that runs from object A to object $B$, the angular momentum that associates with an occupied orbit is one of $-l_{o} \hbar$ and $+l_{o} \hbar$. (We exclude - for the occupied orbit that associates with $l_{o}$ - values of $l_{z}$ for which $-l_{o}<l_{z}<+l_{o}$.) The angular momentum that associates with an unoccupied orbit is $0 \hbar$. Relative to the axis, the total angular momentum that associates with an object C is the sum - over the occupied orbits - of the respective $\pm l_{o} \hbar$.

Regarding modeling that we discuss below, the following notions pertain.

- Individual objects C do not associate directly with elementary particles.
- Mathematical modeling regarding an object C associates with a so-called (mathematical) solution. Some sets of solutions associate with modeling for LRI (or, long-range interaction) elementary particles such as the photon. Some sets of solutions associate with modeling for SRI (or, short-range interaction) elementary bosons such as the W boson. Some sets of solutions associate with modeling for ELF
(or, elementary fermion) elementary fermions such as the electron. We use the three-word phrase simple elementary particles to denote SRI elementary bosons and ELF elementary fermions.
- We suggest that the modeling explains data (some of which other modeling seems not to explain), echoes useful modeling (that other people have developed), and predicts possibly reasonable data that might result from future observations and experiments.
- This essay does not necessarily point directly to a basis - for the modeling - that people would consider to be fundamental.


### 2.1.2. We develop a mathematical basis for model-

 ing that this essay features.We focus on mathematics that associates with the discussion above about objects C and orbits. We do not suggest a notion of direct physics relevance for such objects.

The discussion about hypothetical objects C suggests expressions of the form $\sum_{o \in O}\left( \pm l_{o}\right)$. Here, $O$ associates with the set of occupied orbits. The integer o denotes a member of $O$. The association $o=l_{0}$ pertains.

We use the symbol $\Gamma$ to denote an ascendingorder list of the relevant $o \in O$. Within a list $\Gamma$, we separate values of $l_{o}$ by using the symbol '. The symbol $l_{\text {max }}$ denotes the maximum value of $l_{o}$ in $\Gamma$. For example, $\Gamma=1^{‘} 3$ associates with $l_{\max }=3$ and with $1 \in \Gamma, 2 \notin \Gamma$, and $3 \in \Gamma$.

We define $l_{\Sigma}$ to be the sum of the various values of $\pm l_{o}$. Equation (1) pertains. We associate the word expression with equation (1).

$$
\begin{equation*}
l_{\Sigma}=\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right) \tag{1}
\end{equation*}
$$

We define $\Sigma$ to be the absolute value of the sum of the various values of $\pm l_{o}$. Equation (2) pertains. We associate the word solution with equation (2).

$$
\begin{equation*}
\Sigma=\left|l_{\Sigma}\right|=\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right| \tag{2}
\end{equation*}
$$

The two-word term Diophantine equations associates with the modeling that we pursue.

Table 2 alludes to all $l_{\Sigma}=\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)$ expressions for which $1 \leq l_{o} \leq l_{\max } \leq 4$ and no two values of $l_{o}$ are the same. The rightmost five columns discuss solutions $\Sigma=\left|l_{\Sigma}\right|=\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right|$. Regarding table 2, the notions of monopole, dipole, quadrupole, and octupole associate with the numbers of contributors to sets of mathematically related solutions. In table 2, the notions of monopole, dipole, quadrupole, and octupole do not necessarily associate with uses - in physics - of the terms

Table 2: $\Sigma=\left|l_{\Sigma}\right|=\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right|$ solutions, assuming that $1 \leq l_{o} \leq l_{\max } \leq 4$ and that no two values of $l_{o}$ are the same. The columns labeled $l_{1}$ through $l_{4}$ show contributions toward expressions $\bar{l}_{\Sigma}=\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)$. In those four columns, the symbol 0 is a placeholder for an unused pair, $-l_{o}$ and $+l_{o}$, of values. The symbol $n_{0}$ denotes the number of times the symbol 0 associates with an $l_{o}$ for which $1 \leq l_{o} \leq l_{\max } \leq 4$. The symbol $n_{\Gamma}$ denotes the number of elements in the list $\Gamma$. For each row, there are $2^{n_{\Gamma}}$ possible ways to assign signs regarding the set of $n_{\Gamma}$ terms. There are $2^{n_{\Gamma}}$ expressions of the form $l_{\Sigma}=\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)$. Thus, there are $2^{n_{\Gamma}-1}$ solutions $\Sigma=\left|l_{\Sigma}\right|=\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right|$. The $\Sigma$ column shows values of $\Sigma$ that associate with solutions. For example, for $l_{\max }=2$ and $\Gamma=1^{\prime} 2$, the two solutions feature, respectively, $\Sigma=1$ (as in $1=|-1+2|$ ) and $\Sigma=3$ (as in $3=|+1+2|$ ). (The two expressions that associate with $1=|-1+2|$ are $-1+2$ and $+1-2$. The two expressions that associate with $3=|+1+2|$ are $+1+2$ and $-1-2$.) The number $n_{\Sigma g \Gamma}$ equals $2^{n_{\Gamma}^{-1}}$ and states the number of solutions. The column for which the one-word label is notion associates with the number of solutions. For a row for which exactly one solution pertains, the column shows the word monopole. For a row for which exactly two solutions pertain, the column shows the word dipole. For a row for which exactly four solutions pertain, the column shows the word quadrupole. For a row for which exactly eight solutions pertain, the column shows the word octupole. For the case of octupole, each one of $\Sigma=2$ and $\Sigma=4$ associates with two solutions. Regarding $\Sigma=2$, $|-1+2-3+4|=2=|-1-2-3+4|$. Regarding $\Sigma=4,|-1-2+3+4|=4=|+1+2-3+4|$.

| $l_{\text {max }}$ | $\Gamma$ | $l_{1}$ | $l_{2}$ | $l_{3}$ | $l_{4}$ | $\Sigma$ | $n_{0}$ | $n_{\Gamma}$ | $n_{\Sigma \mathrm{g} \Gamma}$ | Notion |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | $\pm 1$ | - | - | - | 1 | 0 | 1 | 1 | Monopole |
| 2 | 2 | 0 | $\pm 2$ | - | - | 2 | 1 | 1 | 1 | Monopole |
| 2 | 1'2 | $\pm 1$ | $\pm 2$ | - | - | 1,3 | 0 | 2 | 2 | Dipole |
| 3 | 3 | 0 | 0 | $\pm 3$ | - | 3 | 2 | 1 | 1 | Monopole |
| 3 | $1 \times 3$ | $\pm 1$ | 0 | $\pm 3$ | - | 2,4 | 1 | 2 | 2 | Dipole |
| 3 | $2^{\prime} 3$ | 0 | $\pm 2$ | $\pm 3$ | - | 1,5 | 1 | 2 | 2 | Dipole |
| 3 | $1^{\prime} 2^{\prime} 3$ | $\pm 1$ | $\pm 2$ | $\pm 3$ | - | 0,2,4,6 | 0 | 3 | 4 | Quadrupole |
| 4 | 4 | 0 | 0 | 0 | $\pm 4$ | 4 | 3 | 1 | 1 | Monopole |
| 4 | 1'4 | $\pm 1$ | 0 | 0 | $\pm 4$ | 3,5 | 2 | 2 | 2 | Dipole |
| 4 | $2^{\prime} 4$ | 0 | $\pm 2$ | 0 | $\pm 4$ | 2,6 | 2 | 2 | 2 | Dipole |
| 4 | 34 | 0 | 0 | $\pm 3$ | $\pm 4$ | 1,7 | 2 | 2 | 2 | Dipole |
| 4 | 1'2‘4 | $\pm 1$ | $\pm 2$ | 0 | $\pm 4$ | 1,3,5,7 | 1 | 3 | 4 | Quadrupole |
| 4 | $1{ }^{6} 34$ | $\pm 1$ | 0 | $\pm 3$ | $\pm 4$ | 0,2,6,8 | 1 | 3 | 4 | Quadrupole |
| 4 | $2^{\prime} 3^{〔} 4$ | 0 | $\pm 2$ | $\pm 3$ | $\pm 4$ | 1,3,5,9 | 1 | 3 | 4 | Quadrupole |
| 4 | $1^{\prime} 2^{\prime} 3^{\prime} 4$ | $\pm 1$ | $\pm 2$ | $\pm 3$ | $\pm 4$ | 0,2,2,4,4,6,8,10 | 0 | 4 | 8 | Octupole |

monopole, dipole, quadrupole, and octupole. (We anticipate that discussion regarding table 3 suggests uses regarding physics.)

We use the symbol $\Sigma g \Gamma$ to denote the combination of a list $\Gamma$ and a relevant value of $\Sigma$. The letter g anticipates an association with electromagnetism and an association with gravity. (Perhaps, think of g as in gamma rays and g as in gravity. We anticipate that $1 \mathrm{~g} \Gamma$ solutions associate with electromagnetism and that $2 \mathrm{~g} \Gamma$ solutions associate with gravity. We anticipate that our modeling associates - with each other - aspects of electromagnetism and aspects of gravitation. Other people discuss possibilities for relationships between electromagnetism and gravity. For example, reference [1] explores notions of a coupling between electromagnetism and gravity. Reference [2] and reference [3] discuss Einstein-Maxwell equations that suggest combining electromagnetic stress-energy tensors and the Einstein field equations, which have origins in modeling regarding gravitation.)

We associate the symbol $\Sigma \mathrm{g}$ with solutions of the form $\Sigma g \Gamma$. We associate the symbol $\Sigma g^{\prime}$ with $\Sigma \mathrm{g}$ solutions for which $\Sigma \in \Gamma$. We associate the symbol $\Sigma g^{\prime \prime}$ with $\Sigma$ g solutions for which $\Sigma \notin \Gamma$.
2.1.3. We develop modeling that associates with intrinsic electromagnetic and gravitational properties of objects and with aspects of electromagnetic and gravitational fields.
We explore the notion that some solutions that table 2 lists associate with long-range interactions (or, LRI) and with properties - of physical objects such as planets or elementary particles - that people do infer or might infer via observations based on information carried by electromagnetic fields and gravitational fields.

Regarding observations - via electromagnetism - pertaining to an object with nonzero charge, people might infer both a size of a charge of the object and a velocity with which the object moves. (Other inferences, such as the magnetic moment of the object might also pertain.) We associate charge with a notion of intrinsic property and velocity with a notion of extrinsic property.

We deploy the symbol PROP to associate with $\Sigma g \Gamma$ solutions that we associate with intrinsic properties of objects. We deploy the symbol CURR to associate with $\Sigma g \Gamma$ solutions that we associate with currents of properties.

We explore the notion that some solutions that table 2 lists associate with intrinsic properties - such as charge - of objects. (Later, we explore extrinsic properties such as velocity.)

We assume the following associations. 1 g associates with electromagnetism. 2 g associates with gravitation. Each $\Sigma g \Gamma$ solution (or, $\Sigma=\left|l_{\Sigma}\right|=$ $\left|\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)\right|$ solution) associates with two $l_{\Sigma}=$ $\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)$ expressions. We associate $l_{\Sigma}>0$ with left-circular polarization. We associate $l_{\Sigma}<0$ with right-circular polarization. (We anticipate that, for non-LRI elementary particles, the notion of two expressions per solution associates with two handednesses. See discussion related to table 8 . For LRI elementary particles, it is not necessarily inappropriate to associate left-circular polarization with left-handedness and right-circular polarization with right-handedness.)

Table 3 discusses interpretations - regarding properties of an object - regarding $\Sigma g^{\prime}$ solutions for which $1 \leq \Sigma \leq 2$ and $1 \leq l_{\max } \leq 4$.

We anticipate the following notions. For gravity produced by an object like the Sun, $2 g^{\prime}$ solutions other than 2 g 2 associate with adjustments with respect to the gravity that would associate with a mass that is both spherically symmetric and nonrotating. Regarding large-scale gravitation, 2 g ' solutions other than 2 g 2 can associate with gravitational effects that dominate gravitational effects that would associate with the combination of spherically symmetric and non-rotating.

Table 3 suggests two uses for the words monopole, dipole, quadrupole, and octupole. One use associates with mathematics and with table 2. (This use does not necessarily associate directly with physics.) One use associates with physics and with the dependence of potentials that associate with modeling regarding components of LRI-centric interactions (or, LRI forces). (Reference [4] discusses a similar application regarding gravitation and modeling that associates with general relativity. Reference [5] discusses an application - of notions of monopole, dipole, and so forth - regarding acoustics.)

We assume that a solution for which $\Sigma \geq 1$ and for which the notion of $\Sigma g^{\prime}$ pertains associates with an RDP of the form $\Xi^{-n_{\Sigma g \Gamma}}$. RDP stands for radial dependence of potential. Here, we consider Newtonian modeling for potentials (as in potential energy) that associate with fields (such as the electromagnetic field and the gravitational field) that an object produces. For a solution other than a monopole solution, the potential can (and generally does) vary based on angular coordinates (as well as based on a radial coordinate). We assume that $\Xi^{-1}=r^{-1}$, in which $r$ is the spatial distance from the object. (We provide a cautionary note regarding terminology. Per table 3, we associate the solution for which $\Sigma$ is one and $\Gamma$ is $1^{\prime} 2^{‘} 4$ with each one of the following: $\Xi^{-3}$ and hence mathematical quadrupole, $r^{-3}$ and hence a behavior of potential that associates with a notion of quadrupole, and a physics object
that associates with a magnetic dipole for which the axis rotates around an axis that does not equal the axis that associates with the magnetic dipole. One way to think about the seeming tension between quadrupole and dipole is to associate the factor $\Xi^{-1}$ that associates with $l_{o}=4$ with $(c t)^{-1}$ instead of with $r^{-1}$. Here, $c$ denotes the speed of light and $t$ denotes the time that light takes to go from the magnetic-dipole object to the distance $r$ from the object. This interpretation has consistency with the notion that the relevant quadrupole component of the electromagnetic field associates with an object that people might characterize as having the properties of a magnetic dipole.)

### 2.1.4. We extend our modeling to include extrinsic properties of objects.

We anticipate extending the notions of PROP and CURR to apply widely regarding modeling regarding LRI. We anticipate that, for each one of most LRI PROP solutions, there is an LRI CURR solution.

Notions of three degrees of freedom seem to pertain regarding solutions that table 3 shows.

The following examples - of three degrees of freedom - pertain regarding $1 g^{\prime}$ solutions. Regarding 1 g 1 '2, three degrees of freedom pertain. Two degrees of freedom associate with the orientation of the magnetic moment 3 -vector. One degree of freedom associates with the magnitude of the magnetic moment 3 -vector. Compared to $1 \mathrm{~g} 1^{‘} 2,1 \mathrm{~g} 1^{〔} 2^{‘} 4$ has three more degrees of freedom. Two degrees of freedom associate with the orientation of the angular velocity 3 -vector. One degree of freedom associates with the magnitude of the angular velocity 3 -vector.

Regarding each of the solutions that table 3 shows, $l_{o}=4$ seems to associate - regarding rotation - with three degrees of freedom.

To explore CURR solutions, we want to add three degrees of freedom that associate with the CURR aspects of PROP-and-CURR 4 -somes. (We use the one-element term 4 -some and not the oneelement term 4 -vector. For modeling centric to special relativity and scalar PROP aspects, the notion of 4 -vector might be appropriate. However, the following notions pertain. This essay discusses PROP aspects that do not necessarily associate with the notion of scalar. This essay includes CURR aspects that - depending on choices regarding kinematics modeling - do not necessarily need to comport with special relativity.) We assume that modeling - regarding the three degrees of freedom that we want to add - associates with one $l_{o}$.

The limit $l_{\text {max }} \leq 4$ does not necessarily allow for enough degrees of freedom. For example, regarding the PROP solutions $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 4 \mathrm{x}$ (with x denoting either v or w ), one could not have CURR solutions.

We assume that $l_{o}=6$ and $l_{o}=8$ have rel-

Table 3: Interpretations - regarding properties of an object - regarding $\Sigma g^{\prime}$ solutions for which $1 \leq \Sigma \leq 2$ and $1 \leq l_{\text {max }} \leq 4$. We assume the following notions. 1 g 1 associates with a component - of the electromagnetic field that the object produces - that associates with the object's charge. The word scalar associates with this solution. 1g1'2 associates with the object's magnetic field. An axis associates with that field. The one-element term 3 -vector associates with this solution. (We assume that, for a bar magnet, an association with 1 g 1 ' 2 pertains, even though the notions of charge and rotation do not necessarily pertain.) $1 \mathrm{~g} 1 ‘ 2 ‘ 4$ associates with a combination of magnetic field and rotation (over time) of the axis of the magnetic field. (The Earth is an object for which the axis of the magnetic field rotates. That rotation associates with the notion that the axis of the Earth's rotation does not equal the axis of the magnetic field.) The one-element term 3-vector associates with that rotation of the axis that associates with the magnetic field. Overall, two 3 -vectors pertain regarding 1 g 1 '2'4. 2 g 2 associates with the object's mass. The word scalar associates with this solution. 2 g 2 ' 4 associates with rotation of the object's mass. An axis associates with that rotation. The one-element term 3-vector associates with this solution. (Regarding general relativity, this solution associates with aspects of rotational frame dragging.) 2 g 1 ' 2 ‘ 3 associates with a non-spherically-symmetric distribution of mass. The one-element term 3-vector associates with each of a possible - axis and magnitude of a - minimal moment of inertia and a possible - axis and magnitude of a - maximal moment of inertia. $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{v}$ associates with rotation of the minor axis of moment of inertia. The one-element term 3-vector associates with that rotation. $2 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 4 \mathrm{w}$ associates with rotation of the major axis of moment of inertia. The one-element term 3 -vector associates with that rotation. (Regarding general relativity, each of $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{v}$ and $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{w}$ might associate with aspects of rotational frame dragging.)

| $\Sigma$ | Monopole | Dipole | Quadrupole | Octupole |
| :--- | :--- | :--- | :--- | :--- |
| 1 | 1 g 1 | $1 \mathrm{~g} 1^{‘} 2$ | $1 \mathrm{~g} 1^{‘} 2^{‘} 4$ | - |
| 2 | 2 g 2 | $2 \mathrm{~g} 2^{‘} 4$ | $2 \mathrm{~g} 1^{‘} 2^{‘} 3$ | $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 4 \mathrm{v}, 2 \mathrm{~g} 1^{‘} 2^{‘} 33^{‘} 4 \mathrm{w}$ |

evance. For example, either one of $1 g 1^{〔} 2^{‘} 4^{〔} 6$ and $1 \mathrm{~g} 1 \cdot 2^{\prime} 4^{\prime} 8$ could be a CURR solution that associates with the PROP solution $1 g 1^{\prime} 2^{‘} 4$. We assume that - overall (except regarding a possibility regarding dark energy density of the universe) $-l_{\max } \leq 8$ pertains.

We assume that each of $l_{o}=5$ and $l_{o}=7$ does not have relevance. This assumption associates with the following discussion. Consider a PROP solution $\Sigma g \Gamma$. Adding an odd value of $l_{o}$ to a list $\Gamma$ produces a list that we denote by $\Gamma^{\prime}$. Arithmetically, for each solution $\Sigma^{\prime} g \Gamma^{\prime}, \Sigma^{\prime}$ cannot equal $\Sigma$. With respect to the PROP solution $\Sigma g \Gamma$, no relevant CURR solutions associate with $\Gamma^{\prime}$. We anticipate that the assumption also associates with a likely limit $-\Sigma \leq 4$ - regarding relevant solutions $\Sigma \mathrm{g} \Gamma$. (See table 17 .)

We discuss notions regarding $l_{o}=3.1 \mathrm{~g} 1^{\prime} 2$ is a PROP solution that associates with intrinsic (nominal) magnetic moment. (See table 3.) Here, $\Sigma=1$ and $\Gamma=1^{\prime} 2$ pertain. For $\Gamma=1^{\prime} 2, \Sigma=3$ can pertain. We anticipate that a PROP $3 \mathrm{~g} 1 ‘ 2$ solution associates with intrinsic anomalous magnetic moment. (See table 5.) People measure anomalous magnetic moments for charged leptons (which are elementary fermions). For charged leptons, anomalous magnetic moments vary with fermion flavour. We anticipate that, for some modeling, this essay associates - at least indirectly - $l_{o}=3$ with (at least) the property of flavour for leptons. (The word leptons associates with some - but not all - elementary fermions. See discussion related to equation (4). This essay de-emphasizes, but does not rule out, notions that similar modeling associating with $l_{o}=3 \mathrm{might}$ associate with aspects beyond lepton flavour.)

We discuss notions regarding $l_{o}=6$ and $l_{o}=8$. We anticipate that - for some SRI solutions $-l_{o}=$

6 can associate with interactions with elementary fermions and $l_{o}=8$ can associate with interactions with elementary bosons. (See table 8,) We anticipate that - for elementary fermions $-6 \in \Gamma$ pertains regarding PROP solutions. (See table 9.)

### 2.1.5. We discuss modeling that might associate

 with PROP solutions for which $8 \in \Gamma$.We anticipate modeling - in which some $\Sigma=0$ solutions associate with simple elementary particles (or, SRI elementary bosons and ELF elementary fermions) - in which (for at least the known elementary particles) the following notions might pertain regarding some quantum kinematics modeling. (See table 8 and table 9.) Simple particles for which PROP modeling associates with $8 \notin \Gamma$ can model (for some circumstances) as not entangled. Simple particles for which PROP modeling associates with $8 \in \Gamma$ seemingly do not model (for any circumstances) as not entangled.

For convenience, we associate the word free with aspects that associate with modeling that associates with PROP solutions for which $8 \notin \Gamma$. For convenience, we associate the word entwined with aspects that associate with modeling that associates with PROP solutions for which $8 \in \Gamma$. We choose the word entwined so as to avoid using the word entangled. (In some circumstances, free elementary particles model as entangled. An example features an electron that is part of an atom.)

