# Schrödinger Equation and Free Particle Wave Function

#### Karl De Paepe

#### Abstract

Using the wave function of a free particle we obtain a solution of the Schrödinger equation for a class of potentials.

### **1** Time dependent accelerating frame of reference

Consider an accelerating frame of reference  $\mathcal{F}'$  with coordinates x', t' and an inertial frame of reference  $\mathcal{F}$  with coordinates x, t. The coordinates of the frames being related by

$$x' = x - f(t) \qquad t' = t \tag{1}$$

Since dx' = dx and position probabilities are the same for  $\mathcal{F}'$  and  $\mathcal{F}$  we have for the wave function  $\psi(x, t)$  with respect to  $\mathcal{F}$  and corresponding wave function  $\psi'(x', t')$  with respect to  $\mathcal{F}'$  that [1]

$$|\psi'(x',t')|^2 = |\psi(x,t)|^2 \tag{2}$$

Consequently there is a real valued function  $\beta(x, t)$  such that

$$\psi'(x',t') = e^{-\frac{i}{\hbar}\beta(x,t)}\psi(x,t) \tag{3}$$

With respect to  $\mathcal{F}$  let the wave function  $\psi(x, t)$  satisfies the Schrödinger equation

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2}(x,t) = i\hbar\frac{\partial\psi}{\partial t}(x,t)$$
(4)

With respect to  $\mathcal{F}'$  we have an additional force  $m\ddot{f}(t)$  and hence an additional potential  $m\ddot{f}(t)x' + V_0(t')$ . The wave function  $\psi'(x',t')$  then satisfies the Schrödinger equation

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi'}{\partial x'^2}(x',t') + \left(m\ddot{f}(t')x' + V_0(t')\right)\psi'(x',t') = i\hbar\frac{\partial\psi'}{\partial t'}(x',t')$$
(5)

Now

$$\frac{\partial}{\partial x'} = \frac{\partial}{\partial x} \qquad \frac{\partial}{\partial t'} = \dot{f} \frac{\partial}{\partial x} + \frac{\partial}{\partial t}$$
(6)

and on substituting (3) in (5) and using (4) and (6) gives

$$\begin{bmatrix} \frac{i\hbar}{2m} \frac{\partial^2 \beta}{\partial x'^2} + \frac{1}{2m} \left( \frac{\partial \beta}{\partial x'} \right)^2 + m\ddot{f}(x-f) + V_0 - \dot{f} \frac{\partial \beta}{\partial x} - \frac{\partial \beta}{\partial t} \end{bmatrix} \psi + \frac{i\hbar}{m} \left[ \frac{\partial \beta}{\partial x} - m\dot{f} \right] \frac{\partial \psi}{\partial x} = 0$$
(7)

We have

$$\beta(x,t) = m\dot{f}(t)x + \int_0^t [V_0(s) - mf(s)\ddot{f}((s) - \frac{1}{2}m\dot{f}(s)^2]ds + C$$
(8)

is the unique solution of (7) satisfying the initial condition [2]

$$\beta(x,0) = m\dot{f}(0)x + C \tag{9}$$

### 2 Space and time dependent velocity

Let  $v_{\epsilon}(x,t)$  be a smooth function in variables  $\epsilon, x, t$ . Require  $v_{\epsilon}(x,0) = 0$ . Define  $X_{\epsilon}(u;t)$  to be the curve x(t) such that

$$\frac{dx}{dt} = v_{\epsilon}(x, t) \tag{10}$$

and x(0) = u. Require that the curves are defined for all t and the curves do not intersect. We then have a frame of reference  $\mathcal{F}_{\epsilon}$  with coordinates  $x_{\epsilon}, t_{\epsilon}$  such that

$$x_{\epsilon} = X_{\epsilon}(x;t) \qquad t_{\epsilon} = t \tag{11}$$

Let  $\psi(x,t)$  satisfy (4). Let  $V_{\epsilon}(x_{\epsilon},t_{\epsilon})$  be the potential in these coordinates. We have

$$\frac{1}{m}\frac{\partial V_{\epsilon}}{\partial x_{\epsilon}}(x_{\epsilon}, t_{\epsilon}) = v(x, t)\frac{\partial v}{\partial x}(x, t) + \frac{\partial v}{\partial t}(x, t)$$
(12)

Let  $\psi_{\epsilon}(x_{\epsilon}, t_{\epsilon})$  be the wave function satisfying the Schrödinger equation in  $x_{\epsilon}, t_{\epsilon}$  coordinates and  $\psi_{\epsilon}(x, 0) = \psi(x, 0)$ . Let  $B(x_0; \epsilon)$  be the set of points  $x_0 - \epsilon < x < x_0 + \epsilon$ . Choose  $v_{\epsilon}(x, t)$  so that for  $u \in B(x_0; \epsilon)$ 

$$X_{\epsilon}(u;t) = X_0(x_0;t) + u - x_0 \tag{13}$$

Let  $\widehat{\mathcal{F}}$  be a frame of reference with coordinates  $\hat{x}, \hat{t}$  related to coordinates x, t of  $\mathcal{F}$  by

$$\hat{x} = x - X_0(x_0; t) \qquad \hat{t} = t$$
 (14)

The potential in these coordinates is  $m\ddot{X}_0(\hat{x}_0:\hat{t})\hat{x} + V_0(\hat{t})$ . Let  $\hat{\psi}(\hat{x},\hat{t})$  be the wave function satisfying the Schrödinger equation with this potential and  $\hat{\psi}(x,0) = \psi(x,0)$ . We have by (8) a  $\hat{\beta}(x,t)$  such that

$$\frac{\widehat{\psi}(\widehat{x},\widehat{t})}{\psi(x,t)} = e^{-\frac{i}{\hbar}\widehat{\beta}(x,t)} \qquad \frac{\partial\widehat{\beta}}{\partial x}(x,t) = m\dot{X}_0(x_0;t) \tag{15}$$

hence for points  $(X_{\epsilon}(u;t),t)$  where  $u \in B(x_0;\epsilon)$  we have

$$\frac{\psi_{\epsilon}(x_{\epsilon}, t_{\epsilon})}{\psi(x, t)} = e^{-\frac{i}{\hbar}\widehat{\beta}(x, t)} \qquad \frac{\partial\widehat{\beta}}{\partial x}(x, t) = m\dot{X}_{0}(x_{0}; t)$$
(16)

Define coordinates  $x' = x_0, t' = t_0$ . Let  $\psi'(x', t') = \psi_0(x', t')$ . Now  $x_0$  is arbitrary and let  $\beta(x, t)$  be the limit of  $\hat{\beta}(x, t)$  as  $\epsilon \to 0$  so we get

$$\frac{\partial\beta}{\partial x}(x,t) = mv_0(x,t) \tag{17}$$

Require  $v(x,t) \to 0$  as  $v \to -\infty$ . We then have  $\beta(x,t) \to 0$  as  $x \to -\infty$  hence by (17)

$$\beta(x,t) = \int_{-\infty}^{x} v_0(u,t) du \tag{18}$$

Consequently

$$\psi'(x',t') = e^{-\frac{i}{\hbar} \int_{-\infty}^{x} v_0(u,t) du} \psi(x,t)$$
(19)

## References

- [1] K. De Paepe, Physics Essays, September 2008
- [2] A. Colcelli, G. Mussardo, G. Sierra, A. Trombettoni, Phys. Rev. Lett. 123, 130401 (2019)