### 2.1.6. We list and discuss $\Sigma g^{\prime}$ solutions that associate with $l_{\max } \leq 8$.

Table 4 extends table 3 and lists PROP and CURR $\Sigma g^{\prime}$ solutions, for which $1 \leq \Sigma \leq 4$ and $l_{\max } \leq 8$. We anticipate discussing the extent to which $3 \mathrm{~g} \Gamma$ solutions might associate with an elementary particle and the extent to which $4 g \Gamma$ solu-

Table 4：PROP and CURR $\Sigma g^{\prime}$ solutions，for which $1 \leq \Sigma \leq 4$ and $l_{\max } \leq 8$ ．The symbol $n_{\Gamma, \text { PROP }}$ denotes the number of elements in the $\Gamma$ that associates with PROP．Table 4 lists an RDF－or radial dependence of force－for each PROP solution． For a CURR $\Sigma g \Gamma$ solution，the RDP（or，radial dependence of potential）equals $\Xi^{-1}$ times the RDP for the associated PROP $\Sigma g \Gamma$ solution．For each one of PROP and CURR，the RDF equals $\Xi^{-1}$ times the RDP．For example，for each of $1 g 1$ and 2 g 2 ，the RDF is $\Xi^{-2}$ ，which is $r^{-2}$ ．An $\mathrm{x}-\mathrm{as}$ in $\Sigma \mathrm{g} \Gamma \mathrm{x}$－denotes the notion that more than one solution pertains．Table 4 shows properties－of objects that associate with $\Sigma \mathrm{g} \Gamma$ components of $\Sigma \mathrm{g}$－that associate with the PROP solution．The table attempts to use familiar symbols．The table associates the symbols with phrases．$S$ denotes spin，as in the expression $S(S+1) \hbar^{2}$ ．The symbol TBD denotes at least one of the three－word phrase to be determined and the notion that this essay discusses the aspect elsewhere．

| $\Sigma$ | $\Sigma \mathrm{g} \Gamma$ PROP | $\Sigma \mathrm{g} \Gamma$ CURR | $n_{\Gamma, \mathrm{PROP}}$ | RDF PROP | PROP－associated properties |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1g1 | 1g1＇2 | 1 | $\Xi^{-2}$ | $q$－Charge |
| 1 | $1 \mathrm{~g} 1 \times 2$ | $1 \mathrm{~g} 1^{\prime}{ }^{\prime} 4$ | 2 | $\Xi^{-3}$ | $\mu$－Magnetic dipole moment （including from rotating charge） |
| 1 | $1 \mathrm{~g} 1^{\prime}{ }^{\prime} 4$ | $\begin{aligned} & \lg 1^{\prime} 2^{‘} 4^{‘} 6, \\ & 1 \mathrm{~g} 1^{‘} 2^{\prime} 4^{\prime} 8 \end{aligned}$ | 3 | $\Xi^{-4}$ | $\mu, \omega$－Magnetic dipole moment and internal angular velocity |
| 1 | 1g1＇2‘4 ${ }^{6} \mathrm{x}$ | $1 \mathrm{~g} 1 \cdot 2^{\prime} 4^{6} 6^{\prime} 8 \mathrm{x}$ | 4 | $\Xi^{-5}$ | TBD |
| 1 | 1g1＇2＇4＇8 | $1 \mathrm{~g} 1^{\prime} \mathrm{C}^{\prime} 4^{6}{ }^{\text {¢ }} 8 \mathrm{x}$ | 4 | $\Xi^{-5}$ | TBD |
| 1 | $1 \mathrm{~g} 1^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ | None | 5 | $\Xi^{-6}$ | TBD |
| 1 | $1 \mathrm{~g} 1^{4}{ }^{\circ} 6$ | $\begin{aligned} & \lg 1^{‘} 2^{‘} 4^{6} 6 x \\ & \lg 1^{‘} 4^{‘} 6^{‘} 8 \end{aligned}$ | 3 | $\Xi^{-4}$ | TBD |
| 1 | $\lg 1 \times 4 \times 6 \times 8$ | $1 \mathrm{~g} 12^{\prime} 4^{6} 6^{\prime} 8 \mathrm{x}$ | 4 | $\Xi^{-5}$ | TBD |
| 1 | $1 \mathrm{~g} 1 \times 6 \times 8$ |  | 3 | $\Xi^{-4}$ | TBD |
| 1 | $1 \mathrm{~g} 1^{\prime} 2^{\prime} 6^{\prime} 8 \mathrm{x}$ | $1 \mathrm{~g} 1 \cdot 2^{\prime} 4^{6} 6{ }^{\text {8 }} \mathrm{x}$ | 4 | $\Xi^{-5}$ | TBD |
| 2 | 2g2 | 2g2 4 | 1 | $\Xi^{-2}$ | $m$－Mass |
| 2 | 2g2 ${ }^{\text {4 }}$ | $2 \mathrm{~g} 2^{\prime} 4^{\prime} 8$ | 2 | $\Xi^{-3}$ | $m, \omega$－Rotating（spherically symmetric aspects of）mass |
| 2 | 2g2＇4＊8 | None | 3 | $\Xi^{-4}$ | TBD |
| 2 | $2 \mathrm{~g} 1^{\prime}{ }^{\prime} 3$ | $\begin{aligned} & 2 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 4 \mathrm{v} \\ & 2 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 4 \mathrm{w}, \\ & 2 \mathrm{~g} 1^{\prime} 2^{‘} 6 \\ & 2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{‘} 8 \end{aligned}$ | 3 | $\Xi^{-4}$ | $I_{C}$－Moments of inertia |
| 2 | $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3 \times 4 \mathrm{v}$ | $\begin{aligned} & 2 \mathrm{~g}^{‘} 2^{‘} 3^{‘} 4^{{f300082d8-a683-488c-9910-c058b9049e34}} 8 \mathrm{x} \\ & 2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{‘} 6^{‘} 8 \mathrm{x} \end{aligned}$ | 4 | $\Xi^{-5}$ | TBD |
| 2 | 2g1＇2‘3 ${ }^{4}{ }^{\text {c }}$ 6x | $2 \mathrm{~g} 12^{\prime} 3^{\prime} 4^{6} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-6}$ | TBD |
| 2 | $2 \mathrm{~g} 1^{\prime} 3^{\prime} 3^{\prime}{ }^{\prime} 8 \mathrm{x}$ | $2 \mathrm{~g} 1{ }^{\prime} 3^{〔} 4^{6} 6{ }^{\text {c }} 8 \mathrm{x}$ | 5 | $\Xi^{-6}$ | TBD |
| 2 | $2 \mathrm{~g} 1^{\prime} 3^{\prime} 3^{\prime}{ }^{\prime} 8 \mathrm{x}$ | $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-6}$ | TBD |
| 2 | $2 \mathrm{~g} 1^{\prime}{ }^{\prime} 3 \cdot 4^{\prime} 6^{\prime} 8 \mathrm{x}$ | None | 6 | $\Xi^{-7}$ | TBD |
| 3 | 3 g 3 | $3 \mathrm{~g} 3^{6} 6$ | 1 | $\Xi^{-2}$ | TBD（a function of elementary fermion flavour） |
| 3 | $3 \mathrm{3} \times 6$ | None | 2 | $\Xi^{-3}$ | TBD |
| 3 | $3 \mathrm{~g} 2^{\prime}{ }^{〔} 4$ | $\begin{aligned} & 3 \mathrm{~g} 2^{‘} 3^{‘} 4^{‘} 6 \mathrm{x} \\ & 3 \mathrm{~g} 2^{‘} 3^{‘} 4^{‘} 8 \end{aligned}$ | 3 | $\Xi^{-4}$ | TBD |
| 3 | $3 \mathrm{~g} 2 \times 3 \times 4 \times \mathrm{x}$ | $3 \mathrm{~g} 2 \times 3 \times 4^{6} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-5}$ | TBD |
| 3 | 3 g 2 ＇3＇4＇8 | $3 \mathrm{~g} 2 \cdot 3 \cdot 4^{6} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-5}$ | TBD |
| 3 | $3 \mathrm{~g} 2^{\prime} 3^{4} 6^{6} 8 \mathrm{x}$ | None | 6 | $\Xi^{-7}$ | TBD |
| 4 | 4 g 4 | $4 \mathrm{~g} 4^{\text {¢ }}$ | 1 | $\Xi^{-2}$ | $S$－Angular momentum（scalar quantity） |
| 4 | $4 \mathrm{~g} 4^{\text {¢ }} 8$ | None | 2 | $\Xi^{-3}$ | TBD |
| 4 | $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{x}$ | $\begin{aligned} & 4 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 46 \mathrm{x} \\ & 4 \mathrm{~g} 1^{‘} 2^{\prime} 3^{‘} 4^{\wedge} 8 \mathrm{x} \end{aligned}$ | 4 | $\Xi^{-5}$ | TBD |
| 4 | $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{4} 4^{6} 6 \mathrm{x}$ | $4 \mathrm{~g} 12^{\prime} 3^{\prime} 4^{6} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-6}$ | TBD |
| 4 | $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 8 \mathrm{x}$ | $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{6} 6^{\prime} 8 \mathrm{x}$ | 5 | $\Xi^{-6}$ | TBD |
| 4 | $4 \mathrm{~g} 1 \cdot 2^{\prime} 3 \cdot 4^{\prime} 6^{\prime} 8 \mathrm{x}$ | None | 6 | $\Xi^{-7}$ | TBD |

tions might associate with an elementary particle． （See discussion related to table 11．）

## 2．1．7．We discuss a notion of cascades of solutions．

In table 4 for each CURR solution，the list $\Gamma$ includes each of the elements from the list $\Gamma$ for the associated PROP solution．Each CURR $\Gamma$ includes exactly one element that does not occur in the $\Gamma$ for the associated PROP solution．Such a CURR solution associates with the notion of associating－ with the relevant PROP solution－a PROP－and－ CURR 4－some．The three CURR components asso－ ciate with notions of velocity．

We say that each such CURR solution cascades from the associated PROP solution．

For example，CURR 1g1＇2 associates with linear movement of charge（which associates with 1 g 1 ）and （for some modeling）with a contribution to a mag－ netic field．Also，PROP solution $1 \mathrm{~g} 1 ‘ 2$（occurs fur－ ther down in the table and，for some modeling）as－ sociates with rotating charge（for which a notion of angular velocity can be appropriate）and（for some modeling）with a contribution to a magnetic field． PROP 1g1‘2 also associates with effects of a bar magnet（in which a notion of angular velocity is not necessarily appropriate）and（for some model－ ing）with a contribution to a magnetic field．

We extend the notion of cascading to include the notion that a solution that appears as a CURR solution appears（generally further down）in table 4 as a PROP solution．The notion of velocity that pertains for the CURR use of the solution associates with a notion of angular velocity for the PROP use of the solution．

In table 4 the $1 g^{\prime}$ cascade starters are 1 g 1 ， $1 \mathrm{~g} 1^{\prime} 4^{‘} 6$ ，and $1 \mathrm{~g} 1^{‘} 6^{\prime} 8$ ．The $2 g^{\prime}$ cascade starters are 2 g 2 and $2 \mathrm{~g} 1^{\prime} 2^{‘} 3$ ．The 3 g ＇cascade starters are 3 g 3 and $3 \mathrm{~g} 2^{‘} 3^{\circ} 4$ ．The $4 g^{\prime}$ cascade starters are $4 g 4$ and $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{x}$ ．

## 2．1．8．We discuss Og solutions－the three Og solu－

 tions in table 2 and other $0 g$ solutions that cascade from the three Og solutions in table 2 －that associate with all known non－LRI el－ ementary particles and all non－LRI elemen－ tary particles that we suggest．We anticipate the following three associations－ the $0 \mathrm{~g} 1^{‘} 2^{‘} 3$ solution and the W boson，the $0 \mathrm{~g} 1^{‘} 3^{‘} 4$ solution and the Z boson，and the $0 \mathrm{~g} 1^{‘} 2^{‘} 3^{〔} 4$ solution and the Higgs boson．（See table 2 and table 8．）

We anticipate that other non－LRI elementary particles－including each non－LRI elementary par－ ticle of which people know and each non－LRI el－ ementary particle that we suggest－associate with solutions that cascade from the set of solutions that consists of the $0 \mathrm{~g} 1^{\prime} 2^{\prime} 3$ solution，the $0 \mathrm{~g} 1^{\prime} 3^{‘} 4$ solution， and the 0 g $1 ‘{ }^{\prime} \times 3^{‘} 4$ solution．（See table 8 and table 9．）

## 2．1．9．We discuss modeling that might associate

 with $\Sigma g^{\prime}$ cascades．Regarding an elementary particle，the cascades that associate with $1 \mathrm{~g} 1,2 \mathrm{~g} 2,3 \mathrm{~g} 3$ ，and 4 g 4 might be adequate to associate with relevant nominal in－ trinsic properties．These cascades associate with charge，nominal magnetic moment，mass，flavour （for elementary fermions），and spin．The cas－ cades that associate with $1 \mathrm{~g} 1^{〔} 4^{〔} 6,1 \mathrm{~g} 1^{〔} 6^{〔} 8,2 \mathrm{~g} 1^{〔} 2^{〔} 3$ ， $3 \mathrm{~g} 2^{\prime} 3^{‘} 4$ ，and $4 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 4 \mathrm{x}$ might have little relevance regarding modeling for intrinsic properties of ele－ mentary particles．We anticipate that some cas－ cades that associate with $\Sigma g^{\prime \prime}$ solutions might have relevance regarding modeling for anomalous intrin－ sic properties．（Regarding anomalous magnetic mo－ ments，see table 5．）

For an object（such as a proton or a galaxy） that includes more than one elementary particle， the cascades that associate with 1 g 1 and 2 g 2 can be relevant regarding modeling．The cascades that associate with 3 g 3 and 4 g 4 might not have（much or any）relevance．The cascades that associate with $1 \mathrm{~g} 1^{‘} 4^{〔} 6$ and $1 \mathrm{~g} 1^{‘} 6^{‘} 8$ might have relevance regard－ ing modeling that associates with notions of orbital angular momentum and spin angular momentum． The cascade that associates with $2 \mathrm{~g} 1{ }^{\prime} 2^{\prime} 3$ might have relevance regarding modeling that associates with notions of stress－energy．This essay does not ad－ dress the topic of the possible relevance of mod－ eling that would associate with the cascades that associate with $3 \mathrm{~g} 2^{\prime} 3^{\prime} 4$ or with $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{‘} 4 \mathrm{x}$ ．

The end of each $\Sigma g^{\prime}$ cascade associates with a row in table 4 for which no CURR solutions per－ tain．（The solutions $1 \mathrm{~g} 1^{‘} 2^{\circ} 4^{〔} 6^{\circ} 8 \mathrm{x}$ end the cascade that associates with 1 g 1 ，the cascade that associates with $1 \mathrm{~g} 1^{〔} 4^{〔} 6$ ，and the cascade that associates with 1 g1＇6‘8．）Regarding such a row，the PROP solution might associate with modeling regarding changes to properties that associate with the relevant cas－ cade．For example， $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{‘} 4^{〔} 66^{‘} 8 \mathrm{x}$ might associate with notions that the energy or the stress－energy can change．

## 2．1．10．We list and discuss $\Sigma g$＂solutions that might associate with anomalous properties．

Table 5 lists some $\Sigma g^{\prime \prime}$ solutions．The PROP so－ lution $3 \mathrm{~g} 1^{〔} 2$ associates with the 1 g ＇solution $1 \mathrm{~g} 1^{〔} 2$ ， which associates with the property of nominal mag－ netic moment．Table 5 previews the notion that 3g1＇2 associates with modeling regarding anoma－ lous magnetic moment．（See table 10 and equation （4）．）The PROP solution $4 g 1 ' 2 ‘ 3$ associates with the 2 g ＇solution $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3$ ，which associates with non－ spherically－symmetric distribution of mass．Ta－ ble 5 previews the notion that $4 \mathrm{~g} 1 ‘ 2^{‘} 3$ might asso－ ciate with modeling regarding（anomalous）neutrino masses or regarding neutrino mass mixing．（See ta－ ble 10 and table 15 ）

Table 5：Some $\Sigma g$＂solutions．The table lists some $\Sigma \mathrm{g}$＂PROP solutions．Each one of the PROP solutions 3 g 1 ＇ 2 ‘ 4 and 3 g 1 ＇2＇ 6 cascades from the PROP solution 3 g 1 ＇ 2 ．

| $\Sigma$ | PROP $\Gamma$ | PROP $\Sigma=\ldots$ | CURR $\Sigma=\ldots$ | $n_{\Gamma, \text { PROP }}$ | Association |
| :--- | :--- | :--- | :--- | :--- | :--- |
| 3 | $1^{‘} 2$ | $\|+1+2\|$ | $\|+1-2+4\|$, | 2 | Anomalous magnetic moment |
|  |  |  | $\|-1-2+6\|$ |  |  |
| 3 | $1^{{f8b7af74c-6682-447f-a414-f3372f1de717}} 2^{‘} 3$ | $\|-1+2+3\|$ | $\|-1-2+3+4\|$, | 3 | Anomalous gravitational property |
|  |  |  | $\|+1+2-3+4\|$, |  |  |
|  |  |  | $\|-1+2-3+6\|$, |  |  |
|  |  |  | $\|+1-2-3+8\|$ |  |  |

For each PROP $\Gamma$ that table 5 shows， $1 \in \Gamma$ and $2 \in \Gamma$ ．A term such as the two－word term anoma－ lous property might be more appropriate than a possible distinction between the three－word term anomalous magnetic moment and the three－word term anomalous gravitational property．

## 2．1．11．We provide perspective about modeling that associates with attractive and repulsive components of gravity．

We explore the extent to which each 2 g ＇PROP solution associates with gravitational attraction and the extent to which each $2 g^{\prime}$ PROP solution associates with gravitational repulsion．We suggest results，based on generalizing from a series of cases．

We consider modeling regarding an object－ob－ ject A－and the gravity that object A produces．We consider an object B that interacts with the gravity that object A produces．We assume that each one of object A and object B does not exhibit transla－ tional motion．

We include cases for which PROP solutions can pertain for components of object A or for the entire object．

One case associates with object A associating with the notion of a point mass．The PROP $l_{\max }$ is two．Object A has no internal components．The CURR solution $2 \mathrm{~g} 2^{\prime} 4$ associates with the motion of object A and，per an assumption，is zero．

One case associates with the notion（regarding object A）of a non－rotating spherically symmetric distribution of mass．Object A has internal compo－ nents．For each component，we assume that mod－ eling based on a PROP $l_{\max }$ of two suffices．The 2 g 2 solution pertains．The CURR solution $2 \mathrm{~g} 2^{〔} 4$ as－ sociates with motion of the component．Across the components，the motions of components might help to overcome internal gravitational collapse．（For ex－ ample，for some objects，one might consider that the motions associate with thermal energy．）Across the components，the contributions to the CURR 2g2‘4 solution relevant to object A average to zero．Ob－ ject B senses no first－order effects that would asso－ ciate with the CURR $2 \mathrm{~g} 2^{‘} 4$ solution that associates with object A．

One case associates with the notion of a non－ rotating non－spherically－symmetric distribution of mass．Here，we assume that three new（compared to the previous case）degrees of freedom associate with a magnitude and axis that associate with a maximal moment of inertia for object A and that three other new（compared to the previous case）degrees of free－ dom might associate with a magnitude and axis that associate with a minimal moment of inertia for ob－ ject A．The PROP solutions 2 g 2 and $2 \mathrm{~g} 1 ‘ 2^{‘} 3$ per－ tain regarding object A．Over time，object A might evolve to become more spherically symmetric．En－ ergy that associates with having maintained at least one non－zero moment of inertia（which associates with $2 g 1^{〔} 2^{〔} 3$ ）would drain from the $2 g 2$ for object A．Compared to the previous case，this case illus－ trates the notion that the relevance of the $2 \mathrm{~g} 1^{‘} 2^{`} 3$ solution associates with increased（compared to the previous case） 2 g 2 and therefore with more（com－ pared to the previous case）gravitational attraction （as experienced by object B）．

The previous two cases illustrate the notion that a 2 g ＇PROP $\Gamma$ solution for which $n_{\Gamma, \mathrm{PROP}}=3$ as－ sociates with gravitational attraction．

One case associates with the notion of a uni－ formly rotating spherically symmetric distribution of mass．For this case，one angular velocity $\omega$ per－ tains．The angular velocity is with respect to an axis that runs through the center of object A．The angular velocity pertains regarding each component of object A．The angular velocity associates with three degrees of freedom－two of which associate with the axis that associates with $\omega$ and one of which associates with the magnitude of $\omega$ ．The PROP solutions 2 g 2 and $2 \mathrm{~g} 2^{〔} 4$ pertain．The CURR solutions $2 \mathrm{~g}^{〔} 4$ and $2 \mathrm{~g} 2^{‘} 4^{‘} 8$ pertain．From the per－ spective of an object $B$ that does not lie on an exten－ sion of the axis，effects of the values of the $2 \mathrm{~g} 2^{〔} 4^{〔} 8$ for the components do not necessarily sum to zero． If the object $A$ existed in nature，object $A$ might tend（over time）to become oblate．The transition process would release－from 2 g 2 －energy that，in effect，had maintained the spherical symmetry．

One case associates with the notion of such a rotating oblate object．

The last two cases illustrate the notion that，for modeling based on 2 g＇$^{\prime}$ PROP solutions，$n_{\Gamma, \mathrm{PROP}}=$ 2 associates with reducing gravitational effects and thereby associates with gravitational repulsion．

We extrapolate．
Table 6 pertains regarding modeling that has bases in 2g＇PROP solutions．

## 2．1．12．We discuss possible limits on the applicabil－ ity of $\Sigma g^{\prime}$ solutions．

This essay suggests that－regarding $\Sigma \mathrm{g}^{\prime}$ solu－ tions that might have relevance regarding modeling for LRI－that $\Sigma$ does not exceed four．The sugges－ tion associates with an extrapolation regarding data pertaining to the relative strengths of electromag－ netism and gravity．（See table 17．）That extrapo－ lation suggests that a 5g5 solution would associate with a force strength of zero．We assume that，if $5 g 5$ associates with zero force strength，each $\Sigma g^{\prime}$ so－ lution for which $\Sigma \geq 5$ is not relevant regarding LRI physics．The suggestion also associates with the no－ tion that－for PROP solutions $\Sigma g \Gamma$ for which $\Sigma \neq 0$ $-5 \notin \Gamma$ ．（See table 4）The suggestion also asso－ ciates with the notion that－regarding the reaches of instances of $\Sigma$ g＇solutions－a maximum value for the integer $n_{0}$ might be three．（See discussion re－ garding equation（3）．）The value for $n_{0}$ for the 5 g 5 solution would be four．

We discuss $3 g^{\prime}$ solutions and $4 g^{\prime}$ solutions．
Regarding 3g＇，table 4 suggests two cascades． One cascade starts with the $3 g 3$ PROP solution． We suggest－regarding the $3 \mathrm{~g} 3{ }^{〔} 6$ PROP solution －that the lack of a partner CURR solution asso－ ciates with a notion that non－zero effects that as－ sociate with 3 g 3 associate with a lack of structure that can exhibit nonzero angular velocity．We as－ sume that 3 g 3 associates with interactions with（no more than）elementary particles．We assume that the 6 in the $\Gamma$ for the CURR $3 g 3{ }^{〔} 6$ solution asso－ ciates with three flavors and not necessarily with three degrees of freedom regarding motion．We as－ sume that a 3 g 3 interaction would associate with the flavour of one elementary fermion．Based on such notions，we think that the cascade that starts with $3 \mathrm{~g} 2^{‘} 3^{‘} 4$ does not have relevance for the physics modeling that this essay discusses．

Regarding 4g＇，table 4 suggests two cascades． Reasoning similar to reasoning that we just used regarding $3 g^{\prime}$ suggests that the cascade that begins with the $4 g 1^{\prime} 2^{\prime} 3 ‘ 4 \mathrm{x}$ PROP solutions is not relevant regarding our physics modeling．Also， 4 g 4 might associate with interactions with single non－LRI el－ ementary particles．Thus，a 4 g 4 interaction might associate with（no more than）the scalar $S$ that as－ sociates with one elementary particle．（ $S$ denotes spin，as in the expression $S(S+1) \hbar^{2}$ ．）Possibly， a notion that three values associate with the 8 in the $\Gamma$ for the CURR solution $4 g 4^{〔} 8$ associates with
three allowed spins－（in units of $\hbar$ ）zero，one－half， and one－regarding simple elementary particles．

We discuss possible limits regarding $2{ }^{\prime}$＇solu－ tions．

We consider two objects that are some distance apart from each other．We consider doubling linear dimensions－that is doubling the distance between the objects and doubling the diameters of the ob－ jects－while maintaining，for each object，a constant mass per unit volume．A PROP RDF $\Xi^{-6}$ force after the doubling of linear dimensions equals the PROP RDF $\Xi^{-6}$ force before the doubling of lin－ ear dimensions．Possibly，this invariance regarding scaling suggests reasons not to pursue－regarding interactions between pairs of large objects－model－ ing regarding PROP RDF $\Xi^{-l_{r}}$ for which $l_{r}$ exceeds six．Such a limit seems to be consistent with the no－ tion－that table 4 shows－that no CURR solution partners with $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ PROP solutions．

Regarding 2g＇，table 6 associates PROP RDF $\Xi^{-6}$ solutions with gravitational attraction．PROP RDF $\Xi^{-6}$ solutions might associate with a property that associates with energy．（See table 10．）Possi－ bly，the PROP RDF $\Xi^{-7}$ solutions $2 g 1^{\prime} 2^{\prime} 3^{‘} 4^{‘} 6^{〔} 8 \mathrm{x}$ associate with repulsion．（See table 6．）Possibly， the PROP RDF $\Xi^{-7}$ solutions $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{4} 4^{〔} 6^{‘} 8 \mathrm{x}$ asso－ ciate with changes in the energies of objects．

This essay notes－but does not much discuss－ the notion that the PROP solutions $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3 \cdot 4^{\prime} 6^{\prime} 8 \mathrm{x}$ might associate with spontaneous decay by elemen－ tary particles．（See table 10 ．This essay notes－ but does not much discuss－the notions that the PROP solutions $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 4^{〔} 6^{‘} 8 \mathrm{x}$ might associate with repulsion within a small object and might－along with solutions $1 \mathrm{~g} 1^{\prime} 2^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$－associate with aspects such as thermal radiation．Also，this essay does not further discuss possible associations of the notion of $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ repulsion with notions such as en－ tropy and arrow of time．）

## 2．2．Elementary particles

This unit develops and deploys modeling that points to all elementary particles of which people know and to all other elementary particles that our work suggests．

## 2．2．1．We discuss symbols for families of elemen－

 tary particles．We use the following method to catalog elemen－ tary particles．（We anticipate that each one of table 7．table 8，table 9，and table 11 shows uses of the method．）A symbol of the form $S \Phi$ associates with a so－called family of elementary particles．Each el－ ementary particle associates with one family．Each family associates with one of one，three，or eight ele－ mentary particles．For a family，the value $S$ denotes the spin（in units of $\hbar$ ）for each elementary parti－ cle in the family．$S$ associates with the expression

Table 6: 2g' PROP solutions and the extents to which gravitational attraction and gravitational repulsion pertain. An $x$ - as in $\Sigma g \Gamma x$ - denotes the notion that more than one solution pertains. The symbol $\dagger$ associates with the notion that no CURR solutions pertain.

| $2 \mathrm{~g}^{\prime} n_{\Gamma, \mathrm{PROP}}$ | Gravitational ... | RDF PROP | N-pole | Examples of 2g' PROP solutions |
| :---: | :---: | :---: | :---: | :---: |
| 1 | Attraction | $\Xi^{-2}$ | Monopole | 2g2 |
| 2 | Repulsion | $\Xi^{-3}$ | Dipole | 2g2'4 |
| 3 | Attraction | $\Xi^{-4}$ | Quadrupole | 2g1'2‘3 |
| 4 | Repulsion | $\Xi^{-5}$ | Octupole | $2 \mathrm{~g} 1^{\prime} 3^{\prime} 3^{\prime} 4 \mathrm{x}$ |
| 5 | Attraction | $\Xi^{-6}$ | 16-pole | $2 \mathrm{~g} 1 \cdot 2^{\prime} 36^{6} 8 \mathrm{x}$ |
| 6 | Repulsion $\dagger$ | $\Xi^{-7}$ | 32-pole | $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ |

$S(S+1) \hbar^{2}$ that associates with angular momentum. Values of $S$ include $0,0.5$, and 1 and might include 2,3 , and 4 . The symbol $\Phi$ associates with a symbol of the form $\mathrm{X}_{Q}$, in which X is a capital letter and $Q$ is the magnitude of charge (in units of $\left|q_{e}\right|$, in which $q_{e}$ denotes the charge of an electron) for each particle in the family. For cases for which $Q=0$, this essay omits - from the symbols for families the symbol $Q$. (See table 11.)

### 2.2.2. We list known and possible LRI elementary bosons.

Table 7 alludes to sets of solutions that might associate with LRI (or, long-range interaction) elementary particles. We anticipate that discussion related to table 11 indicates that notions of 3 L and 4 L might not necessarily associate with elementary particles.

### 2.2.3. We develop modeling that matches and suggests all elementary particles (other than LRI bosons) that this essay discusses.

We explore the notion that solutions for which $\Sigma=0$ associate with known simple (or, non-LRI) elementary particles and with possible simple elementary particles. Based on arithmetic, for each $\Sigma=0$ solution, $n_{\Gamma}$ is at least three. We associate the symbol SRI (as in short-range interaction or as in elementary boson that does not associate with a long-range interaction) with non-LRI elementary bosons. We associate the symbol ELF (as in elementary fermion) with fermion elementary particles.

For each $\Sigma=0$ solution, there are two expressions of the form $0=l_{\Sigma}=\sum_{l_{o} \in \Gamma}\left( \pm l_{o}\right)$. (See equation (1), equation (2), and discussion related to table 22) For each solution for which $\Sigma>0$, there are two expressions. (See discussion related to table 3, For each solution for which $\Sigma>0$, we associate one expression with left-handedness and with left-circular polarization. For each solution for which $\Sigma>0$, we associate one expression with right-handedness and right-circular polarization. For each $\Sigma=0$ solution, we assume that one expression associates with the notion of lefthandedness and the other expression associates with the notion of right-handedness.

Table 8 shows $0 g \Gamma$ solutions that might associate with elementary bosons that are not LRI bosons. (Reference [6] discusses the inflaton particle.)

Some Standard Model modeling suggests an association between a Higgs field and the notion that some elementary particles have nonzero masses. We anticipate that, for elementary bosons, relationships between masses feature squares of masses. (See table 12,) We anticipate that, for elementary fermions, relationships between masses feature logarithms of masses. (See table 13 and table 14 ) We note - but do not further discuss - the notion that regarding the Higgs boson - differences between the CURR solution $0 \mathrm{~g} 1{ }^{\prime} 2^{‘} 3^{\prime} 4^{\prime} 8$ and the CURR solution $0 \mathrm{~g} 1^{‘} 2^{\prime} 3^{‘} 4^{`} 6$ might associate (from a perspective of modeling) with differences between the two types of relationships between elementary particle masses.

Regarding table 8, we note - but do not discuss further - the notion that the number, one or two, of CURR solutions associates with the extent to which each SRI boson might be its own antiparticle. Regarding the case of just one CURR solution, the following sentence pertains. For each of + and the antiparticle for the $\mathrm{W}^{ \pm}$boson is the $\mathrm{W}^{\mp}$ boson. Regarding each case that associates with two CURR solutions, the SRI boson is - or might be its own antiparticle.

Table 9 shows $0 g \Gamma$ solutions that might associate with elementary fermions.

Regarding table 9 , we note - but do not discuss further - the notion that the number - one or two of CURR solutions associates with the number - between particles and possible antiparticles - of handednesses. Similarly, the notion of Dirac fermions might associate with cases for which two CURR solutions pertain and the notion of Majorana fermions might associate with the case for which only one CURR solution pertains.

### 2.2.4. We discuss aspects regarding the entwined

 simple elementary particles that we suggest and people have yet to find.Each one of table 8 and table 9 shows simple particles that - within the context of our work model as entwined.

After very early in the timeline that concordance

Table 7: Sets of solutions that might associate with LRI (or, long-range interaction) elementary particles. The symbol $n_{E P}$ denotes the number of elementary particles. Items that the table shows in parentheses might - depending on future data or on interpretations of vocabulary and modeling - associate with elementary particles. TBD denotes the three-word phrase to be determined.

| $\Sigma g \Gamma$ for PROP | $\Sigma g \Gamma$ for CURR | Family | Boson | $n_{E P}$ |
| :--- | :--- | :--- | :--- | :--- |
| $1 g \Gamma$ | $1 g \Gamma$ | 1 L | Photon | 1 |
| $2 g \Gamma$ | $2 g \Gamma$ | $(2 \mathrm{~L})$ | (Graviton) | $(1)$ |
| $3 g \Gamma$ | $3 g \Gamma$ | $(3 \mathrm{~L})$ | (TBD) | $(1)$ |
| $4 g \Gamma$ | $4 g \Gamma$ | $(4 \mathrm{~L})$ | (TBD) | $(1)$ |

Table 8: Solutions that might associate with SRI (or, non-LRI) elementary bosons. Each column with a label $0=\ldots$ shows a calculation that produces the $\Sigma=0$ that associates with a $0 g \Gamma$ solution. Each integer that the calculation includes is a member of $\Gamma$. No other integer is a member of $\Gamma$. The symbol $n_{\Gamma, \mathrm{PROP}}$ denotes the number of $l_{o}$ that appear in the $\Gamma$ for the PROP solution. The symbol $n_{E P}$ denotes the number of elementary particles. For each of the Higgs boson, the Z boson, and the W boson, $\left(n_{\Gamma, \mathrm{PROp}}\right)^{2}$ associates with an aspect related to the mass of the boson. (See the column - with the one-word label sum - in table 12 ) The $W$ boson is the only charged elementary boson. In table 8 , only one PROP solution does not have - as members of $\Gamma$ - all three of 1,3 , and 4 . Based on the previous two sentences, we associate $0 \mathrm{~g} 1^{〔} 2^{‘} 3$ with the W boson and we associate $0 \mathrm{~g} 1^{`} 3^{`} 4$ with the Z boson. The symbol $\dagger$ associates with the possibility that - for each one of the Higgs boson, the Z boson, and the aye boson - there might be differences between interactions with elementary bosons ( $\dagger \mathrm{b}$ ) and interactions with elementary fermions ( $\dagger \mathrm{f}$ ). For each one of the Higgs, Z, and W bosons, the PROP solution appears in table 2 For each one of the aye boson, the jay boson, and the gluons, the parenthesized item in the leftmost column of table 8 suggests that the PROP solution cascades from a CURR solution in table 8 (Regarding the notion of cascading, see discussion regarding table 4) We assume that the aye (or, 0I boson) associates with notions of an inflaton. Inflatons would be zero-mass zero-charge bosons that might have played key roles during a hypothesized inflationary epoch, early in the evolution of the universe. The table suggests that one CURR solution that associates with the Z boson equals the PROP solution that associates with the aye boson. The table suggests that one CURR solution that associates with the Higgs boson equals the PROP solution that associates with the jay boson and the eight gluons. The symbol $\ddagger$ calls attention to the notion that the two CURR solutions - that we associate with the jay boson and the eight gluons - associate with nine (or, three times three) degrees of freedom (or with nine choices). We associate one of the nine degrees of freedom (or choices) with the jay boson and the other eight degrees of freedom (or choices) with the eight gluons. We assume that the jay boson associates with notions of Pauli repulsion. Pauli repulsion associates with the notion that two fermions (whether elementary fermions or not elementary fermions) cannot occupy the same state. Pauli repulsion associates with repulsive aspects of the residual strong force. We suggest the possibility (but do not necessarily require) that some modeling might associate one of the gluon CURR solutions with aspects (of interactions) that erase - for example from a quark - one of three color charges and might associate the other of the gluon CURR solutions with aspects (of interactions) that paint - for example onto a quark - one of the three color charges.

| $0=\ldots$, re $0 \mathrm{~g} \Gamma$ for PROP | $0=\ldots$, re $0 \mathrm{~g} \Gamma$ for CURR | $n_{\Gamma, \text { PROP }}$ | Family | Bosons | $n_{E P}$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $\|+1-2-3+4\|$ | $\|+1-2-3-4+8\|(\dagger \mathrm{b})$, | 4 | 0 H | Higgs | 1 |
|  | $\|-1+2-3-4+6\|(\dagger \mathrm{f})$ |  |  |  |  |
| $\|-1-3+4\|$ | $\|-1-3-4+8\|(\dagger \mathrm{b})$, | 3 | 1 Z | Z | 1 |
|  | $\|+1-3-4+6\|(\dagger \mathrm{f})$, |  |  |  |  |
| $\|-1-2+3\|$ | $\|-1-2-3+6\|$ | 3 | $1 \mathrm{~W}_{1}$ | W | 1 |
| $\|-1-3-4+8\|(1 \mathrm{Z} \dagger \mathrm{b})$ | $\|+1-2-3-4+8\|(\dagger \mathrm{b})$, | 4 | 0 I | Aye | 1 |
|  | $\|-1+2-3-4+6\|(\dagger \mathrm{f})$ |  |  |  |  |
| $\|+1-2-3-4+8\|(0 \mathrm{H} \dagger \mathrm{b} ; \ddagger)$ | $\|+1-2+3-4-6+8\|$, | 5 | 1 J | Jay | 1 |
|  | $\|-1-2-3+4-6+8\|$ |  |  |  |  |
| $\|+1-2-3-4+8\|(0 \mathrm{H} \dagger \mathrm{b} ; \ddagger)$ | $\|+1-2+3-4-6+8\|$, | 5 | 1 G | Gluons | 8 |
|  | $\|-1-2-3+4-6+8\|$ |  |  |  |  |

Table 9: Solutions that might associate with elementary fermions (or, ELF elementary particles). We assume, regarding PROP 0 g solutions, that $6 \in \Gamma$ associates with elementary fermions. (The might-be three degrees of freedom that might associate with $l_{6}$ might associate with three choices regarding flavours for elementary fermions.) Paralleling notions pertaining to SRI elementary bosons, if, and only if, one of 1,3 , and 4 is not a member of a PROP $\Gamma$, an elementary particle that associates with table 9 has nonzero charge. The symbol $n_{E P}$ denotes the number of elementary particles. The leftmost column in table 9 alludes to relevant PROP solutions and alludes to - in parentheses - at least one CURR solution from which a PROP solution cascades. (Regarding the notion of cascading, see discussion related to table 4 ) We discuss solutions that associate with quarks. $0.5 \mathrm{Q}_{1 / 3}$ particles and $0.5 \mathrm{Q}_{2 / 3}$ particles are the only known particles for which $0<Q<1$. Here, $Q$ denotes the magnitude of the charge, in units of $\left|q_{e}\right| . q_{e}$ denotes the charge of the electron. The notion of PROP solution might associate with each of the following two solutions: $0=|-1+2-3-6+8|$ and $0=|-1+2-3-4+6|$. The notion of $Q=1$ associates with the first of the two PROP solutions. The notion of $Q=0$ associates with the second of the two PROP solutions. The notions of $Q=1 / 3$ and $Q=2 / 3$ might associate with states that mix - and, regarding charge, lie between $-Q=0$ and $Q=1$. (See aspects of table 13 and table 14 ) The symbol $\ddagger$ associates with the notion that the PROP solution that associates with the 0.5 M particles equals one of the PROP solutions that associates with quarks. This essay does not further address the notion that this would-be reuse of a PROP solution might associate with the notion that nature might not include heavy neutrinos.

| $0=\ldots$, re $0 g \Gamma$ for PROP | $0=\ldots$, re $0 g \Gamma$ for CURR | Families | Fermions | $n_{E P}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\|+1-3-4+6\|(1 \mathrm{Z} \dagger \mathrm{f})$ | $\|-1+3-4-6+8\|$ | 0.5 N | Neutrinos | 3 |
| $\|-1-2-3+6\|\left(1 \mathrm{~W}_{1}\right)$ | $\|-1+2-3-4+6\|$, | $0.5 \mathrm{C}_{1}$ | Charged leptons | 3 |
|  | $\|-1+2-3-6+8\|$ |  |  | 3 |
| $\|-1+3-4-6+8\|(0.5 \mathrm{~N})$ | $\|-1-2-3+4-6+8\|$, | 0.5 R | Arcs | 3 |
|  | $\|+1-2+3-4-6+8\|$ |  |  | 6 |
| $\|-1+2-3-6+8\|\left(0.5 \mathrm{C}_{1}\right)$, | $\|-1-2-3+4-6+8\|$, | $0.5 \mathrm{Q}_{y / 3}$, | Quarks | 6 |
| $\|-1+2-3-4+6\|(0 \mathrm{H} \dagger \mathrm{f}$, | $\|+1-2+3-4-6+8\|$ | $y=1$ or 2 |  |  |
| $\left.0 \mathrm{I} \dagger \mathrm{f}, 0.5 \mathrm{C}_{1}\right)$ |  |  | Heavy neutrinos | 3 |
| $\|-1+2-3-4+6\|(0 \mathrm{H} \dagger \mathrm{f}$, | $\|-1-2-3+4-6+8\|$, | 0.5 M |  |  |
| $\left.0 \mathrm{I} \dagger \mathrm{f}, 0.5 \mathrm{C}_{1} ; \ddagger\right)$ | $\|+1-2+3-4-6+8\|$ |  |  |  |

cosmology features, 0.5 R particles might exist only in hadron-like particles that include 1G (or, gluon) particles and that are somewhat analogs to known hadrons. In the very early universe, 0.5 R and 1 G particles might model as being components of seas.

After very early in the timeline that concordance cosmology features, OI particles might contribute effects that are negligible compared to effects that other bosons contribute. In the very early universe, OI particles might fulfill the role that concordance cosmology posits for the hypothesized inflaton elementary particle.

### 2.3. Isomers and dark matter

This unit suggests that most dark matter has bases in five isomers of the elementary particles that are not LRI elementary particles and that ordinary matter has bases in one (other) isomer of most elementary particles that are not LRI elementary particles.

### 2.3.1. We discuss the notion that, if nature includes only one isomer of each elementary particle, modeling might not suffice to explain known data about dark matter.

Discussion above points to two types of elementary particles that would measure as dark matter or that would provide bases for dark matter. 0.5 M fermions associate with the notion of free and would measure as dark matter. 0.5 R fermions associate with the notion of entwined. Hadron-like particles containing gluons and 0.5 R fermions would contain
no charged particles and would measure as dark matter.

We use the term DMAI to denote stuff that has bases in 0.5 M elementary fermions or in 0.5 R elementary fermions. DM abbreviates the two-word term dark matter. AI abbreviates the two-word term all isomers. (Here, we allude to a notion of multiple isomers of some elementary particles. For the moment we assume that nature includes just one isomer.)

We use notation of the form DM:OM to denote an inferred ratio of DM effects to OM effects. OM abbreviates the two-word term ordinary matter.

Measurements suggest seemingly significant DM:OM ratios. (For information about the ratios and for relevant references, see table 22 and discussion related to table 22, Ratios of approximately $5^{+}: 1$ pertain regarding densities of the universe, galaxy clusters, and many galaxies. Seemingly significant ratios of $1: 0^{+}, 0^{+}: 1$, and $\sim 4: 1$ pertain regarding some galaxies. A ratio of $1: 1$ might pertain regarding some depletion of cosmic microwave background radiation (or, CMB). We know of no other such seemingly significant DM:OM ratios.

We suggest that, if DMAI is the only type of dark matter, DMAI might not suffice to explain various seemingly significant ratios of dark matter to ordinary matter. We suggest that the notion of DMAI might not suffice to explain dark matter.

2．3．2．We discuss the notion that nature includes six isomers of each elementary particle that is not an LRI boson．
We suggest that nature includes six isomers of the SRI and ELF elementary particles－or，six iso－ mers of the set of elementary particles that asso－ ciates with all non－LRI elementary bosons and all elementary fermions．（See table 8 and table 9 ）One isomer associates with ordinary matter plus one iso－ mer of DMAI．That one isomer of DMAI measures as dark matter．Each one of the other five isomers of the set of non－LRI elementary particles measures as dark matter．Regarding densities of the universe， the five isomers of non－DMAI that measure as dark matter associate with the 5 in the DM：OM ratio of $5^{+}: 1$ ．The six isomers of DMAI associate with the ${ }^{+}$in the DM：OM ratio of $5^{+}: 1$ ．

We use a two－word phrase isomer number to de－ note one isomer．Here，number can be any one of zero，one，．．．，and five．We associate the two－word term isomer zero with the isomer that includes ordi－ nary matter．We use the two－word phrase alt isomer to denote any one of the five isomers that does not associate with ordinary matter．

2．3．3．We discuss modeling－regarding simple ele－ mentary particles－that might associate with the notion of six isomers of simple elemen－ tary particles．
Table 8 and table 9 point to all simple elemen－ tary particles that this essay features．For each PROP solution，each one of $1 \in \Gamma$ and $3 \in \Gamma$ per－ tains．We suggest that－relative to one of those two membership（in $\Gamma$ ）associations，the other member－ ship（in $\Gamma$ ）association associates with three choices． （Here，the notion of three choices associates－for other circumstances－with notions that an $l_{o}$ asso－ ciates with three degrees of freedom．See discussion －regarding table 9 －regarding $6 \in \Gamma$ and three flavours．）Each PROP solution in table 8 and table 9 associates with two expressions．We assume that one expression associates with left－handedness and one expression associates with right－handedness． We point to the possibility that the combination of three choices and two handednesses associates with six isomers．We anticipate using this notion regard－ ing the six isomers．（See discussion related to table 18．）

## 2．3．4．We discuss modeling－for long－range inter－ actions－that associates with the notion of six isomers of simple elementary particles．

All six isomers produce and interact with a com－ mon notion of gravity．We suggest that one instance of the PROP solution 2 g 2 associates with interac－ tions between all six isomers．We say that one in－ stance of the PROP solution 2 g 2 has a reach of
six，as in six isomers．We suggest that each iso－ mer associates with its own instance of the PROP solution 1 g 1 and its own instance of the PROP so－ lution 1g1＇2．（Reference［7］suggests the notions of dark matter charges and dark matter photons．） We say that each instance of the PROP solution 1 g 1 has a reach of one，as in one isomer．Each in－ stance of the PROP solution $1 \mathrm{~g} 1^{〔} 2$ has a reach of one．Each isomer－including the ordinary matter isomer－scarcely interacts with any other isomer via electromagnetism．

We address the topic of reach for each PROP solution $\Sigma \mathrm{g} \Gamma$ to which table 2 alludes．

Based on the reach of the PROP solution 1 g 1 and the reach of the PROP solution $1 \mathrm{~g} 1{ }^{\prime} 2$ ，we sug－ gest that $n_{0}=0$ associates with a reach of one． Based on the reach of the PROP solution 2g2，we suggest that $n_{0}=1$ associates with a reach of six．

We assume that，for $n_{0} \geq 1$ ，equation（3）com－ putes the reach．Here，$\rho_{I}$ denotes the reach．Here， $\operatorname{gen}\left(S U\left(l_{g r}\right)\right)$ denotes the number of generators of the group $S U\left(l_{g r}\right)$ ．For $l_{g r} \geq 2$ ，gen $\left(S U\left(l_{g r}\right)\right)$ equals $\left(l_{g r}\right)^{2}-1$ ．

$$
\begin{equation*}
\rho_{I}=\operatorname{gen}(S U(7)) / \operatorname{gen}\left(S U\left(2 n_{0}+1\right)\right) \tag{3}
\end{equation*}
$$

The reach that associates with $n_{0}=2$ is two． The reach that associates with $n_{0}=3$ is one．

The number of instances of a PROP $\Sigma \mathrm{g} \Gamma$ com－ ponent of an LRI elementary particle is six divided by the reach that associates with the PROP $\Sigma g \Gamma$ solution．

We assume that the reach of a CURR counter－ part solution to a PROP $\Sigma g \Gamma$ solution is the same as the reach of the PROP $\Sigma g \Gamma$ solution．

We address the reach of the $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 6^{‘} 8 \mathrm{x}$ PROP solutions．For $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 6^{‘} 8$ ，each of 1,2 ，and 3 ap－ pears in $\Gamma$ and 4 does not appear in $\Gamma$ ．We assume that $n_{0}=1$ ．The reach for each $2 g 1^{\prime} 2^{\prime} 3^{‘} 6^{\prime} 8 \mathrm{x}$ is six． We assume that stress－energy is the relevant PROP－ associated property．Stress－energy－and－velocity 4－ somes can associate with stuff－such as a galaxy cluster－that associates with all of the six iso－ mers．Stress－energy－and－velocity 4 －somes can as－ sociate with stuff－such as a star－that associates with less than all of the six isomers．

We address the reach of the $2 \mathrm{~g} 1^{〔} 2^{〔} 3^{〔} 4^{〔} 6^{〔} 8 \mathrm{x}$ PROP solutions．For $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 4^{‘} 6^{‘} 8$ ，each of 1,2 ， 3 ，and 4 appears in $\Gamma$ ．We assume that $n_{0}=0$ ． The reach for each $2 \mathrm{~g}^{\prime} 2^{\prime} 3^{\prime} 4^{\prime} 6^{\prime} 8 \mathrm{x}$ solution is one． We might assume that abilities to change stress－ energy associates with relevant PROP－associated properties．Stress－energy changes might－if model－ ing based on them has physics relevance－associate with single－isomer stuff－such as an atomic nucleus．

Table 10 shows the reach $\left(\rho_{I}\right)$ for－and other information about－each one of some solutions that table 4 and table 5 list．Discussion that relates to

Table 10: Reaches and other information regarding some solutions that associate with electromagnetism, gravity, 3L, and 4L. $\rho_{I}$ denotes reach. $\Sigma \in \Gamma$ associates with the symbol g'. $\Sigma \notin \Gamma$ associates with the symbol g". TBD denotes to be determined. NYN denotes not yet named. Table 22 suggests that the PROP solution $1 \mathrm{~g} 1 ‘ 4 ‘ 6 ‘ 8$ associates with (at least) hyperfine states for atoms.

| $S$ | $\Sigma$ | PROP solution | $\rho_{I}$ | Solution type | PROP RDF | Properties or other associations |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | 1g1 | 1 | $\Sigma \in \Gamma$ | $\Xi^{-2}$ | Charge |
| 1 | 1 | 1g1'2 | 1 | $\Sigma \in \Gamma$ | $\Xi^{-3}$ | Magnetic moment |
| 1 | 1 | $1 \mathrm{~g} 1^{\prime}{ }^{\prime} 4$ | 6 | $\Sigma \in \Gamma$ | $\Xi^{-4}$ | Rotating axis of magnetic moment |
| 1 | 1 | $1 \mathrm{~g} 1^{6}{ }^{6} 6$ | 2 | $\Sigma \in \Gamma$ | $\Xi^{-4}$ | Charge, rotation, and velocity |
| 1 | 1 | 1g1'4 $6^{\prime} 8$ | 2 | $\Sigma \in \Gamma$ | $\Xi^{-5}$ | Charge, rotation, and angular velocity |
| 1 | 1 | $1 \mathrm{~g} 1^{\prime} 4^{4} 6^{6} 8 \mathrm{x}$ | 6 | $\Sigma \in \Gamma$ | $\Xi^{-6}$ | Changes (re systems that include charged particles) re internal angular momenta |
| 2 | 2 | 2g2 | 6 | $\Sigma \in \Gamma$ | $\Xi^{-2}$ | Mass |
| 2 | 2 | 2g2 4 | 2 | $\Sigma \in \Gamma$ | $\Xi^{-3}$ | Rotating (spherically symmetric aspects of) mass |
| 2 | 2 | 2g1'2'3 | 1 | $\Sigma \in \Gamma$ | $\Xi^{-4}$ | Moments of inertia (stress-energy) |
| 2 | 2 | $2 \mathrm{~g} 1^{\prime} 3^{\prime} 3^{\prime} 4 \mathrm{v}$ | 1 | $\Sigma \in \Gamma$ | $\Xi^{-5}$ | Rotating axis of moment of inertia |
| 2 | 2 | $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\text {4 }} 4 \mathrm{w}$ | 1 | $\Sigma \in \Gamma$ | $\Xi^{-5}$ | Rotating axis of moment of inertia |
| 2 | 2 | $2 \mathrm{~g} 1^{\prime} \mathrm{C}^{\prime} 6^{6} 88 \mathrm{x}$ | 6 | $\Sigma \in \Gamma$ | $\Xi^{-6}$ | Stress-energy |
| 2 | 2 | $2 \mathrm{~g} 1 \cdot 2^{\prime} 3^{\prime} 4^{6} 68 \mathrm{x}$ | 1 | $\Sigma \in \Gamma$ | $\Xi^{-7}$ | Changes re stress-energy |
| 3 | 3 | 3 g 3 | 2 | $\Sigma \in \Gamma$ | $\Xi^{-2}$ | TBD (a function of elementary fermion flavour), NYN |
| 3 | 3 | $3 \mathrm{~g} 1 \times 2$ | 1 | $\Sigma \notin \Gamma$ | $\Xi^{-3}$ | Anomalous magnetic moment |
| 3 | 3 | $3 \mathrm{~g} 1^{\prime}{ }^{\prime} 4$ | 6 | $\Sigma \notin \Gamma$ | $\Xi^{-4}$ | Anomalous magnetic moment |
| 3 | 3 | $3 \mathrm{~g} 1^{\prime}{ }^{〔} 6$ | 2 | $\Sigma \notin \Gamma$ | $\Xi^{-4}$ | Anomalous magnetic moment |
| 4 | 4 | 4 g 4 | 1 | $\Sigma \in \Gamma$ | $\Xi^{-2}$ | Elementary particle angular momentum (scalar quantity) |
| 4 | 4 | $4 \mathrm{~g} 1^{\prime} 2^{〔} 3$ | 1 | $\Sigma \notin \Gamma$ | $\Xi^{-4}$ | Anomalous gravitational property |

table 4 and to table 7 suggests that some items that table 4 lists are not necessarily relevant to modeling that pertains to relevant physics. Table 10 does not include some not necessarily relevant items.

Regarding the notion of a reach, $\rho_{I}$, of two, there are three instances of the PROP solution. We number the isomers so that one instance of the PROP $2 \mathrm{~g} 2^{\prime} 4$ solution associates with interactions between isomer zero and isomer three. One instance of the PROP 2g2‘4 solution associates with interactions between isomer one and isomer four. One instance of the PROP $2 \mathrm{~g} 2^{‘} 4$ solution associates with interactions between isomer two and isomer five.

We use notation of the form $\Sigma\left(\rho_{I}\right) g \Gamma$ to denote a PROP $\Sigma \mathrm{g} \Gamma$ solution and the reach $\rho_{I}$ that associates with one modeling use that features an instance of the solution. For example, 2(2)g2‘4 pertains regarding the PROP $2 \mathrm{~g} 2{ }^{\prime} 4$ solution. We extend use of such notation to non-LRI elementary particles. For non-LRI elementary particles, the reach is one and notation of the form $S(1) \Phi$ pertains.

We assume that - for each PROP $\Sigma(2) \mathrm{g} \Gamma$ solution - one instance of the solution associates with interactions between isomer zero and isomer three. One instance of the solution associates with interactions between isomer one and isomer four. One
instance of the solution associates with interactions between isomer two and isomer five.

### 2.3.5. We start to discuss the extent to which the

 properties of any one isomer's elementary particles differ from the properties of elementary particles that associate with other isomers.If the stuff that associates with each of the five all-dark-matter isomers evolved similarly to ordinary matter, our suggestions regarding dark matter might not adequately comport with observations regarding the Bullet Cluster collision of two galaxy clusters. We anticipate that the isomers of ELF elementary particles differ in ways such that our suggestions regarding dark matter do not necessarily disagree with observations pertaining to the Bullet Cluster. (See discussion related to table 18.)

### 2.3.6. We discuss notions regarding excitations and de-excitations of LRI fields (or, of LRI elementary particles).

An excitation of a $\Sigma \mathrm{g}$ LRI field associates with the value of $\Sigma$ and with a set of isomers that associate with an object A that contributes to - or excites - the field. The notion of the two-word term active attribute pertains regarding object A. (Here, the word attribute includes PROP properties and

CURR aspects that associate with translational motions of the PROP properties.) Regarding the notion that an excitation of a field encodes information that associates with the excitation, we assume that an excitation encodes information that associates with the relevant set of isomers.

A set of isomers associates with an object B that takes from - or de-excites - the field. The notion of the two-word term passive attribute pertains regarding object B.

Ten types of de-excitations exist. One type consists of de-excitations that associate with reach-six solutions. Three types consist of de-excitations that associate with reach-two solutions. One of the three types associates with isomer zero and isomer three. Another one of the three types associates with isomer one and isomer four. The other one of the three types associates with isomer two and isomer five. Six types consist of de-excitations that associate with reach-one solutions. Each one of the six types associates with exactly one isomer.

We consider some cases that associate with electromagnetism.

We consider an object B that consists entirely of ordinary matter.

- The attributes charge, translational motion of charge, magnetic moment, and translational motion of magnetic moment associate with reaches of one. (For information about reaches - for this case - and for other cases, see table 10.) Regarding these attributes, object $B$ interacts only with the aspects - of the electromagnetic field - that associate with active contributions from isomer zero aspects of objects A. For example, an ordinary matter object $B$ engages - via the passive attributes of charge, translational motion of charge, magnetic moment, and translational motion of magnetic moment - only in interactions that would associate with electromagnetism that isomer zero generates.
- The attributes that associate with hyperfine state (for example, of a hydrogen atom) and translational motion of a hyperfine state associate with reaches of two. An ordinary matter object B that exhibits hyperfine states engages - via those states and translational motion of those states - in interactions that would associate with electromagnetism that isomer zero generates and in interactions that would associate with electromagnetism that isomer three generates.

We consider some cases that associate with gravitation.

We consider an object B that consists entirely of ordinary matter.

- The attributes that associate with mass and translational motion of mass associate with reaches of six. (For information about reaches - for this case - and for other cases, see table 10.) Regarding these attributes, object B interacts with the aspects - of the gravitational field - that associate with active contributions from all objects $A$, independent of the isomeric composition of the objects A. Regarding mass, an RDF of $\Xi^{-2}$ pertains. Regarding translational motion of mass, an RDF of $\Xi^{-3}$ pertains.
- The attributes that associate with non-spherically-symmetric mass and motion of non-spherically-symmetric mass associate with reaches of one. An ordinary matter object $B$ that exhibits aspects of non-spherically-symmetric mass and motion of non-spherically-symmetric mass engages via those attributes - only in interactions that would associate with gravitation that isomer zero generates. Regarding non-spherically-symmetric mass, an RDF of $\Xi^{-4}$ pertains. Regarding translational motion of non-spherically-symmetric mass, an RDF of $\Xi^{-5}$ pertains. A lack of effects that would associate with active contributions - to gravity by object A - of five dark-matter isomers and with an RDF of $\Xi^{-4}$ or an RDF of $\Xi^{-5}$ is - for some circumstances - not necessarily significant compared to reach-six aspects that associate with an RDF of $\Xi^{-2}$ or an RDF of $\Xi^{-3}$.


## 3. Results

This unit discusses explanations for known data and discusses suggestions regarding possible data that people have not yet measured. The discussion includes explanations and suggestions regarding elementary particles, dark matter, galaxies, dark energy, and the cosmos.

### 3.1. Elementary particles

This unit lists elementary particles that associate with our modeling and discusses relationships between properties of elementary particles.

### 3.1.1. We list all elementary particles of which people know or that we suggest.

Table 11 consolidates and summarizes information about all elementary particles of which people know or that this essay suggests. (See table 7 table 8, and table 9.)

Reference [8] notes that modeling based on QFT (or, quantum field theory) suggests that massless elementary particles cannot have spins that exceed

Table 11: Elementary particles. The symbol $Q$ associates with magnitude of charge. The columns labeled $Q>0$ and $Q=0$ have entries in the form of a name of one particle or a name of a set of more than one particle, followed (in parentheses) by a number of particles, followed by a symbol for the family of particles. NYN denotes not yet named. NYD denotes not yet detected. One might assert that people know of some NYD particles, at least indirectly. The word free associates with modeling that features PROP solutions for which $8 \notin \Gamma$. The word entwined associates with modeling that features PROP solutions for which $8 \in \Gamma$. For 1L, some modeling might associate with entwined. (See table 4 and discussion - regarding cascades - related to table 4) For example, notions of entwined might pertain regarding electromagnetism within an atom or regarding light in a laser cavity. We associate the word mixed with $\Sigma \mathrm{L}$ for which some relevant components associate with PROP solutions for which $8 \notin \Gamma$ and some relevant components associate with PROP solutions for which $8 \in \Gamma$. For $2 L$, some modeling might associate with entwined. For example, notions of entwined might pertain regarding gravitation within a black hole. For each of 3L and 4L, notions of entwined might not pertain. (See discussion - regarding cascades related to table 4.)

| $S$ | $m$ | $Q>0$ | $Q=0$ | Status | $\Sigma$ | Free / Entwined |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | $>0$ | - | Higgs boson (1), 0H | Known | 0 | Free |
| $1 / 2$ | $>0$ | Charged leptons (3), $0.5 \mathrm{C}_{1}$ | Neutrinos (3),0.5N | Known | 0 | Free |
| $1 / 2$ | $>0$ | - | Heavy neutrinos (3), 0.5M | NYD | 0 | Free |
| 1 | $>0$ | W boson (1), $1 \mathrm{~W}_{1}$ | Z boson (1), 1Z | Known | 0 | Free |
| 1 | $=0$ | - | Photon (1), 1L | Known | 1 | Mixed |
| 2 | $=0$ | - | Graviton (1), 2L | NYD | 2 | Mixed |
| 3 | $=0$ | - | NYN (1), 3L | NYD | 3 | Free |
| 4 | $=0$ | - | NYN (1), 4L | NYD | 4 | Free |
| 0 | $=0$ | - | Aye boson (1), 0I | NYD | 0 | Entwined |
| $1 / 2$ | $>0$ | Quarks $(3), 0.5 \mathrm{Q}_{1 / 3}$ | - | Known | 0 | Entwined |
| $1 / 2$ | $>0$ | Quarks $(3), 0.5 \mathrm{Q}_{2 / 3}$ | - | Known | 0 | Entwined |
| $1 / 2$ | $>0$ | - | Arcs (3), 0.5R | NYD | 0 | Entwined |
| 1 | $=0$ | - | Jay boson $(1), 1 \mathrm{~J}$ | NYD | 0 | Entwined |
| 1 | $=0$ | - | Gluons (8), 1G | Known | 0 | Entwined |

two. In our work, 3L might associate with nonzero anomalous magnetic moments for at least charged leptons. (See table 5 and discussion related to equation (4).) Modeling based on QFT suggests - without assuming elementary particles with spins of more than one - values for some anomalous magnetic moments. In our work, 4L might associate with notions - for at least neutrinos - of anomalous gravitational properties and mass mixing. (See table 5 and table 15 .) We suggest the notion that a (possible but for now hypothetical) QFT that includes gravity might successfully estimate aspects that people associate with neutrino mass mixing. Our work is not necessarily incompatible with notions that nature does not include zero-mass elementary bosons that have spins that exceed two.

### 3.1.2. We explore relationships among properties of elementary bosons.

Table 12 discusses relationships between properties of elementary bosons. (Reference [9] provides data regarding the masses of the Higgs, Z, and W bosons.)

Table 12 points to possibly deeper (than people might otherwise suggest) relationships between the physics properties of spin, mass, and charge. (Also, regarding the non-zero-mass elementary bosons, a notion that non-zero spin might associate with - in effect - reduction in mass seems not to be incompatible with discussion related to table 6.)

We mention the following notions regarding
anomalous properties. (See discussion related to table5.) The notion - for the Higgs, Z, and W bosons - that $l_{m s}$ is not zero might associate with a notion of anomalous property. To the extents that the ra$\operatorname{tios}\left(m_{\mathrm{Higgs}}\right)^{2}:\left(m_{\mathrm{Z}}\right)^{2}:\left(m_{\mathrm{W}}\right)^{2}$ do not measure as exactly $17: 9: 7$, anomalous properties might associate with variations from exactness. (Here, $m$ denotes mass.)

### 3.1.3. We explore relationships among properties of charged elementary fermions.

We consider hypothetical elementary fermions for which $Q=1$. For some value of mass, the gravitational attraction between two identical such hypothetical elementary fermions would equal the electrostatic repulsion between the two fermions. Our work shows that a mass - for which we use the expression $m(18,3)$ - seems to have meaning beyond the notion that - for the mass $m(18,3)$ - gravitational attraction between two $Q=1$ identical elementary fermions would be three-quarters of the electrostatic repulsion between the two identical elementary fermions. (See table 14 and table 17 .)

Table 13 discusses relationships between properties of known charged elementary fermions. (Reference [9] provides the data that underlies table 13.)

Table 14 shows equations that underlie aspects of table 13 (Reference 9 provides the data that underlies table 14 .)

Table 13 and table 14 might point to possibly deeper (than people might otherwise suggest) rela-

Table 12: Relationships between properties of elementary bosons. $Q$ denotes the magnitude of charge, in units of $\left|q_{e}\right| . m$ denotes mass, in units of $m_{\text {Higgs }} / 17^{1 / 2}$ or in units of $m_{\mathrm{Z}} / 9^{1 / 2}$. $S$ denotes spin, as in the expression $S(S+1) \hbar^{2}$. $l_{m s}$ equals -1 for $m>0$ and equals 0 for $m=0$. The sum is the sum of the numbers in the preceding four columns. Each sum is the square of an integer. For each nonzero-mass particle, the integer equals $n_{\Gamma, \text { Prop }}$. (See table 8 ) There are no-nonzero mass elementary bosons for which the integer equals one or two. (For a $\Gamma$ that includes just one value of $l_{o}$ or that includes just two values of $l_{o}, \Sigma \neq 0$ pertains.) NYN denotes the three-word phrase not yet named. Of the non-zero masses to which table 12 alludes, the most accurately known mass is that of the Z boson. Using the mass of the Z boson and numbers in table 12 one can calculate a nominal mass for the Higgs boson and a nominal mass for the W boson. The calculated mass for the Higgs boson differs from the experimentally determined mass by less than two (experimental) standard deviations. The calculated mass for the W boson differs from the experimentally determined mass by less than four (experimental) standard deviations. To the extent that one uses the notion that ruling out an equality requires a difference of at least five standard deviations, experimental results do not seem to rule out relationships that table 12 states.

| Bosons | Family | $Q(Q+1)$ | $m^{2}$ | $S^{2}$ | $l_{m s}$ | Sum |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| Higgs | 0 H | 0 | 17 | 0 | -1 | 16 |
| Aye | 0 I | 0 | 0 | 0 | 0 | 0 |
| Z | 1 Z | 0 | 9 | 1 | -1 | 9 |
| W | $1 \mathrm{~W}_{1}$ | 2 | 7 | 1 | -1 | 9 |
| Jay | 1 J | 0 | 0 | 1 | 0 | 1 |
| Gluons | 1G | 0 | 0 | 1 | 0 | 1 |
| Photon | 1L | 0 | 0 | 1 | 0 | 1 |
| Graviton | 2L | 0 | 0 | 4 | 0 | 4 |
| NYN | 3L | 0 | 0 | 9 | 0 | 9 |
| NYN | 4L | 0 | 0 | 16 | 0 | 16 |

Table 13: Values of $\log { }_{10}\left(m_{\text {particle }} / m_{e}\right)$ for known charged elementary fermions. Regarding "flavour," this table generalizes, based on terminology that associates with charged leptons and with neutrinos. For example, people use the term electronneutrino. The symbol $l_{f}$ numbers the three flavours. The " $l_{f}\left(0.5 \mathrm{C}_{1}\right)$ " terms pertain for fermions in the $0.5 \mathrm{C}_{1}$ family. The symbol $0.5 \mathrm{Q}_{>0}$ denotes the pair $0.5 \mathrm{Q}_{1 / 3}$ and $0.5 \mathrm{Q}_{2 / 3}$. The " $l_{f}\left(0.5 \mathrm{Q}_{>0}\right)$ " terms pertain for quarks (or, elementary particles in the two families $0.5 \mathrm{Q}_{2 / 3}$ and $\left.0.5 \mathrm{Q}_{1 / 3}\right) . l_{m}$ is an integer parameter. The domain $-6 \leq l_{m} \leq 18$ might have relevance regarding modeling. $Q$ denotes the magnitude of charge, in units of $\left|q_{e}\right|$. The family $0.5 \mathrm{C}_{1}$ associates with $Q=1$. The family $0.5 \mathrm{Q}_{2 / 3}$ associates with $Q=2 / 3$. The family $0.5 \mathrm{Q}_{1 / 3}$ associates with $Q=1 / 3$. Regarding the rightmost four columns, items show $\log { }_{10}\left(m_{\text {particle }} / m_{e}\right)$ and - for particles that nature includes - the name of an elementary fermion. For each $\dagger$ case, no particle pertains. Each number in the column with label $Q=1 / 2$ equals the average of the number in the $Q=2 / 3$ column and the number in the $Q=1 / 3$ column. The notion of geometric mean pertains regarding the mass of the $Q=2 / 3$ particle and the mass of the $Q=1 / 3$ particle. Regarding each $\dagger$ case, a formula for $m\left(l_{m}, l_{q}\right)$ calculates the number. Regarding the formula, the domain $0 \leq l_{q} \leq 3$ pertains. Regarding table $13 l_{q}=3 Q$ pertains. Table 14 shows the formula.

| $l_{f}\left(0.5 \mathrm{C}_{1}\right)$ | $l_{f}\left(0.5 \mathrm{Q}_{>0}\right)$ | $l_{m}$ | $Q=1$ | $Q=2 / 3$ | $Q=1 / 2$ | $Q=1 / 3$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 1 (Electron) | 1 (Up, Down) | 0 | 0.00 Electron | 0.66 Up | $0.80 \dagger$ | 0.94 Down |
| - | 2 (Charm, Strange) | 1 | $1.23 \dagger$ | 3.36 Charm | $2.83 \dagger$ | 2.29 Strange |
| $2(\mathrm{Mu})$ | 3 (Top, Bottom) | 2 | 2.32 Muon | 5.52 Top | $4.72 \dagger$ | 3.92 Bottom |
| 3 (Tau) | - | 3 | 3.54 Tau | - | - | - |

Table 14: Equations that underlie aspects of table 13 This table shows equations that may pertain regarding all known charged elementary fermions, the known 0.5 N neutrinos, and the suggested 0.5 R arcs. (Regarding 0.5 N neutrinos, see table 15 Regarding 0.5R arcs, see table 16 )

tionships between the physics properties of mass， charge，and flavour．

## 3．1．4．We show modeling that might estimate the anomalous magnetic moment for the tau el－ ementary particle．

We explore modeling regarding anomalous mag－ netic moments for $0.5 \mathrm{C}_{1}$ elementary particles（or， charged leptons）．

Table 5 associates two CURR solutions with the $3 \mathrm{~g} 1 ‘ 2$ PROP solution．The $3 \mathrm{~g} 1{ }^{\prime} 2^{‘} 6 \mathrm{CURR}$ solution includes a 6 in $\Gamma$ ．We assume that the strength of $3 \mathrm{~g} 1^{〔} 2^{〔} 6$ can vary based on elementary fermion flavour，but not based on charge．The $3 \mathrm{~g} 1^{‘} 2^{〔} 4$ CURR solution does not include a 6 in $\Gamma$ ．We as－ sume that the strength of $3 \mathrm{~g} 1^{\prime} 2^{\prime} 4$ can vary based on charge，but not based on elementary fermion flavour．

We explore the notion that one can approxi－ mate $a_{c l}$ ，the anomalous magnetic moment for the $c l$ charged lepton，via equation（4）．

$$
\begin{equation*}
a_{c l} \approx a_{4}+a_{6} t_{c l} \tag{4}
\end{equation*}
$$

Here，$a_{4}$ might vary only with charge and would be a constant with respect to a choice between $c l=e($ for the electron $), c l=\mu$（for the muon）， and $\mathrm{cl}=\tau$（for the tau）．Here，$a_{6}$ might vary only with fermion flavour．We assume that $t_{\mathrm{cl}}$ is $\left(\log \left(m_{\mathrm{cl}} / m_{e}\right)\right)^{2}$ ．（Perhaps，compare with table 13 and with aspects－that comport with squares of properties－of table 14．The notion of squares of properties might associate with notions of self－ interactions．）Based on data that reference［9］pro－ vides regarding the electron and the muon，we cal－ culate $a_{4}$ and $a_{6}$ ．Then，we calculate a value，$a_{\tau, \mathrm{PM}}$ ， for $a_{\tau}$ ．Here，PM denotes the two－word term pro－ posed modeling．PM associates with our work．Ref－ erence［10］provides，based on Standard Model mod－ eling techniques，a first－order result－which we call $a_{\tau, \mathrm{SM}}$－for $a_{\tau}$ ．Here，SM denotes the two－word term Standard Model．The value of $a_{\tau, \text { PM }}$ results in a value of $\left(a_{\tau, \mathrm{PM}}-a_{\tau, \mathrm{SM}}\right) / a_{\tau, \mathrm{SM}}$ of approximately -0.00228 ．Each one of $a_{\tau, \mathrm{PM}}$ and $a_{\tau, \mathrm{SM}}$ comports with experimental data that reference［9］provides．

Regarding anomalous magnetic moments，this essay does not explore quantifying aspects that as－ sociate with higher－order Standard Model terms or aspects that might associate with the PROP solu－ tions $3 \mathrm{~g} 1^{‘} 2^{〔} 4$ and $3 \mathrm{~g} 1^{‘} 2^{〔} 6$ ．（Regarding the PROP solutions $3 \mathrm{~g} 1^{\prime} 2^{‘} 4$ and $3 \mathrm{~g} 1^{‘} 2^{‘} 6$ ，see table 5．）

## 3．1．5．We discuss the masses of neutrinos．

Table 15 suggests rest energies that may pertain regarding the 0.5 N neutrinos．This table extends aspects of table 13 and table 14 ．（Reference［9］pro－ vides data that underlies aspects of table 13，table 14 and table 15 ．Reference［11］discusses the notion of neutrino mass mixing．）

3．1．6．We discuss possible masses for the zero－ charge quark－like elementary fermions that we suggest．
Table 16 suggests rest energies that might per－ tain regarding the suggested 0.5 R arcs．This table extends aspects of table 13 and table 14 （Reference ［9］provides data that underlies aspects of table 13 ， table 14 and table 16．）

We explore two alternatives regarding values of $d^{\prime}(0), d^{\prime}(1)$ ，and $d^{\prime}(2)$ ．（See table 14．）Changing those numbers would impact the calculated masses for quarks and the calculated suggested masses for arcs．（Changing those numbers would not impact the calculated masses for charged leptons．）Regard－ ing each of the two alternatives，if one excludes one of three methods for estimating the mass of the top quark，the calculated mass for each of the six quarks is within five standard deviations of the experimen－ tal mass．（Reference［9］discusses the three meth－ ods．）For the third method for estimating the mass of the top quark，the value that we calculate for the mass of the top quark would be less than eleven standard deviations below the mass people have cal－ culated．

One alternative has bases in the notions of $d^{\prime}(-1)=0^{2} / 2^{2}, d^{\prime}(0)=1^{2} / 2^{2}, d^{\prime}(1)=-2^{2} / 2^{2}$ ， and $d^{\prime}(2)=-(2 \times 3) / 2^{2}$ ．For this alternative， the three arc rest energies would，respectively，be $\approx 8.14 \mathrm{MeV}, m(1,3) c^{2}$ ，and $m(2,3) c^{2}$ ．

The other alternative has bases in the notions of $d^{\prime}(0) \approx 0.264825, d^{\prime}(1)=-2^{2} / 2^{2}$ ，and $d^{\prime}(2)=$ $-(2 \times 3) / 2^{2}$ ．For this alternative，the three arc rest energies would，respectively，equal $m(1,3) c^{2}$ ， $m(1,3) c^{2}$ ，and $m(2,3) c^{2}$ ．Across the three $0.5 \mathrm{C}_{1}$ el－ ementary fermions and the three 0.5 R elementary fermions，$m(0,3) c^{2}$ would pertain once，$m(1,3) c^{2}$ would pertain twice，$m(2,3) c^{2}$ would pertain twice， and $m(3,3) c^{2}$ would pertain once．Regarding $d^{\prime}(0)$ ， one might consider the possibility that the following two notions pertain．$d^{\prime}(0)=+1 / 2^{2}$ pertains．The notion of anomalous property pertains．

## 3．1．7．We discuss possible masses for heavy neutri－ nos．

For purposes of estimating or calculating masses，the known neutrinos associate with a value of $l_{m}$ for which $-6 \leq l_{m} \leq-3$ ．Charged leptons associate with $0 \leq l_{m} \leq 3$ ．If heavy neutrinos as－ sociate with $6 \leq l_{m} \leq 9$ ，a lower bound on rest energies for heavy neutrinos might be $m(6,3) c^{2} \sim$ $6 \times 10^{3} \mathrm{GeV}$ ，which might be large enough to com－ port with limits that associate with observations． （References［12］and［13］discuss limits that obser－ vations may set．People have not detected 0.5 M particles．）To the extent the lower bound asso－ ciates with $m(6,3 / 2) c^{2}$ ，the lower bound would be $\sim 2.5 \times 10^{9} \mathrm{GeV}$ ．

Table 15: Rest energies that may pertain regarding the 0.5 N neutrinos.

| Topic | Note |
| :---: | :---: |
| $l_{m}=-1$ | $m(-1,3)=m(-1,3 / 2)$ - Comports with the equation underlying the main calculation regarding the masses of charged elementary fermions. |
| Assumption | $m\left(l_{m}, 3 / 2\right)$ pertains - regarding elementary fermions - for $l_{m} \leq-1$. |
| Neutrinos | We suggest masses for the three 0.5 N neutrinos. <br> People suggest - based on observations - that the sum of the three neutrino rest energies is at least approximately 0.06 eV and not more than approximately 0.12 eV . We note two possibilities. <br> - $m c^{2}=m(-4,3 / 2) c^{2} \approx 3.4 \times 10^{-2} \mathrm{eV}$ pertains for each of the three neutrinos. - $m c^{2}=m(-4,3 / 2) c^{2} \approx 3.4 \times 10^{-2} \mathrm{eV}$ pertains for each of two neutrinos. For one neutrino, one of $m(-6,3 / 2) c^{2} \approx 4.2 \times 10^{-6} \mathrm{eV}$ and $m(-5,3 / 2) c^{2} \approx 4.4 \times 10^{-4} \mathrm{eV}$ might pertain. |
| Neutrinos | We suggest aspects regarding possible differences between mass eigenstates and interaction eigenstates for the three 0.5 N neutrinos. <br> Regarding interactions between some LRI bosons and an elementary fermion, the following notions pertain. Interactions between 2 L and an elementary fermion conserve the mass of the elementary fermion, but do not necessarily conserve the flavour of the elementary fermion. Interactions between 3L and an elementary fermion conserve the flavour of the elementary fermion, but do not necessarily conserve the mass of the elementary fermion. Interactions between 4 L and an elementary fermion do not necessarily conserve the mass of the elementary fermion or the flavour of the elementary fermion. <br> Standard Model notions of mass mixing might associate with effects that associate with the 4 g " PROP solution $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3$ and some of the CURR solutions $4 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{x}$, $4 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 6$, and $4 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 8$. (See table 5 and table 10 ) <br> A QFT (or, quantum field theory) that includes gravity might - within a limit (regarding $\Sigma \mathrm{L}$ ) of $\Sigma \leq 2$ - associate mass mixing with a notion of anomalous gravitational properties and might successfully estimate aspects that people associate with mass mixing. |

Table 16: Rest energies that might pertain regarding the suggested 0.5 R arcs.

| Topic | Note |
| :--- | :--- |
| Arcs | Our work suggests (but does not necessarily require) some specific masses for the three arc |
|  | particles. |
|  | $l_{q}=0-$ This notion comports with the notion - for arcs - that $Q=0$. |
|  | $m\left(l_{m}, 0\right)=m\left(l_{m}, 1\right) \cdot\left(m\left(l_{m}, 1\right) / m\left(l_{m}, 2\right)\right)$ - This essay assumes this equation. |
|  | $m(0,0) c^{2} \approx 10.7 \mathrm{MeV}, m(1,0) c^{2} \approx 6.8 \mathrm{MeV}, m(2,0) c^{2} \approx 102 \mathrm{MeV}$. |

3.1.8. We discuss a possible limit regarding the spins of LRI elementary particles.
Table 17 suggests the possibility that - for LRI elementary particles $\Sigma \mathrm{L}-\Sigma$ might be no greater than four.

A limit - for LRI elementary particles $\Sigma \mathrm{L}$ - of $\Sigma \leq 4$ seems to be consistent with other aspects of our modeling.

### 3.2. Dark matter

This unit suggests specifications for dark matter.

### 3.2.1. We discuss - for the six isomers -elementary-fermion masses, flavours, and handedness.

Regarding each $l_{I}$ that is at least one, we assume that the elementary particles in isomer $l_{I}$ match with respect to mass - the elementary particles in isomer zero.

For $0 \leq l_{I} \leq 5$, we associate the quarks in isomer $l_{I}$ with three values of $l_{m}$. (See table 13 and table 14 .) The values are $3 l_{I}+0,3 l_{I}+1$, and $3 l_{I}+2$. Across the six isomers, quarks associate with each value of $l_{m}$ that is in the range $0 \leq l_{m} \leq 17$. Regarding quarks and flavours, we assume that within isomer $l_{I}$ - flavour 1 associates with $l_{m}=3 l_{I}$, flavour 2 associates with $l_{m}=3 l_{I}+1$, and flavour 3 associates with $l_{m}=3 l_{I}+2$.

Aspects of table 13 and table 14 point to the possibility that means for matching flavours and masses for charged leptons do not match means for matching flavours and masses for quarks. For charged leptons, isomer zero does not have a charged lepton that associates with $l_{m}=1$ and does have a charged lepton that associates with $l_{m}=3$. We assume that - for each $l_{I}$ - a charged lepton associates with each of $l_{m}=3 l_{I}+0, l_{m}=3 l_{I}+2$, and $l_{m}=3 l_{I}+3$.

We assume that - for each isomer $l_{I}$ such that $1 \leq l_{I} \leq 5$ - the charged-lepton flavour that associates with $l_{m}=3\left(l_{I}\right)+0$ equals the flavour that associates with the isomer $l_{I}-1$ charged lepton that associates with the same value of $l_{m}$ and - thus with $l_{m}=3\left(l_{I}-1\right)+3$. We assume that across the six isomers, one cyclical order pertains regarding flavours for charged leptons.

Table 18 shows, for isomers of charged elementary fermions, matches between masses and flavours.

Beyond the topic of flavours, the topic of handedness exists. Ordinary matter associates with lefthandedness. Our modeling suggests the possibility - and we assume - that isomers 0,2 , and 4 associate with left-handedness and that isomers 1,3 , and 5 associate with right-handedness.

### 3.2.2. We prepare to discuss the evolution of stuff

 that associates with each isomer.We associate the symbol OMSE with all SRI elementary particles and all ELF elementary particles except 0.5 M and 0.5 R elementary particles. OMSE abbreviates the three-element phrase ordinary-matter-similar elementary particles. We associate the symbol DMAI with the 0.5 M and 0.5R elementary particles. DMAI abbreviates the five-word phrase dark matter regarding all isomers. DMAI associates with the notion that - regarding isomer zero - these particles measure as being dark matter and do not measure as being ordinary matter.

We use the three-element term isomer number stuff to denote objects (including SRI elementary particles, ELF elementary particles, hadron-like particles, clumps of stuff, and stars) that associate with the isomer number set of simple elementary particles.
0.5R particles model as entwined. (See table 9.) We suggest that - at least after the inflationary epoch - 0.5R-based stuff consists of hadron-like particles. Each 0.5R-based-stuff hadron-like particle includes gluons and at least two arcs. (We deemphasize discussing roles that jay bosons might play.) Our work does not suggest an extent to which 0.5R-based stuff might form primordial black holes. Our work does not necessarily suggest that a two-or-three-hadron hadron-like particle can include both at least one quark and at least one arc.
0.5 M particles model as free. (See table 9 .) Our work does not suggest an extent to which $0.5 \mathrm{M}-$ based stuff might form primordial black holes.

Regarding each one of the six isomers, we suggest that stuff made from DMAI behaves within bounds for dark matter that associate with concordance cosmology.

### 3.2.3. We discuss - for each dark matter isomer -

 the evolution of stuff that associates with that isomer.We discuss the evolution of isomer 1, 2, 4, and 5 OMSE stuff.

Here, we use the two-word term alt isomer to designate an isomer other than isomer zero and isomer three.

A charged baryon that includes exactly three flavour 3 quarks is more massive than the counterpart zero-charge baryon that includes exactly three flavour 3 quarks. (For example, two tops and a bottom have a larger total mass than do one top and two bottoms.) Alt isomer flavour 3 charged leptons are less massive than isomer zero flavour 3 charged leptons. When flavour 3 quark states are much populated (and based on interactions mediated by W bosons), the alt isomer converts more charged baryons to zero-charge baryons than does

Table 17: The possibility that - for LRI elementary particles $\Sigma L-\Sigma$ might be no greater than four.

| Topic | Note |
| :--- | :--- |
| $l_{m}=18$ | $\left(\left(q_{e}\right)^{2} /\left(4 \pi \varepsilon_{0}\right)\right) /\left(G_{N}(m(18,3))^{2}\right)=4 / 3$. |
| Monopole properties | A force strength factor of 4 seems to associate with 1 g 1 and a force strength |
|  | factor of 3 seems to associate with 2 g 2 . (See, above, the equation |
|  | $\left.(4 / 3) \times\left(\beta^{2}\right)^{6}=\left(\left(q_{e}\right)^{2} /\left(4 \pi \varepsilon_{0}\right)\right) /\left(G_{N}\left(m_{e}\right)^{2}\right).\right)$ Possibly, other force strength |
|  | factors would be 2 for $3 \mathrm{~g} 3,1$ for 4 g 4 , and 0 (or, zero) for 5g5. Possibly, the |
|  | notion of zero force strength regarding 5 g 5 associates with a lack of relevance |
|  | for (and a lack of monopole properties that would associate with) solutions |
|  | $\Sigma \mathrm{g} \Sigma$ for which $\Sigma \geq 5$ and with a lack of LRI elementary particles $\Sigma \mathrm{L}$ for |
|  | which $\Sigma \geq 5$. |

Table 18: Matches between masses and flavours, for isomers of charged elementary fermions. The symbol $0.5 \mathrm{Q}_{>0}$ denotes the pair $0.5 \mathrm{Q}_{1 / 3}$ and $0.5 \mathrm{Q}_{2 / 3}$. The symbol $l_{f}$ numbers the three flavours. (See table 13 .)

| Isomer | $l_{m}\left(0.5 \mathrm{Q}_{>0}\right)$ | Respective $l_{f}\left(0.5 \mathrm{Q}_{>0}\right)$ | $l_{m}\left(0.5 \mathrm{C}_{1}\right)$ | Respective $l_{f}\left(0.5 \mathrm{C}_{1}\right)$ |
| :--- | :--- | :--- | :--- | :--- |
| 0 | $0,1,2$ | $1,2,3$ | $0,2,3$ | $1,2,3$ |
| 1 | $3,4,5$ | $1,2,3$ | $3,5,6$ | $3,1,2$ |
| 2 | $6,7,8$ | $1,2,3$ | $6,8,9$ | $2,3,1$ |
| 3 | $9,10,11$ | $1,2,3$ | $9,11,12$ | $1,2,3$ |
| 4 | $12,13,14$ | $1,2,3$ | $12,14,15$ | $3,1,2$ |
| 5 | $15,16,17$ | $1,2,3$ | $15,17,18$ | $2,3,1$ |

isomer zero. Eventually, in the alt isomer, interactions that entangle multiple W bosons result in the alt isomer having more neutrons and fewer protons than does isomer zero. The sum of the mass of a proton and the mass of an alt isomer flavour 1 charged lepton exceeds the mass of a neutron. Compared to isomer zero neutrons, alt isomer neutrons scarcely decay. The IGM (or, intergalactic medium) that associates with the alt isomer scarcely interacts with itself via electromagnetism.

We discuss the evolution of isomer three OMSE stuff.

The following two possibilities pertain. The evolution of isomer three OMSE stuff parallels the evolution of ordinary matter (or, isomer zero OMSE stuff). The evolution of isomer three OMSE stuff does not parallel the evolution of ordinary matter (or, isomer zero OMSE stuff). The second possibility might associate with - for example - a difference in handedness - with respect to charged leptons or with respect to W bosons - between isomer three and isomer zero. (See discussion related to table 18.)

### 3.3. Formation and evolution of the universe

This unit suggests eras - two of which would precede cosmic inflation - in the rate of expansion of the universe and suggests mechanisms that associate with the eras.

### 3.3.1. We discuss perspective regarding the rate of expansion of the universe.

Concordance cosmology points to three eras in the rate of expansion of the universe. The eras fea-
ture, respectively, rapid expansion; continued expansion, with the rate of expansion decreasing; and continued expansion, with the rate of expansion increasing.

This essay suggests using the notion of eras regarding the separating from each other of clumps - that, today, people would consider to be large of stuff. Examples of such clumps might include galaxy clusters and possibly even larger clumps.

### 3.3.2. We provide perspective regarding long-range

 interactions between objects.As two objects move away from each other, the relative effect of an RDF $\Xi^{-(k+1)}$ component decreases compared to the effect of an RDF $\Xi^{-k}$ component. One might associate the two-word phrase time period with a time range in which an RDF $\Xi^{-l_{r}}$ component provides dominant effects. Assuming that objects move away from each other and that one time period associates with $\Xi^{-(k+1)}$ and another time period associates with $\Xi^{-k}$, the time period that associates with $\Xi^{-(k+1)}$ comes before the time period that associates with $\Xi^{-k}$. Two smaller objects (such as galaxies) transit similar time periods more quickly than do two larger objects (such as galaxy clusters).

### 3.3.3. We discuss known and suggested eras in the history of the universe.

Table 19 discusses eras in the rate of separating of large clumps. (For discussion about possibilities leading up to a Big Bang, see reference [14.) For discussion about the possible inflationary epoch, see references [15], [6], [16], and [17]. For data and discussion about the two multi-billion-years eras, see

Table 19: Eras regarding the rate of separating of large clumps. The rightmost two columns suggest eras. (Table 20 discusses aspects that associate with each one of some eras.) In table 19 subsequent rows associate with later eras. The word inflation names the era that associates with the third row in the table. Regarding eras that would precede inflation, our modeling points to the possibility for the two eras that the table discusses. Concordance cosmology suggests inflation and the next two eras. Regarding inflation, people hypothesize this era. People suggest that the inflationary era started about $10^{-36}$ seconds after the Big Bang. People suggest that the inflationary era ended between $10^{-33}$ seconds after the Big Bang and $10^{-32}$ seconds after the Big Bang. Possibly, no direct evidence exists for this era. Observations support the notions of the two billions-of-years eras. TBD denotes to be determined. The symbol $\dagger$ denotes a possible association between the relevant era and the notion of a Big Bang. The leftmost four columns describe phenomena that our modeling suggests as noteworthy causes for the eras. (Regarding phenomena that associate with gravitation, table 19 echoes aspects including aspects regarding attraction and repulsion - that table 6 and table 10 show.) An RDF associates with the PROP solutions. Generally, a noteworthy cause associates with notions of acceleration. Generally, an era associates with a range of velocities. The symbol $\rightarrow$ associates with the notion that a noteworthy cause may gain prominence before an era starts.

| Force | PROP solutions | RDF | $\rho_{I}$ | $\rightarrow$ | Rate of separating | Duration |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| Attractive | $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 6^{‘} 8 \mathrm{x}$ | $\Xi^{-6}$ | 6 | $\rightarrow$ | Is negative | TBD |
| Repulsive | $0 \mathrm{~g} 1^{‘} 3^{‘} 4^{〔} 6$ | - | 1 | $\rightarrow$ | Turns positive $\dagger$ | TBD |
| Repulsive | $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 4 \mathrm{x}$ | $\Xi^{-5}$ | 1 | $\rightarrow$ | Increases rapidly | Fraction of a second |
| Attractive | $2 \mathrm{~g} 1^{‘} 2^{‘} 3$ | $\Xi^{-4}$ | 1 | $\rightarrow$ | Decreases | Billions of years |
| Repulsive | $22^{‘} 4$ | $\Xi^{-3}$ | 2 | $\rightarrow$ | Increases | Billions of years |
| Attractive | 2 g 2 | $\Xi^{-2}$ | 6 | $\rightarrow$ | Would decrease | - |

references [18], [19], [20], and [21]. For discussion of attempts to explain the rate of expansion of the universe, see reference [22].)

Table 20 suggests details regarding eras to which table 19 alludes.

Presumably, effects that associate with PROP solutions $2 \mathrm{~g} 1^{\prime} 2^{‘} 3^{‘} 4 \mathrm{x}$ and $2 \mathrm{~g} 2^{\prime} 4$ associate with concordance cosmology notions of dark energy pressures.

### 3.3.4. We discuss aspects regarding interactions between isomers and regarding possibilities for

 conversions between dark matter and ordinary matter.Our work suggests that - for times after early in the history of the universe - the following notions pertain.

Stuff associating with any one isomer interacts with stuff associating with a second isomer via gravitational interactions that associate with components of 2 L for which the reaches are six and, for some pairs of isomers, via gravitational interactions that associate with components of 2 L for which the reaches are two. Some interactions between isomers associate with components of 1L (or, electromagnetism) for which the reaches are six or two.

Some conversions between dark matter stuff and ordinary matter stuff might involve isomer zero dark matter. Our work does not suggest that nature includes leptoquarks. Conversions between heavy neutrinos and leptons might be possible. Conversions between hadron-like particles that contain 0.5 R particles and hadron particles (such as neutrons) might be possible.

Such conversions might be improbable and might be hard to detect.

Beyond such conversions, the following notions might pertain. Cross-isomer conversions might as-
sociate with effects of high-energy photons. Ordinary matter would not decay into dark matter and dark matter would not decay into ordinary matter.

### 3.3.5. We discuss aspects that might associate with baryon asymmetry (or, matter-antimatter asymmetry).

Our work suggests that isomer zero associates with left-handedness and that isomer three associates with right-handedness. All aspects that associate with isomer zero and with a reach (or, $\rho_{I}$ ) of two couple isomer zero and isomer three.

This essay suggests, but does not further explore, that a notion of matter-antimatter symmetry might pertain across the combination of isomer zero and isomer three.

### 3.3.6. We discuss aspects that might associate with an upper bound on the dark energy density of the universe.

Work above associates with values of $l_{\max }$ that are no larger than eight. For larger values of $l_{\max }$, the notion of a 9 g 9 solution might have relevance. Mathematically, per equation (3), the reach of 9 g 9 might be one-sixth.

Such math might suggest a notion of six isomers of the known six-isomer universe.

Possibly, during at least one of the first era and the second era that table 19 suggests (and possibly during later eras, possibly including the current era), the notion of entwined associates with interactions between our six-isomer universe and the other would-be five six-isomer universes.

Possibly, observations of dark energy densities associate with effects of such interactions.

Such notions suggest an upper bound of five on the ratio of dark energy density of the universe to the sum of all other densities of the universe.

Table 20：Details regarding eras regarding the rate of separating of large clumps．Table 19 discusses the eras．Each of the symbols $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{x}$ and $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{y}$ denotes either or both of $2 \mathrm{~g} 1^{\prime} 2^{\prime} 33^{\prime} 4 \mathrm{v}$ and $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 4 \mathrm{w}$ ．Our work does not necessarily specify the elementary fermions for which isomers form during the era that associates with the two－word phrase is negative．To the extent that the first significant appearance of most known elementary particles occurs during or just after the inflationary era，our work suggests that isomers of at least one of 0.5 M and 0.5 R form during the era that associates with the two－word phrase is negative．The symbol $\dagger$ associates with some aspects for which the involvement of 0.5 M or 0.5 R might pertain． We base 2 L －aspects of table 20 on table 10 ．Other $2 \mathrm{~g} \Gamma$ solutions might also pertain．（Perhaps，contrast table 4 with table 10 ）

| Rate of separating | Note |
| :---: | :---: |
| Is negative | Possibility： $2 \mathrm{~g} 1{ }^{\prime} 2^{\prime} 3^{6} 6^{\prime} 8 \mathrm{x}$ and their compacting of＂some form of energy＂lead to conditions suitable for the universe to form and evolve． <br> Possibility：The value of six for $\rho_{I}$（for $2 \mathrm{~g} 1^{‘} 2^{‘} 3^{‘} 6^{〔} 8 \mathrm{x}$ ）associates with setting up a system for which roughly equal creation of isomers pertains． <br> Possibility：Isomers of some elementary fermions and of 1J form．$\dagger$ <br> Possibility：The following interactions might characterize this era．For each interaction，the net circular polarization for each of before and after the interaction might be zero．$\dagger$ Presumably，the formation of gluons（or，1（1）G） could associate with the formation of arcs（or， $0.5(1) \mathrm{R})$ ）． <br> － $2(6) \mathrm{g} 1^{‘} 2^{‘} 3^{‘} 6^{‘} 8 \mathrm{x}+2(6) \mathrm{g} 1^{‘} 2^{{f2fd8865d-9ca1-406e-87e5-b670e6dbca8f}} 8 \mathrm{x} \rightarrow 0.5(1) \mathrm{M}+0.5(1) \mathrm{M}$ ． <br> －2（6）g1‘ $2^{\prime} 3^{‘} 6^{‘} 8 \mathrm{x}+2(6) \mathrm{g} 1^{\prime} 2^{\prime} 3^{\prime} 6^{\prime} 8 \mathrm{x} \rightarrow 0.5(1) \mathrm{R}+0.5(1) \mathrm{R}$ ． <br> －2（6）g1‘ $2^{\prime} 3^{\prime} 6^{{fc36baa8e-4a73-4703-9912-8ccd03561ac4}} 4 \mathrm{y}$ ． |
| Decreases | － |
| Increases | － |
| Would decrease | This essay does not try to explore the possibility that（or to estimate a time at which）a transition－for the largest observable objects－from repulsion based on $2 \mathrm{~g} 2^{\prime} 4$ to attraction based on 2 g 2 might occur． |

Concordance cosmology suggests a ratio of dark energy density of the universe to the sum of all other densities of the universe that is approximately 2.2 . (Reference [9] provides the data that we use.)

This essay does not further consider possible relationships (or possible lack of close relationships) between the notion of dark energy density and notions of dark energy pressure. This essay does not further discuss the topic of dark energy density of the universe.

### 3.4. Formation and evolution of galaxies

This unit suggests that our notions regarding long-range interactions and our specifications for dark matter combine to provide insight regarding galaxy formation and galaxy evolution.

### 3.4.1. We suggest aspects regarding events leading to the formation of a galaxy.

Reference [23] suggests that galaxies form around early clumps of stuff. The reference associates the word halo with such clumps.

Table 19 suggests that single-isomer stuff - such as stuff that features 0.5 R particles - forms during an era in which PROP solutions $2 \mathrm{~g} 1^{\prime} 2^{\prime} 3^{\prime} 6^{\prime} 8 \mathrm{x}$ which associate with attraction - dominate regarding prototype large clumps. Smaller-scale clumps might form before larger-scale clumps. Effects that associate with the PROP solution $2 \mathrm{~g} 1 ‘{ }^{\prime}{ }^{\prime} 3$ - which is attractive - might contribute to the formation of smaller-scale clumps. The reach that associates with $2 \mathrm{gl}^{\prime} 2^{\prime} 3$ is one.

We suggest that each one of many early halos associates with one isomer. We associate with such early halos the three-element term one-isomer original clump. We know of no reason why the six isomers would not form such clumps approximately equally. (Some concordance cosmology suggests that known elementary fermions form early in the era in which effects that associate with $2 \mathrm{~g} 1^{‘} 2^{〔} 3$ dominate regarding large-scale phenomena. Per remarks above, we suggest that that era starts after the formation of halos. Also, we suggest that our scenario does not necessarily depend on whether or when 0.5 M particles first form.)

Table 21 discusses suggestions regarding the formation and early evolution of a galaxy for which a notion of a one-isomer original clump pertains.

Presumably, some galaxies form based on two or more clumps, for which all of the clumps associate with just one isomer. Presumably, some galaxies form based on two or more clumps, for which some clumps associate with isomers that are not the same as the isomers that associate with some other clumps.

### 3.4.2. We suggest aspects regarding the evolution of galaxies.

We suggest two eras regarding the evolution of galaxies. The first era associates with the first two rows in table 21. The second era associates with the 2 g 2 attractive force that associates with the third row in table 21

Some galaxies do not exit the first era and do not significantly collide with other galaxies.

Many galaxies result from aspects associating with the 2 g 2 attractive force that associates with the third row in table 21. We discuss three cases. (Mixed cases and other cases might pertain.)

- Each one of some era one galaxies does not collide with other galaxies. Such a galaxy accumulates (via 2 g 2 attraction) stuff associating with various isomers that have representation in nearby IGM (or, intergalactic medium). The galaxy becomes an era two galaxy. The galaxy might include stuff that significantly associates with as many as five isomers.
- Each one of some era two galaxies merges (via 2g2 attraction) mainly just with galaxies that feature the same five isomers. The galaxy that merged, in effect, loses it status of being a galaxy. The resulting larger object is an era two galaxy. The galaxy might include stuff that significantly associates with as many as five isomers.
- Each one of some era one or era two galaxies merges (via 2 g 2 attraction) with other galaxies. The galaxy that merged, in effect, loses its status of being a galaxy. The resulting larger object is an era two galaxy. The galaxy might include stuff that significantly associates with as many as six isomers.


### 3.4.3. We suggest an explanation for the quenching of star formation within some galaxies and the stopping of the accrual of matter by some galaxies.

People report the notion that some galaxies seem to stop forming stars. (See reference [24] and reference [25].) Such quenching might take place within three billion years after the Big Bang, might associate with a relative lack of hydrogen atoms, and might pertain to half of the galaxies that associate with the notion of a certain type of galaxy. (See reference [25].) Reference [26] discusses a galaxy that seems to have stopped accruing both ordinary matter and dark matter about four billion years after the Big Bang.

We suggest that the quenching and the stopping of accruing nearby matter might associate with repulsion that associates with $2(2) \mathrm{g} 2^{\prime} 4$. Quenching might associate with galaxies for which original

Table 21: Stages and other information regarding the evolution of a galaxy for which a notion of a one-isomer original clump pertains. The table suggests stages, with subsequent rows associating with later stages. The rightmost one column describes aspects of the stage. The leftmost four columns in the table describe a component of 2 L that is a noteworthy cause for the stage. (Regarding phenomena that associate with gravitation, table 21 echoes aspects - including aspects regarding attraction and repulsion - that table 6 and table 10 show.) The symbol $\rightarrow$ associates with the notion that a noteworthy cause may gain prominence before a stage starts. Table 21 associates with a scenario in which a galaxy forms based on one original clump and does not significantly collide with other galaxies. The galaxy might retain some stuff that associates with the repelled isomer.

| Force | PROP solution | RDF | $\rho_{I}$ | $\rightarrow$ | Aspects of the stage |
| :--- | :--- | :--- | :--- | :--- | :--- |
| Attractive | $2 \mathrm{~g}^{〔} 2^{‘} 3$ | $\Xi^{-4}$ | 1 | $\rightarrow$ | A one-isomer original clump forms. |
| Repulsive | $2 \mathrm{~g} 2^{‘} 4$ | $\Xi^{-3}$ | 2 | $\rightarrow$ | The original clump repels (some) stuff that <br> associates with the isomer that associates with the <br> original clump and (most) stuff that associates <br> with one other isomer. |
| Attractive 2 g 2 | $\Xi^{-2}$ | 6 | $\rightarrow$ | The original clump attracts stuff that associates <br> with the four not-repelled isomers and stuff that <br> associates with the isomer that associates with the <br> original clump. |  |

clumps featured isomer zero stuff or isomer three stuff. The galaxy that reference [26] discusses might (or might not) associate with the notion of significant presence early on of one of isomers zero and three, one of isomers one and four, and one of isomers two and five. Such early presences might associate with a later lack of nearby stuff for the galaxy to accrue.

### 3.4.4. We suggest an explanation for some data regarding stellar stream GD-1 in the Milky

 Way galaxyData regarding stellar stream GD-1 suggest the possibility of effects from a yet-to-be-detected non-ordinary-matter clump - in the Milky Way galaxy with a mass of $10^{6}$ to $10^{8}$ solar masses. (For data and discussion regarding the undetected object, see references [27] and [28].) We suggest that the undetected object might be a clump of dark matter.

### 3.5. Ratios of dark matter effects to ordinary matter effects

This unit shows that our specification for dark matter seems to explain observed ratios of dark matter effects to ordinary matter effects.

Table 22 provides explanations for observed ratios of dark matter effects to ordinary matter effects. (For data and discussion regarding densities of the universe, see reference [9]. For data and discussion regarding galaxy clusters, see references [29], 30, 31, and [32. For data and discussion regarding absorption of CMB, see references [33], [34], and [35]. For data and discussion regarding observed early galaxies, see references [36] and [37]. Reference [36] influenced our choice of a time range to associate with the word early. For data and discussion regarding the combination of $0^{+}: 1$ and later, see references [38, [39, 40], 41, [42], and 43. For data and discussion regarding
observed dark matter galaxies, see references 23, [44, and 45]. Current techniques might not be capable of observing early dark matter galaxies. References [46] and [47] suggest, regarding galaxy clusters, the existence of clumps of dark matter that might be individual galaxies. Extrapolating from results that references [23] and [48] discuss regarding ultrafaint dwarf galaxies that orbit the Milky Way galaxy might suggest that the universe contains many DM:OM $1: 0^{+}$later galaxies. For data and discussion regarding galaxies for which DM:OM ratios of $\sim 4: 1$ pertain, see references 49 and 50 . For data and discussion regarding later galaxies for which DM:OM ratios of $5^{+}: 1$ pertain, see reference [23]. References 51] and 52] provide data about collisions of galaxies.)

## 4. Discussion

This unit discusses some possibilities regarding specific possible elementary particles and regarding dark matter, discusses some aspects regarding tensions regarding large-scale phenomena, and discusses aspects regarding properties that people infer from data or that people might be able to infer from data. This unit discusses relationships between our work and other work. Other work includes observational research and modeling-centric research. This unit provides a visual recap of some results from our work.

### 4.1. Some hypothesized elementary particles

This unit discusses possibilities regarding the existence and properties of some hypothesized elementary particles.

### 4.1.1. We discuss the possible existence of axions.

This essay seemingly does not suggest an elementary boson that would associate with notions of

Table 22: Explanations for observed ratios of dark matter effects to ordinary matter effects. DM denotes dark matter. OM denotes ordinary matter. DM:OM denotes a ratio of dark matter effects to ordinary matter effects. Inferences of DM:OM ratios come from interpreting data. Regarding densities of the universe, we assume that DMAI stuff associates with the plus in DM:OM $5^{+}: 1$. Stuff - other than DMAI stuff - that associates with isomers one through five associates with the five in DM:OM $5^{+}: 1$. Regarding some galaxy clusters, we assume that galaxy clusters (that have not collided with other galaxy clusters) associate with DM:OM ratios that are similar to DM:OM ratios for densities of the universe. The four-word phrase some absorption of CMB associates with the notion that people measured some specific depletion of CMB (or, cosmic microwave background radiation) and inferred twice as much depletion as people expected based solely on hyperfine interactions with hydrogen atoms. Possibly, half of the depletion associates with DM effects. We assume that isomer three hydrogen-like atoms account for the half of the absorption for which isomer zero (or, ordinary matter) hydrogen atoms do not account. The reach of an instance of $1 \mathrm{~g} 1 ‘ 4{ }^{〔} 6 ‘ 8$ is two isomers. (See table 10 ) The occurrence of $1 \in \Gamma$ associates with electromagnetism. The occurrence of $6 \in \Gamma$ associates with fermion aspects. The occurrences of $4 \in \Gamma$ and $8 \in \Gamma$ associate with spin and orbital angular momentum. (See table 4 and table 10 ) We assume that the PROP solution $1 \mathrm{~g} 1^{\prime} 4^{\prime} 6^{\prime} 8$ associates with (at least) hyperfine states. We assume that PROP 1 g 1 ' $2^{〔} 4^{4} 6{ }^{\prime} 8 \mathrm{x}$ solutions associate with (at least) hyperfine transitions. Regarding galaxies, the notion of early associates with observations that pertain to galaxies that people associate with (or, would, if people could detect the galaxies, associate with) high redshifts. High might associate with $z>7$ and possibly with smaller values of $z$. Here, $z$ denotes redshift. The word later associates with the notion that observations pertain to objects later in the history of the universe. The three-word phrase dark matter galaxy denotes a galaxy that contains much less ordinary matter than dark matter. Possibly, people have yet to directly detect early dark matter galaxies.

| Aspect | DM:OM | Comment |
| :---: | :---: | :---: |
| Densities of the universe | $5^{+}: 1$ | - |
| Some galaxy clusters | $5^{+}: 1$ | - |
| Some absorption of CMB | 1:1 | Half of the absorption might be via DM. |
| Some early galaxies | $0^{+}: 1$ | For each of some early galaxies, each original clump associates with isomer zero. Later, the galaxy might accumulate DM. |
| Some later galaxies | $0^{+}: 1$ | Some early DM:OM $0^{+}: 1$ galaxies survive (without significant collisions with galaxies for which DM:OM is not $\left.0^{+}: 1\right)$ until later times. |
| Some early galaxies | $1: 0^{+}$ | For each of some early galaxies, each original clump associates with an isomer other than isomer zero. Early on, the density of OM stars is small and people do not detect the galaxy. Later, the galaxy might accumulate enough OM to be visible. The term dark matter galaxy pertains. |
| Some later galaxies | $1: 0^{+}$ | Some early DM:OM 1: $0^{+}$galaxies survive (without significant collisions with galaxies for which the DM:OM is not $1: 0^{+}$) until later times. The term dark matter galaxy pertains. |
| Some later galaxies | $\sim 4: 1$ | An original clump might associate with any isomer other than isomer three. (Isomer three repels OM stuff.) Eventually, the galaxy accumulates enough stuff (that does not associate with the isomer that associates with the original clump) to have a DM:OM ratio that is somewhat near 4: 1 . |
| Many later galaxies | $5^{+}: 1$ | Over time, galaxies collide. Collisions tend to result in the formation of larger galaxies that include much stuff from smaller galaxies. A later galaxy that results from enough collisions is likely to associate with somewhat similar - across the six isomers - amounts of stuff from originally one- (or few-) isomer original clump galaxies. |

an axion. People suggest that - under some circumstances - axions might convert into photons. We suggest that observations that people might associate with effects of axions might instead associate with the difference between our notion of $1(6) \mathrm{g} 1^{‘} 2^{〔} 4$ and other modeling notions that might associate with notions of $1(1) \mathrm{g} 1^{\prime} 2^{\prime} 4$. Also, observations that people might associate with effects of axions might instead associate with interactions involving jay (or, 1J) bosons or aye (or, 1I) bosons. (See table 20.)

### 4.1.2. We discuss the possible existence of magnetic monopoles.

This essay does not suggest an elementary particle that would associate with notions of a magnetic monopole. Table 2 and table 4 seem not to suggest a 1L interaction with a monopole other than an electric monopole.

### 4.1.3. We discuss the possible existence and possible properties of right-handed $W$ bosons

Reference [53] discusses a fraction of decays of ordinary matter top quarks for which the decay products include W bosons - that might produce right-handed W bosons. The fraction, $f_{+}$, is $3.6 \times 10^{-4}$. Reference 9 provides a confidence level of 90 percent that the rest energy of a would-be $W_{R}$ (or, right-handed W boson) exceeds 715 GeV . (Perhaps, note also, reference 54.)

Our work suggests that $W_{R}$ bosons associate only with isomers one, three, and five. Our work suggests possibilities for inter-isomer interactions and conversions. (See discussion related to table 19 and table 20.)

We explore a notion that aspects of our modeling might approximately reproduce the above result that Standard Model modeling suggests.

Aspects related to table 14 and table 18 suggest values of calculated masses that do not associate with masses of known or suggested elementary particles. For example, our modeling does not suggest that $m(5,3)$ associates with the inertial mass of an isomer one charged lepton. However, perhaps such mass-like quantities associate with some measurable aspects of nature. For charged leptons and $0 \leq l_{I} \leq 4$ and $0 \leq l_{I}^{\prime} \leq 2, m\left(3\left(l_{I}+1\right)+l_{I}^{\prime}, 3\right)=$ $\beta m\left(3\left(l_{I}+0\right)+l_{I}^{\prime}, 3\right)$. One might conjecture that isomer zero observations of some aspects of isomer one phenomena associate with notions of non-inertial mass-like quantities that are $\beta$ times the inertial masses for isomer zero elementary particles (and that are $\beta$ times inertial masses for the counterpart isomer one elementary particles).

Based on notions of scaling that might calculate non-inertial mass-like quantities (and ignoring some of our suggestions regarding lack of interisomer phenomena), one might conjecture that our modeling suggests that $f_{+} \sim e^{\left(\beta^{-1}\right)}-1 \approx \beta^{-1} \approx$
$2.9 \times 10^{-4}$. This estimate might not be incompatible with results that reference 53 discusses. A notion of $m_{\text {non-inertial }, \mathrm{W}_{R} \text { isomer one }} c^{2}=\beta m_{\mathrm{W}} c^{2} \approx$ $2.8 \times 10^{5} \mathrm{GeV}$ might pertain. Here, the notion of non-inertial mass-like quantity might associate with inferences that associate with 1 L or $1 \mathrm{~W}_{1}$ and do not associate directly with 2 L .

### 4.2. Interactions involving the jay boson

This unit discusses interactions that involve jay bosons.

### 4.2.1. We discuss interactions - that involve jay

 bosons - that might take place before or during inflation.We consider interactions in which two jay bosons move in parallel, interact, and produce one aye boson plus something else. Here, we assume that conservation of angular momentum pertains and that one can de-emphasize angular momentum that is not intrinsic to the relevant elementary particles. We consider two cases. In the first case, the two jay bosons have the same (one of either right or left) circular polarization. Conservation of angular momentum allows an outgoing combination of one 2 L particle and one 0I particle. Conservation of angular momentum might preclude producing one 1L particle and one 0 I particle. In the second case, one jay boson has left circular polarization and the other jay boson has right circular polarization. Conservation of angular momentum allows the production of two 0I particles. Conservation of angular momentum might preclude producing one 1 L particle and one 0 I particle.

The two cases might comport with the notion that gravitation can be significant during inflation. The two cases might comport with the notion that jay bosons form before aye bosons form. (See table 20.)

The two cases might comport with a (possibly not relevant) notion that electromagnetism might become significant essentially only after inflation.

### 4.2.2. We discuss the notion of Pauli repulsion.

Physics includes the notion that two identical fermions cannot occupy the same state. Regarding some modeling, one notion is that repulsion between identical fermions associates with overlaps of wave functions. Another notion features wave functions that are anti-symmetric with respect to the exchange of two identical fermions.

Our modeling might be compatible with such aspects of other modeling and, yet, not necessitate for kinematics modeling - the use of wave functions. Modeling based on jay bosons might suffice.

Modeling based on jay bosons might suggest that prevention of two identical fermions from occupying the same state might associate with, in effect,
trying to change aspects related to the fermions. Notions of changing a spin orientation or, for elementary fermions, changing a flavour might pertain.

### 4.2.3. We discuss Pauli crystals.

Reference 55] and reference 56] report detection of Pauli crystals. We suggest that modeling based on the notion of jay bosons might help explain relevant phenomena.

### 4.2.4. We discuss a possible discrepancy - regarding energy levels in positronium - between modeling and observation.

Reference 57] and reference [58] discuss the transition - between two states of positronium characterized by the expression $2^{3} S_{1} \rightarrow 2^{3} P_{0}$. People discuss the energy that associates with the transition. Four standard deviations below the nominal observed value of energy approximately equals four standard deviations above the nominal value of energy that modeling suggests.

Perhaps, notions regarding jay bosons extend to explain the might-be discrepancy regarding positronium. Compared to modeling that people use, a new notion of virtual charge exchange or a new notion of virtual flavour change might pertain.

To the extent that modeling that people use does not suffice, modeling related to the jay boson might help to close the gap between observation and modeling.

### 4.3. Constraints regarding dark matter

This unit discusses the extent to which our notion of dark matter comports with constraints about the nature of dark matter - that people associate with data about dark matter or with outputs from models - that people use - that have bases in assumptions about dark matter.

### 4.3.1. We discuss aspects related to cosmological models.

Reference [23] summarizes some thinking about constraints on dark matter and about notions of dark matter. The article notes that CDM (or, cold dark matter) might comport well with various models. Some models associate with the one-element term $\Lambda$ CDM. The article notes that people have yet to determine directly whether nature includes CDM stuff. The article notes that people consider that notions of SIDM (or, self-interacting dark matter) might be appropriate regarding nature. People also use other terms, such as the three-word term warm dark matter, to note possible attributes of dark matter. For example, reference [59] suggests that notions of WDM (or, warm dark matter) might reduce discrepancies between data regarding clustering within galaxies and modeling that associates with CDM. Notions such as SIDM and WDM arose
from modeling that differs from our modeling. We are reluctant to try to closely associate terms such as SIDM or WDM with our modeling. (We suggest that isomer zero 0.5 R -based stuff, isomer zero 0.5 M stuff, and all stuff associating with isomers one, two, four, and five might comport with some notions of CDM. We suggest that the remaining dark matter stuff - or, isomer three OMSE stuff - might associate with some notions of WDM and with some notions of SIDM.)

We suggest that our notion of dark matter is not necessarily incompatible with constraints - that have bases in cosmological models - on dark matter.

### 4.3.2. We discuss aspects related to collisions of pairs of galaxy clusters.

In particular we discuss the Bullet Cluster collision of two galaxy clusters. (Reference 60 discusses the Bullet Cluster.) Presumably, observations regarding other such collisions might pertain.

Observations suggest two general types of trajectories for stuff. Most dark matter - from either one of the clusters - exits the collision with trajectories consistent with having interacted just gravitationally with the other cluster. Also, ordinary matter stars - from either cluster - exit the collision with trajectories consistent with having interacted just gravitationally with the other cluster. However, ordinary matter IGM (or, intergalactic medium) from either cluster - lags behind the cluster's ordinary matter stars and dark matter. That ordinary matter IGM interacted electromagnetically with the other cluster's ordinary matter IGM, as well as gravitationally with the other cluster.

Our work suggests that - regarding each cluster - essentially all dark matter - except isomer three IGM - passes through without interacting significantly electromagnetically with stuff from the other cluster. Our work suggests that isomer three IGM that associates with each cluster might interact significantly with isomer three IGM that associates with the other cluster. Isomer three IGM might follow trajectories similar to trajectories for isomer zero IGM.

We are uncertain as to the extent to which observational data might suggest that the amounts of dark matter that lags the bulk of dark matter are sufficiently small that our nominal notions regarding isomer three IGM do not comport with observations.

Should the actual fraction of lagging dark matter be too small, we might need to reconsider the extent to which isomer three differs from isomer zero. We note some examples of possible reconsideration. For one example, possibly isomer three has right-handed elementary fermions but interactions involving such fermions model as retaining aspects of left-handed-centric interactions that asso-
ciate with isomer zero. For another example, possibly isomer three does not evolve adequately similarly to isomer zero. To the extent that isomer three adequately differs from or does not evolve similarly to isomer zero, our explanation regarding CMB depletion via - in part - interactions with dark matter hydrogen-like atoms might be inaccurate (for example, based on an inaccurate estimate of the number of isomer three hydrogen-like atoms).

We suggest that our notion of dark matter is not necessarily incompatible with constraints - that have bases in observations of collisions of galaxy clusters - on dark matter.

### 4.4. Tensions regarding large-scale phenomena

This unit suggests means to resolve tensions between data and modeling - regarding the rate of expansion of the universe and other large-scale phenomena.

### 4.4.1. We suggest an explanation for the notion that concordance cosmology underestimates recent increases in the rate of expansion of the universe.

Table 19 and table 20 discuss possible and known eras in the history of the universe.

People suggest that concordance cosmology modeling underestimates - for the second multi-billion-years era - increases in the rate of expansion of the universe. (See references [61, [62], [63], [64], [65], 66], 67], and 68].)

We suggest the following explanation for such underestimates.

When using modeling based on general relativity, people might try to extend the use of an equation of state (or use of a cosmological constant) that works well regarding early in the first multi-billionyears era. Regarding that time, our modeling suggests dominance by attractive effects that associate with the $2 \mathrm{~g} 1{ }^{\prime} 2^{‘} 3$ component of gravity. The notion of a reach of one pertains. The symbol $2(1) \mathrm{g} 1^{\prime} 2^{〔} 3$ pertains. Our modeling suggests that - later in the first multi-billion-years era - repulsive effects that associate with $2(2) \mathrm{g} 2^{‘} 4$ become significant. Dominance by $2(2) g 2^{‘} 4$ pertains by the time the second multi-billion-years era starts. However, people's use of an equation of state that has roots in the time period in which $2(1) \mathrm{g} 1^{‘} 2^{‘} 3$ dominates would - at best extrapolate based on a notion of $2(1) \mathrm{g} 2^{4} 4$ (and not a notion of $\left.2(2) \mathrm{g} 2^{‘} 4\right)$. That modeling would underestimate the strength of the key gravitational driver - of expansion - by a factor of two.

We point - conceptually - to the following possible remedy.

People might change (regarding the stressenergy tensor or the cosmological constant) the aspects that would associate with repulsion and the $2 \mathrm{~g} 2 \cdot 4$ component of gravity. The contribution - to
the pressure - that associates with $2 \mathrm{~g} 2^{4} 4$ needs to double (compared to the contribution that would associate with $\left.2(1) \mathrm{g} 2^{〔} 4\right)$.

### 4.4.2. We suggest an explanation for the notion that concordance cosmology overestimates

 large-scale clumping of matter.People suggest that concordance cosmology modeling overestimates large-scale clumping of matter - ordinary matter and dark matter. (For data and discussion, see references [69, [70, [71, and [64].)

We suggest that concordance cosmology modeling associates with a repulsive component $-2(1) \mathrm{g} 2^{`} 4$ - of gravity. Our modeling suggests that $2(2) \mathrm{g} 2^{`} 4$ pertains. (That is, for each instance of $2 \mathrm{~g} 2^{‘} 4$, a reach of two isomers pertains.) The additional (compared to concordance cosmology modeling) repulsion might explain the overestimating - of clumping, per concordance cosmology modeling - that people suggest.

### 4.4.3. We suggest an explanation for the notion

 that concordance cosmology might not account for some observations about effects within individual galaxies - of the gravity associated with nearby galaxies.People suggest that concordance cosmology modeling might not account for some observations about effects - within individual galaxies - of the gravity associated with nearby galaxies. (For data and discussion, see reference [72].)

We suggest that concordance cosmology modeling associates with a repulsive component - 2(1)g2‘4 - of gravity. Our modeling suggests that 2(2)g2‘4 pertains. The additional (compared to concordance cosmology modeling) repulsion might explain at least some aspects of the observations that people report.

### 4.5. Inferable properties

This unit discusses properties that people might infer from data.

Table 10 suggests properties (of objects) that people might infer from data. Table 10 has a basis in $\Sigma g^{\prime}$ solutions. Examples of properties include charge, nominal magnetic moment, mass, and isomer or isomers.

Table 5 suggests properties (of objects) that people might infer from data. Table 5 has a basis in $\Sigma g^{\prime \prime}$ solutions. An example of a property is anomalous magnetic moment.

Discussion - for example, regarding PROP-andCURR 4-somes - associates with the notion that people might infer the properties of velocity and angular velocity.

Table 9 points to the notion that people might infer - from data and for an object - numbers of each of the elementary fermions.

This essay de-emphasizes discussing the following notions - inferable properties that might associate with 0 g boson solutions (See table 8), inferable properties that might associate with properties of LRI fields, and other possible inferable properties to which this essay does not directly allude.

### 4.6. Associations with other work

This unit discusses associations between our work and other work. Other work includes observational research and modeling-centric research. This unit discusses briefly aspects of other work. This unit provides references - regarding other work to review articles and other information. This unit suggests context for associating our work with other work. This unit does not necessarily explore thoroughly relationships between our work and other work.
4.6.1. We discuss relationships - between our work and other work - regarding physics constants and physics properties.
People discuss possibilities for relationships between electromagnetism and gravity. For example, reference [1] explores notions of a coupling between electromagnetism and gravity. People discuss possibilities for modeling that blends modeling used regarding electromagnetism and modeling used regarding gravity. Reference [2] and reference [3] discuss Einstein-Maxwell equations that suggest combining electromagnetic stress-energy tensors and the Einstein field equations, which have origins in modeling regarding gravitation.

Our work seems to interrelate some physics constants. (See table 12 and table 14) Our work seems to interrelate some properties, including via modeling that catalogs physics properties. (See table 4 and table 10 )

We might offer new approaches to estimating some physics properties. This essay points to masses - that would comport with recent experimental results and that would have smaller standard deviations than standard deviations that associate with recent experiments - for each of the tau elementary fermion and the Higgs boson. (See respectively table 14 and table 12 .) Our work suggests - regarding the anomalous magnetic dipole moment of the tau elementary fermion - a possible estimate that might approximate a Standard Model estimate. (See discussion related to table 13 and table 14 )

### 4.6.2. We discuss relationships - between our work and other work - regarding modeling regarding physics properties.

Table 23 discusses approximate relationships between modeling that can deploy elementary-particle properties and aspects of our modeling.

We discuss the extent to which general relativity and the case of $n_{I}=6$ are mutually compatible.

To gain perspective, we consider two cases. Each case involves an isomer zero star and a planet. Across the two cases, the stars are identical and the planets are identical except that one planet associates with isomer zero and one planet associates with an alt (or, non-isomer-zero) isomer. The planets have identical non-spherical distributions of mass. The planets start out on identical orbits. We assume that, over time, each planet evolves toward spherical symmetry and, in so doing, sheds intrinsic energy. Via effects associating with the $2 \mathrm{~g} 1^{\prime} 2^{‘} 3$ PROP solution (and the related CURR solution and possibly solutions that cascade therefrom), the star has some influence on the isomer zero planet that the star does not have on the alt isomer planet.

In general, to the extent that modeling needs to have bases that associate (essentially) only with the PROP-and-CURR pair 2 g 2 and 2g2‘4, modeling based on general relativity might suffice. To the extent that modeling needs to consider effects that our work would associate with 2 L components that do not associate with a reach of six, problems might arise regarding overly relying on modeling based just on general relativity.

### 4.6.3. We discuss relationships - between our work and other work - regarding group theory.

Discussion leading to table 4 features notions of three degrees of freedom. Discussion regarding equation (3) suggests an association between the number of generators of the group $S U\left(2 n_{0}+1\right)$ and the reaches that associate with some solutions that associate with $n_{0}$.

We suggest that - for some aspects of modeling - three degrees of freedom, mathematics associating with two one-dimensional harmonic oscillators, and mathematics associating with the group $S U(2)$ associate with each other. (For integers $l_{g r}$ such that $l_{g r} \geq 2$, reference [73] interrelates mathematics associating with $l_{g r}$ one-dimensional harmonic oscillators and mathematics associating with the group $S U\left(l_{g r}\right)$.) Here, we consider that one oscillator might associate with boson-like excitations regarding a relevant aspect that associates with the relevant value of $-l_{o}$. The other oscillator might associate with boson-like excitations regarding a relevant aspect that associates with the relevant value of $+l_{o}$. The number of generators of the group $S U(2)$ is three.

Regarding equation (3) and $S U\left(2 n_{0}+1\right)$, we suggest that modeling might consider that each $l_{o}$ that associates with no contribution to $\Gamma$ associates with a contribution of two toward the $2 n_{0}+1$ and that a (previously hypothetical) notion of $l_{0}$ (or, $l_{\text {zero }}$ ) associates with a contribution of one toward the $2 n_{0}+1$.

Table 23: Approximate relationships between modeling that can deploy elementary-particle properties and aspects of our modeling. $n_{I}$ denotes a number - one or six - of isomers. Modeling that other people use associates with $n_{I}=1$. Each one of some of the items in the symbol column does not associate with a popular symbol. CNC associates with charge-current 4 -vectors and with Maxwell's equations. Compared to CNC, QED adds associations with magnetic fields created by other than charge currents and adds associations with anomalous magnetic moments. QCD associates with $1 \mathrm{G}, 0.5 \mathrm{Q}_{1 / 3}$, and $0.5 \mathrm{Q}_{2 / 3}$. We suggest the possibility that QCD might extend to associate with 0.5 R . The symbol PEF associates with the three-word phrase Pauli exclusion force. We suggest that PEF associates with 1 J , each $0.5 \Phi$ family, and fermions that are not elementary particles. WIP associates with $1 W_{1}$ and $1 Z$. The symbol $\dagger$ denotes a notion of a (currently) hypothetical analog to QED. Our modeling suggests that a modeling basis might need to encompass the notion of anomalous gravitational property and the notion of six isomers.

| Modeling | Range of $\Sigma$ | $l_{o}$ PROP | $n_{I}$ | Symbol |
| :--- | :--- | :--- | :--- | :--- |
| Newtonian gravity | 2 | 2 | 1 | NEW |
| Moments of inertia | 2 | $1-3$ | 1 | MOI |
| Electrostatics | 1 | 1 | 1 | EST |
| Charge-and-current 4-vectors | 1 | 1 | 1 | CNC |
| Quantum electrodynamics | 1,3 | $1,2,4,6,8$ | 1 | QED |
| Quantum chromodynamics | 0 | $1-4,6,8$ | 1 | QCD |
| Pauli exclusion force | 0 | $1-4,6,8$ | 1 | PEF |
| Weak-interaction phenomena | 0 | $1-4$ | 1 | WIP |
| Suggested by our modeling | $0-4$ | $1-4,6,8$ | 6 | PRM |
| Gravitational analog to QED $\dagger$ | 2,4 | $1-4,6,8$ | 6 | QGD |

### 4.6.4. We discuss relationships - between our work and other work - regarding elementary particles.

We discuss other work that tries to suggest new elementary particles.

Reference 74 lists some types of modeling that people have considered regarding trying to extend the elementary particle Standard Model, including trying to suggest elementary particles that people have yet to find. Reference 75 provides information about some of these types of modeling. References [76, [77], and [78] provide some information about modeling and about experimental results. Reference [9] provides other information about modeling and about experimental results. (See reviews numbered 86, 87, 88, 89, 90, and 94.)

We discuss possible elementary particles that people have yet to find, that we suggest, and that other people might suggest.

Reference [7] suggests the notions of dark matter charges and dark matter photons. We suggest dark matter isomers of charged elementary particles and, in effect, dark matter components - such as components associating with electrostatics and magnetostatics - of electromagnetism. (See discussion related to table 10.)

Reference [6] suggests the notion of a inflaton field. We suggest an inflaton elementary particle. (See table 11 and note the 0I boson.)

People suggest the notion of a graviton. (See, for example, reference [79].) We suggest a graviton. (See table 11.)

Reference [11] discusses notions of sterile neutrinos and heavy neutrinos. We suggest possible elementary particles that might associate with notions of heavy neutrinos. (See table 11.)

We discuss possible elementary particles that people have yet to find, that we suggest, and for which other people might suggest that modeling rules out possible existence.

Reference [8] notes that modeling based on QFT (or, quantum field theory) suggests that massless elementary particles cannot have spins that exceed two. Our work suggests a possible spin-three analog to the photon and the possible graviton. (See table 7.) Our work suggests a possible spin-four analog to the photon and the possible graviton. (See table 7.) Discussion related to table 11 suggests that our work might not be incompatible with notions that nature does not include zero-mass elementary bosons that have spins that exceed two.

We discuss possible elementary particles that people have yet to find and our modeling seems not to suggest.

A symmetry regarding Maxwell's equations suggests that nature might include magnetic monopoles. Reference [80] discusses theory. Reference 78 reviews modeling and experiments regarding magnetic monopoles. We suggest that nature might not include an interaction that would associate with magnetic monopoles. (See table 4.) Reference 81 discusses a search - for magnetic monopoles - that did not detect magnetic monopoles.

Reference 76 reviews modeling and experiments regarding axions. Reference [76] notes modeling that suggests that nature might include axions. We suggest that nature might not include axions. (See table 8.) We suggest that phenomena that people might attribute to axions might not associate with axions. One such phenomenon could be electromagnetic interactions between ordinary matter and dark matter based on, for example, the
$1 \mathrm{~g} 1^{\prime} 2^{`} 4$ component of electromagnetism. (See table 4 and table 10 .)

Reference [77] reviews modeling and experiments regarding leptoquarks. We suggest that nature might not include leptoquarks. (See table 8.)

We discuss neutrino masses.
Reference [11 discusses modeling and data about neutrino masses and neutrino oscillations.

We suggest neutrino masses. (See table 15) As far as we know, our modeling is not incompatible with data that reference 11 discusses. Future experimentation might help validate or refute aspects of our work regarding neutrinos.

We discuss gravitation.
Reference [82] discusses experimental tests of theories of gravity.

We suggest effects - associating with isomers of elementary particles and with reaches of components of gravity - that suggest that other modeling regarding gravity would not be adequately accurate for some circumstances. We are uncertain as to the extent to which aspects that reference [82], reference [83], or reference 84 discuss would tend to validate or refute aspects of our modeling that pertains to gravitation.

We use modeling - regarding gravity - that has some similarities to models that people associate with the term gravitoelectromagnetism. (References [85] and [86] discuss gravitoelectromagnetism.) Our modeling regarding gravity has some similarities to models that use classical physics perturbations regarding Newtonian gravity. (Reference [4] deploys modeling that associates with nonspherical distributions of mass.)

### 4.6.5. We discuss relationships - between our work and other work - regarding cosmology.

We think that - with some exceptions - our work does not necessarily suggest significant changes - to concordance cosmology - regarding the large-scale evolution of the universe. (References [87, [88, [89], and [90] review aspects of concordance cosmology. Reference [22] discusses attempts to explain the rate of expansion of the universe.)

Each exception that this essay discusses associates either with a possible aspect of nature for which people have no observations or with a known gap between observations and concordance cosmology.

One exception pertains regarding before inflation. One exception pertains regarding recent changes in the rate of expansion of the universe. In each case, we suggest noteworthy contributions by a gravitational force component for which each instance (of the component) has a reach that is greater than one isomer. (See table 10.) For times associating with between the two cases, we suggest dominance by gravitational force components that
have reaches of one isomer. For times associating with between the two cases, we do not propose significant incompatibilities between our work and large-scale concordance cosmology.

We discuss a possibility regarding times before inflation.

We think that no direct observations pertain. We suggest two eras before inflation. (See table 19.) The first of those two eras features aspects that the Standard Model and concordance cosmology do not include. (Reference [14] discusses possibilities leading up to a Big Bang. References [16] and 88 ] discuss inflation.) One aspect is the jay boson. (See table 11 and table 19.) The other aspect is the set of $2 \mathrm{~g} 1^{‘}{ }^{\prime} 3^{‘} 6^{‘} 8 \mathrm{x}$ components of gravity. (See table 19) An instance of each component has a reach of six isomers. For purposes of discussion, we assume that the universe transited those two eras. We assume that concordance cosmology can embrace the jay boson. For the first of those two eras, an extrapolation of concordance cosmology techniques might underestimate the strength of the key driver - the $2 \mathrm{~g} 1^{‘} 2^{\prime} 3^{‘} 6^{‘} 8 \mathrm{x}$ components of gravity - by a factor of six.

We discuss phenomena during and after the lead-up to the current multi-billion-years era of increases in the rate of expansion of the universe.

People suggest that concordance cosmology underestimates increases in the rate of expansion. (References 69, 61, 62, 63], and 64] discuss relevant notions.)

We think that we point to a basis for the underestimates. Regarding times before that lead-up, we suggest dominance by an attractive quadrupole gravitational force component $-2 \mathrm{~g} 1^{‘} 2^{‘} 3$. (See table 19.) Each instance of that force component has a reach of one isomer. Before and during the recent multi-billion-years era, the $2 \mathrm{~g} 2{ }^{〔} 4$ gravitational force component gains prominence and then becomes dominant. Each instance of $2 \mathrm{~g} 2^{\prime} 4$ has a reach of two isomers. We suggest that concordance cosmology models that work well regarding times for which reach-one dominance pertains would not necessarily work well after those times. We suggest that extrapolating based on such concordance cosmology modeling would underestimate (conceptually by a factor of two) the strength of the driver for increases in the rate of expansion. We suggest that - to get good results via concordance cosmology modeling - people might adjust the equation of state. In general, for each relevant density, components of pressure that associate with repulsion need to increase.

Our suggested resolution regarding the underestimate seems to differ from possible resolutions based on concordance cosmology modeling. Our suggested resolution focuses on phenomena that would pertain at the times for which concordance
cosmology modeling seems not to be adequate. Other possible resolutions might focus on phenomena early in the history of the universe. (See reference [89].)

### 4.6.6. We discuss relationships - between our work and other work - regarding astrophysics.

We think that our modeling is not necessarily incompatible with astrophysics data or with results based on concordance cosmology modeling. (Here, we assume that the two-word term concordance cosmology includes aspects that associate with dark matter, astrophysics, and effects of gravity on scales as small as one galaxy.)

We discuss properties of dark matter.
Reference 91 suggests the following notions. Most dark matter comports with notions of cold dark matter. Models that associate with the twoword term modified gravity might pertain; but - to the extent that the models suggest long-range astrophysical effects - such models might prove problematic. People suggest limits on the masses of basic dark matter objects. Observations suggest smallscale challenges to the notion that all dark matter might be cold dark matter. People use laboratory techniques to try to detect dark matter. People use astrophysical techniques to try to infer properties of dark matter.

We think that our modeling regarding dark matter comports with such notions. For astrophysical phenomena (and not necessarily regarding the rate of expansion of the universe), components - that have reaches other than six - of gravity play roles locally; however, the impacts do not extend to cosmological scales. The dark matter isomer that might evolve similarly to ordinary matter might provide bases for resolving some of the small-scale challenges.

We discuss observations and models regarding galaxy formation.

Reference 92 discusses galaxy formation and evolution, plus contexts in which galaxies form and evolve. Reference [92] discusses parameters by which people classify and describe galaxies.

We suggest that - regarding galaxies - observations of ratios of dark matter to ordinary matter might tend to cluster near some specific ratios. (See table 22.) Our modeling seems to explain such ratios.

Our modeling suggests that ratios of dark matter to ordinary matter might reflect fundamental aspects - of nature - that concordance cosmology modeling does not include. Here, a key aspect is that of isomers. (See table 22, )

Reference 92 seems not to preclude galaxies that have few ordinary matter stars. Reference 92 seems not to preclude galaxies that have little ordinary matter.

We think that dark matter to ordinary matter ratios that our modeling suggests are not necessarily incompatible with verified concordance cosmology modeling.

We discuss observations and models regarding interactions between galaxies.

Reference [72] suggests that concordance cosmology modeling might not adequately explain gravitational interactions between neighboring galaxies. We suggest that notions pertaining to reaches and isomers might help to bridge the gap between observations and concordance cosmology modeling.

We think that our work points to a possible opportunity to study harmony between results based on established kinematics models and results based on our notions of components of gravity.

### 4.7. Visual recap of results

This unit recaps some results from our work.
Figure 2 alludes to all known elementary particles and to elementary particles that our work suggests. (See table 11) Some successful modeling might not consider notions that our work associates with the would-be spin-three (or, 3L) elementary particle and with the would-be spin-four (or, 4L) elementary particle to associate with elementary particles. (See discussion related to table 11.)

Figure 3 shows quantitative ratios of dark matter effects to ordinary matter effects. (See table 22, Our work suggests quantitative explanations for the ratios. The explanations have bases in our specification for dark matter. The specification has roots in our hypothesis that associates with six isomers of known elementary particles. Each isomer of known elementary particles associates - approximately with its own photon. Each isomer of known elementary particles associates with its own $0 \mathrm{I}, 0.5 \mathrm{M}$, 0.5 R , and 1 J elementary particles.

Figure 4 interrelates elementary particles, isomers, ordinary matter, and dark matter.

Figure 5 suggests details about our explanation for the known ratio - five-plus to one - of dark matter density of the universe to ordinary matter density of the universe. (See table 22.)

Figure 6 interrelates isomers of elementary particles, components of gravity, eras in the evolution of the universe, and eras in the evolution of galaxies.

Figure 7 suggests eras in the evolution of the universe. (See table 19.) As far as we know, direct observations and data associate only with the two multi-billion-years eras.

Figure 8 suggests eras in the evolution of - and various ratios of dark matter to ordinary matter for - galaxies. (See table 21 and table 22.) People have observed galaxies that associate with each one


Figure 2: Known and suggested elementary particles. The figure uses popular symbols for the known elementary particles. The following sentences discuss symbols for and notions about suggested new elementary particles. (The number $S$ in a symbol $S \Phi$, associates with elementary-particle spin in units of $\hbar$.) 0.5 R associates with three spin-one-half zero-charge analogs to quarks. 1J associates with a spin-one zero-charge boson that associates with Pauli repulsion. 0I associates with a spin-zero inflaton. 2L associates with a possible spin-two graviton. 3L associates with a possible spin-three relative of the photon and the graviton. 4L associates with a possible spin-four relative of the photon and the graviton. 0.5 M associates with three spin-one-half heavy neutrinos.

| Ratios of dark matter effects to ordinary matter effects |  |  |
| :---: | :---: | :---: |
| (Ratios that our modeling explains) |  |  |
|  |  | Copyright © 2022 Thomas J. Buckholtz |
| Aspect | DM:OM | Comment |
| Densities of the universe | $5^{+}: 1$ | Observed |
| Some galaxy clusters | $5^{+}: 1$ | Observed |
| Some absorption of CMB | 1:1 | Data observed; Association with DM proposed |
| Some early galaxies | $0^{+}: 1$ | Observed |
| Some later galaxies | $0^{+}: 1$ | Observed |
| Some early galaxies | $1: 0^{+}$ | Possibly too difficult to observe directly |
| Some later galaxies | 1:0 ${ }^{+}$ | Observed |
| Some later galaxies | $\sim 4: 1$ | Observed |
| Many later galaxies | $5^{+}: 1$ | Observed |
| Legend: |  |  |
| DM - dark matter. OM - ordinary matter. |  |  |
| DM:OM - ratio of dark matter effects to | ordinary matter | effects. |

Figure 3: Ratios of dark matter effects to ordinary matter effects. For each ratio except two of the ratios, the following three sentences pertain. People have observed the ratio. People attribute the so-called dark matter effects to dark matter. Our modeling explains the ratio. Regarding some specific depletion of cosmic microwave background radiation (or, CMB), people have observed the ratio, some people speculate that the effects that people might not attribute to ordinary matter are effects of dark matter, and our modeling suggests that non-ordinary-matter effects are effects of dark matter. We use the two-word term early galaxies to include galaxies observed at redshifts of at least (and possibly somewhat less than) seven. Most relevant data about later galaxies pertains to galaxies observed at redshifts considerably less than seven. The three-word term dark matter galaxy pertains to a galaxy for which the DM:OM ratio (or, ratio of dark matter effects to ordinary matter effects) is one to zero-plus. Possibly, current techniques are not adequately sensitive to detect early dark matter galaxies.

# Elementary particles, isomers, ordinary matter, and dark matter <br> (Based on six isomers of all Standard Model elementary particles and of four types of proposed particles) 

Proposed particles: R - zero-charge analogs of quarks; M - heavy neutrinos; I - inflaton; J - spin-1 boson (Pauli repulsion)


Figure 4: Elementary particles, isomers, ordinary matter, and dark matter. For counterpart elementary particles (across the six isomers), the masses are the same. Here, the word counterpart refers to positions in the six similar arrays of symbols for elementary particles. For counterpart elementary particles, the magnitudes of the charges are the same. For counterpart leptons, the flavours are not necessarily the same. For each known elementary particle, this figure uses popular symbols. Here, the lower-case letter $g$ associates with gluons. Here, the symbol $\gamma$ associates with the photon. The figure de-emphasizes the would-be graviton (or, 2L particle), the might-be 3L particle, and the might-be 4L particle. The figure de-emphasizes components - of the photon (or, 1L particle) for which the reaches exceed one.

| Dark matter and ordinary matter |  |  |
| :---: | :---: | :---: |
| (Relative densities of the universe - dark matter : ordinary matter :: $5^{+}: 1$ ) |  |  |
|  |  | Copyright © 2022 Thomas J. Buckholtz |
| Isomer | DMAI | OMSE |
| 0 | 0.06 DM (IGM is not EM self-interactive) | 1.00 OM (IGM is EM self-interactive) |
| 1 | 0.06 DM (IGM is not EM self-interactive) | 1.00 DM (IGM is not EM self-interactive) |
| 2 | 0.06 DM (IGM is not EM self-interactive) | 1.00 DM (IGM is not EM self-interactive) |
| 3 | 0.06 DM (IGM is not EM self-interactive) | 1.00 DM (IGM might be EM self-interactive) |
| 4 | 0.06 DM (IGM is not EM self-interactive) | 1.00 DM (IGM is not EM self-interactive) |
| 5 | 0.06 DM (IGM is not EM self-interactive) | 1.00 DM (IGM is not EM self-interactive) |

Legend:
Isomer - an isomer of all elementary particles except $\Sigma \mathrm{L}$ elementary particles (or, except [mainly] the photon and the graviton).
DM - dark matter. OM - ordinary matter.
DMAI - each isomer's 0.5 R and 0.5 M elementary particles are bases for stuff that measures as DM. \{AI denotes "in all isomers." $\}$
OMSE - denotes "ordinary-matter similar elementary particles" and includes all elementary particles except $\Sigma \mathrm{L}$ particles, 0.5 R , and 0.5 M
Each of 0.06 and 1.00 is a relative density of the universe. The total relative densities are $\sim 5.38 \mathrm{DM}$ and 1.00 OM .
IGM - Intergalactic medium. EM - electromagnetically. "Might be" associates with - for example - details regarding collisions of galaxy clusters.

Figure 5: Dark matter density of the universe and ordinary matter density of the universe. The DM (or, dark matter) relative densities sum to approximately 5.38 times the OM (or, ordinary matter) relative density. Across isomers, the masses of similar elementary particles are identical. However, for charged leptons, associations between flavour and mass are not necessarily identical. Differences in associations between charged-lepton flavours and charged-lepton masses lead to differences - between isomers - in the evolution of stuff that associates with respective isomers. The stuff that associates with at least four DM isomers of elementary particles evolves so that the associated IGM (or, intergalactic medium) does not interact electromagnetically much with itself, compared to the interactivity of OM IGM. The lack - across at least four DM isomers - of much IGM electromagnetic self-interaction might associate with observations regarding the Bullet Cluster collision of two galaxy clusters.

## Isomers, gravity, eras regarding the universe, and eras regarding galaxies

(Including driving forces that led to eras in the history of the universe and eras regarding galaxy formation)


```
Legend:
DM - Dark matter : OM - Ordinary matter ; ~OM - Mostly ordinary matter
2L - Components of gravity; [ }\mp@subsup{\rho}{1}{}]\mathrm{ - Reach [number of isomers] for one instance of a component ; ( }\cdots\mathrm{ pole) ↔2L multipole expansion
\rightarrow \leftarrow : \text { Attraction; } \rightarrow \text { : Repulsion}
1J - A new elementary boson
```

Figure 6: Isomers of elementary particles, components of gravity, eras in the evolution of the universe, and eras in the evolution of galaxies. Some current galaxies did not transit beyond the first era regarding the evolution of galaxies.

| Eras regarding the rate of separating of large clumps |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ("rate of expansion of the universe") |  |  |  |  |  |  |  |
|  |  |  |  |  |  | Copyrigh | 2022 Thomas J. Buckhotz |
| Rate of separating | Duration | Name | $\leftarrow$ | Initiating force | RDF | Reach | PROP solutions |
| Is negative | ? | NYN | $\swarrow$ | Attractive | $\mathrm{r}^{-6}$ | 6 | 2g1`2`3`6`x |
| Turns positive | ? | NYN | $\checkmark$ | Repulsive | - | 1 | 0g1`3 \({ }^{\prime}\) ` 6 |
| Increases rapidly | Fraction of a second | Inflation | $\checkmark$ | Repulsive | $\mathrm{r}^{-5}$ | 1 | $2 \mathrm{~g} 12^{\prime} 3 \times 4 \mathrm{x}$ |
| Decreases | Billions of years | NYN | $\checkmark$ | Attractive | $\mathrm{r}^{-4}$ | 1 | 2g12`3 |
| Increases | Billions of years | NYN | $\swarrow$ | Repulsive | $\mathrm{r}^{-3}$ | 2 | 2g2 4 |
| Legend: |  |  |  |  |  |  |  |
| The era that associates with a row in the table precedes eras that associate with subsequent rows in the table. NYN - not yet named. |  |  |  |  |  |  |  |
| RDF - radial dependence of force [r denotes the distance between two clumps.]. |  |  |  |  |  |  |  |
| Reach (for one instance of the initiating force) - number of isomers. \{The number of instances is six divided by the reach of one instance.\} |  |  |  |  |  |  |  |
| PROP solution - 2 g .. denotes component(s) of gravity [x denotes more than one.]; 0 g ${ }^{1} 3^{1} 4^{1} 6$ associates with Pauli repulsion (or, with the jay - or 1] - boson). |  |  |  |  |  |  |  |

Figure 7: Suggested and known eras regarding the rate of expansion of the universe. The era that associates with a row in the figure precedes eras that associate with subsequent rows in the figure. For each row, the leftmost three columns describe aspects of the era. The rightmost four columns associate with a noteworthy cause for the era. Generally, the noteworthy cause gains prominence before the era starts. Our work proposes the first two eras to which the image alludes. Other work and our work suggest the era of inflation. Other work and our work model aspects of the two multi-billion-years eras. Our work might explain seeming difficulties that other work seems to exhibit regarding modeling aspects of the current multi-billion-years era of increasing rate of separation.

## Eras regarding the formation of some galaxies

(explanations for DM:OM ratios of one to zero-plus, zero-plus to one, about-four to one, and five-plus to one)
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| Era | Phenomena | $\leftarrow$ | Initiating force | RDF | Reach | PROP solution |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| First | A one-isomer original clump (or, halo) forms | $\checkmark$ | Attractive | $\mathrm{r}^{-4}$ | 1 | $2 \mathrm{~g} 1^{\prime} \mathrm{V}^{\prime} 3$ |
| First | The clump repels stuff associating with one other isomer | $\checkmark$ | Repulsive | $\mathrm{r}^{-3}$ | 2 | $2 \mathrm{~g} 2 \times 4$ |
| Second | The clump attracts stuff that associates with five isomers | $\checkmark$ | Attractive | $\mathrm{r}^{-2}$ | 6 | 2g2 |
| Third | Collisions result in a DM:OM ratio of $5^{+}: 1$ | $\checkmark$ | Attractive | $\mathrm{r}^{-2}$ | 6 | 2g2 |

Legend:
The stage that associates with a row in the table precedes stages that associate with subsequent rows in the table.
Initiating forces might tend to gain prominence before - and dominate for at least early parts of - the respective stages.
RDF - radial dependence of force [r denotes the distance between two significantly massive objects.]. Reach (for one instance of the initiating force) - number of isomers.
Some galaxies do not evolve beyond the first era. Some of these galaxies have essentially only ordinary matter. Some have essentially only dark matter.
During the second stage, the repulsive force component also repels some stuff that associates with the isomer that associates with the original clump.
Based on the repulsive force, some galaxies have nearby stuff that associates essentially with only five isomers.
Some galaxies do not evolve beyond the second era. A DM:OM ratio of $\sim 4: 1$ can pertain. \{DM - dark matter. OM - ordinary matter\}
Collisions can merge stuff from galaxies for which the original-clump isomers differ.
Some galaxies might associate with more than one original clump and with more than one original-clump isomer:

Figure 8: Suggested eras and suggested DM:OM ratios for galaxies. The stage that associates with a row in the figure precedes stages that associate with subsequent rows in the figure. For each row, the leftmost two columns associate with aspects of the stage. The rightmost four columns associate with a noteworthy cause for the stage. The noteworthy cause might gain prominence before the stage starts. Some galaxies do not transit beyond some stages. Our work points to possible propensities for nature to form galaxies with DM:OM ratios of approximately one to zero-plus (that is, dark matter galaxies), five-plus to one, four to one, and zero-plus to one. Galaxies that both had more than one original clump and had three original-clump isomers might tend to cease star formation earlier than do some other galaxies.
of the suggested approximate ratios - one to zeroplus, five-plus to one, four to one, and zero-plus to one.

Figure 9 suggests possible relationships between physics properties. (See table 12 and table 14.) This essay leaves open possible opportunities to use these possible relationships to envision possibly more - than this essay explores - fundamental aspects of nature.

Figure 10 suggests possible rest energies for all known elementary fermions (including neutrinos) and all suggested elementary fermions. (See table 13. table 14 table 15, and table 16, This essay associates each rest energy with direct use of or extrapolation from a few formulas. (See table 14 and table 15. )

## 5. Conclusion

This unit summarizes aspects of our work and suggests perspective about our work.

### 5.1. Our modeling

Our modeling features two bases.
One basis unifies and decomposes aspects of electromagnetism and gravity. For each of electromagnetism and gravity, the decomposition seems to associate well with properties - of objects - that people can measure and that other modeling features.

For electromagnetism, the properties include charge and magnetic moment. For gravity and kinematics related to mass, the properties include mass and moments of inertia.

One basis features isomers of elementary particles that do not mediate long-range interactions and features instances of components of long-range interactions.

Our modeling extends from the two bases to do the following. Match all known elementary particles and suggest possible other elementary particles. Describe dark matter. Point to explanations for data that other modeling seems not to explain. Suggest data that might associate with future observations.

We suggest the possibility that the notion that our work explains phenomena that other modeling does not explain points to usefulness for our work. We explain quantitatively eight quantitative data points or approximate data ranges regarding observed ratios of dark matter effects to ordinary matter effects. (See table 22.) Some other explanations have quantitative bases but - to the extent that this essay uses the explanations - are qualitative. For example, this essay suggests an explanation for dark energy pressure. Presumably, people can use simulations to help verify or refute some of our qualitative explanations. Generally, we know of no cases in which our suggestions that address

## Possible relationships between physics properties <br> (Based on properties of some elementary particles)

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Possible relationship - for elementary bosons - between the properties of mass, spin, and charge

$$
(\text { Mass })^{2}=\left(\left(\text { Mass }_{\mathrm{Higss}}\right)^{2} / 17\right) \times\left(\mathrm{j}^{2}+\mathrm{I}_{\mathrm{ms}}-\mathrm{S}^{2}-\mathrm{Q}(\mathrm{Q}+1)\right)
$$

Possible relationship between electromagnetism and gravity (interrelates one charge and two masses)

```
Define \(\beta\) : \((4 / 3) \times\left(\beta^{2}\right)^{6} \equiv\left(\left(q_{e}\right)^{2} /\left(4 \pi \varepsilon_{0}\right)\right) /\left(\mathrm{G}_{\mathrm{N}}\left(\mathrm{m}_{\mathrm{e}}\right)^{2}\right)\)
Define \(\beta^{\prime}: \beta^{\prime} \equiv \mathrm{m}_{\text {tau }} / \mathrm{m}_{\mathrm{e}}\)
Posit: \(\beta^{\prime}=\beta\)
```

Legend:
$\mathrm{j}=4$ for the Higgs boson $(0 \mathrm{H}), 3$ for the Z boson (1Z), 3 for the W boson $\left(1 \mathrm{~W}_{1}\right)$, and $\mathrm{S}^{2}$ for all other elementary bosons.
$\mathrm{S}=\operatorname{spin}($ in units of $\hbar)$.
$\mathrm{l}_{\mathrm{ms}}=1$ for the Higgs boson ( 0 H ), 1 for the Z boson ( 1 Z ), 1 for the W boson $\left(1 \mathrm{~W}_{1}\right)$, and 0 for all other elementary bosons.
$\mathrm{Q}=\mid$ charge $/$ (charge of the electron)|.
$\left(\left(\mathrm{q}_{\mathrm{e}}\right)^{2} /\left(4 \pi \varepsilon_{0}\right)\right) /\left(\mathrm{G}_{\mathrm{N}}\left(\mathrm{m}_{\mathrm{e}}\right)^{2}\right)$ is the ratio - for two electrons - of electrostatic repulsion to gravitational attraction
$\left(\beta^{\prime}\right)^{2}$ is the ratio of the gravitational attraction between two tau to the gravitational attraction between two similarly (to the two taus) distanced electrons
The exponent of 6 - in the definition for $\beta$ - might associate (for some modeling) with the notion of 6 isomers of non-L-family elementary particles.

Figure 9: Possible relationships between physics properties. The boson-centric relationship might pertain for all known elementary bosons and for all new elementary bosons that we suggest. The possible relationship between electromagnetism and gravity enables computing a tau mass that is compatible with experimental results.

## Rest energies for elementary fermions

|  |  |  |  |  | Copyright © 2022 Thomas J. Buckholtz |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| SФ | Elementary particle | Approximate rest energy |  | Note | SФ | Elementary particle | Approximate rest energy |  | Note |
| 0.5C ${ }_{1}$ | Electron | 0.5109989... | MeV | Exp | 0.5 N | Neutrinos - two or three | $3.4 \times 10^{-2}$ | eV |  |
| $0.5 \mathrm{C}_{1}$ | Muon | 105.658... | MeV | Exp | 0.5 N | Neutrinos - one or zero | $\leq 4.4 \times 10^{-4}$ | eV |  |
| $0.5 \mathrm{C}_{1}$ | Tau | $1776.8400 \pm 0.0115$ | MeV | Calc |  |  |  |  |  |
| $0.5 Q_{2 / 3}$ | Up (quark) |  | MeV |  | 0.5R | Arc | 8.6 to 10.7 | MeV |  |
| $0.5 Q_{1 / 3}$ | Down (quark) |  | MeV |  | 0.5R | Arc | 6.8 to 8.6 | MeV |  |
| $0.5 Q_{1 / 3}$ | Charm (quark) | $1.27 \times 10^{3}$ | MeV |  | 0.5R | Arc | 102 to 106 | MeV |  |
| $0.5 \mathrm{Q}_{2 / 3}$ | Strange (quark) | $9.3 \times 10^{1}$ | MeV |  | 0.5 M | Heavy neutrino | $\geq 6 \times 10^{3}$ | GeV |  |
| $0.5 \mathrm{Q}_{2 / 3}$ | Top (quark) | $1.71 \times 10^{5}$ | MeV |  | 0.5 M | Heavy neutrino | $\geq 6 \times 10^{3}$ | GeV |  |
| $0.5 \mathrm{Q}_{1 / 3}$ | Bottom (quark) | $4.18 \times 10^{3}$ | MeV |  | 0.5M | Heavy neutrino | $\geq 6 \times 10^{3}$ | GeV |  |
| Legend: |  |  |  |  |  |  |  |  |  |
| Known particle or value <br> Suggested particle or value |  | Exp - Result from experiments. |  |  |  |  |  |  |  |
|  |  | Calc - Result from a calculation: |  |  | The standard deviation reflects the standard deviation of measurements of $G_{N}$. |  |  |  |  |

Figure 10: Possible rest energies for elementary fermions. Eight standard deviations of the calculated tau rest energy fit within one standard deviation of the measured tau rest energy. The calculated quark rest energies comport with experimental results. Results regarding neutrinos comport with the notion of three (not heavy) neutrinos and with astrophysics data. Our work suggests possible ranges for the rest energies of each of the three possible arc (or, 0.5 R ) elementary fermions. Our work suggests two possible lower limits for the rest energies of the three possible heavy neutrino elementary fermions. The figure shows one of those possible lower limits. The other possible lower limit is $\sim 2.5 \times 10^{9} \mathrm{GeV}$.
possible gaps between other modeling and observations point - compared to other modeling - in a wrong direction regarding closing gaps.

We suggest the possibility that the notion that our work suggests specifications and data that other modeling does not suggest points to possible usefulness for our work. Our suggestions include a specification for dark matter, specifications for new elementary particles, and more (than current measurements provide) accurate masses for neutrinos and some other known elementary particles.

We suggest that the small set of bases for our modeling, the breadth of seemingly coherent scope of our modeling, the simplicity of relevant Diophantine equations, and the possible ease of integrating our modeling and other modeling point to possible usefulness for our work.

### 5.2. Our work

Our work suggests augmentations - to physics modeling - that produce results that may provide progress regarding the following physics opportunities. Complete the list of elementary particles. Describe dark matter. Explain ratios of dark matter to ordinary matter. Explain eras in the history of the universe. Link properties of objects. Interrelate physics models.

We use our modeling to match data that other modeling matches.

We use our modeling to suggest explanations for data that other modeling seems not to explain.

We use our modeling to suggest results regarding data that people have yet to gather.

The breadth and depth of the matched data might suffice to justify using our modeling.

The breadth and unity - within itself and with physics modeling that people use successfully - of our modeling might support the usefulness of our modeling.

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[^0]:    Email address: Thomas.Buckholtz@RoninInstitute.org (Thomas J. Buckholtz)

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