# The strong Bell's theorem which does not use the assumption of locality 

Jiř̌́ Souček<br>Charles University in Prague, Faculty of Arts Nam. Jana Palacha 1/2, 11638, Prague, Czech Republic jiri.soucek@ff.cuni.cz


#### Abstract

The standard Bell's theorem states that quantum mechanics (QM) cannot be locally realistic. Here we prove the strong Bell's theorem which states that QM cannot be realistic. In the proof we use the Mermin's form of the Bell's theorem. Our result solves the old dilemma: non-locality or non-realism. Then we discuss the consequences of the strong Bell's theorem, e.g. that no Bell non-locality of QM exists and that quantum theory is local and non-realistic.


Keywords: Bell's theorem, non-locality, non-realism

## 1. Introduction.

The standard form of Bell's theorem (see [3]) has the following form:
The assumption of local realism of quantum mechanics (QM) implies Bell's inequality which is in the contradiction with experimental facts. As a consequence of this result QM is not local realistic, i.e. either QM is non-local or QM is non-realistic.

We interpret the realism as a counter-factuality: any observable has a value even if it is not measured. (The standard QM rule requires that only the measured observables have a value.)

In this paper we prove the strong form of the Bell's theorem: The realism of QM implies the Bell's theorem and thus the contradiction with the experiment. The conclusion is that QM must be non-realistic.

The difference between the previous assumption (the local realism of QM) and the new assumption (the realism of QM) is enormous. For many years (and tens of years) QM was imprisoned in this dilemma:

## Non-locality or non-realism?

If the non-realism of QM is proved, the problem of non-locality of QM would become irrelevant, i.e. Bell non-locality should not exist.

The proof of non-realism, that we present here, is based on the modification of the Mermin's form of the Bell's theorem (see [1]) and on the simplifications from [2].

The main point is the following: the first part of the proof which usually uses the locality assumption is substituted by the proof based completely on the experimental facts.

The paper is organized as follows. In the second section we define basic concepts. In the third section we present the proof of the strong Bell's theorem. Finally we conclude by revising our arguments and viewing them in a broader context.

## 2. Simple observables, the EPR entangled state

Let us consider the quantum system $S$ with the (finite dimensional) Hilbert space of states $\mathbf{H ( S )}$ and its observable A. We shall call A a simple observable if its eigenvalues are $\{1 .-1\}$. Simple observables are in the $1-1$ relation with projectors given by formulas $P=1 / 2(A+I), A$ $=2 \mathrm{P}-\mathrm{I}$ (where I is the identity operator). Thus simple observables can generate general observables using the spectral decomposition.

The assumption that a simple observable $A$ is realistic is equivalent to the so-called counterfactuality: A is assumed to have a value also in the situation where A is not measured.

The EPR system is the composition of two 2-dimensional systems $S_{1}$ and $S_{2}$ in a special "entangled" state. Let us consider two 2-dimensional systems $S_{1}$ and $S_{2}$ with a simple observable $A_{1}$ in $\mathbf{H}\left(\mathrm{S}_{1}\right)$ and a simple observable $\mathrm{A}_{2}$ in $\mathbf{H}\left(\mathrm{S}_{2}\right)$. There exists orthogonal basis $\left\{u_{+}{ }^{i}, u_{-}^{i}\right\}, i=1,2$ such that $A_{i}\left(u_{+}{ }^{i}\right)=1, A_{i}\left(u_{-}{ }^{i}\right)=-1, i=1,2$.

Then the entangled state ([4]) is defined as

$$
\Psi=2^{-1 / 2}\left(u_{+}{ }^{1} \otimes u_{-}{ }^{2}-u_{-}^{1} \otimes u_{+}{ }^{2}\right)
$$

where the state $\Psi$ is an element of the space $\mathbf{H}\left(\mathrm{S}_{1}\right) \otimes \mathbf{H}\left(\mathrm{S}_{2}\right)$ and $\otimes$ denotes the tensor product.

It is clear that systems $S_{1}$ and $S_{2}$ are in some sense correlated (in fact, anti-correlated). It is assumed that systems $S_{1}$ and $S_{2}$ are then sent in opposite directions and then measured by observers Alice and Bob (further developed in part 3).

## 3. The strong Bell's theorem.

The main goal of this central section is to prove the strong Bell's theorem solely with the assumption of realism. (The usual way is to prove Bell's theorem using the assumption of the local realism.)

The detailed construction of our proof is based on the Mermin's form of the Bell's theorem ([1]).

We shall consider the experiment described in [2], where Alice and Bob are used as independent observers: measurements at Alice in possible three orientations $\mathrm{A}_{1}, \mathrm{~B}_{1}$ and $\mathrm{C}_{1}$ and the measurements in possible three orientations $\mathrm{A}_{2}, \mathrm{~B}_{2}$ and $\mathrm{C}_{2}$ at the part of Bob. The orientations are taken at angles $0^{\circ}, 120^{\circ}$ and $240^{\circ}$. The choice of the orientation at Alice' part (among $\mathrm{A}_{1}, \mathrm{~B}_{1}, \mathrm{C}_{1}$ ) and the choice of the orientation at the Bob's part (among $\mathrm{A}_{2}, \mathrm{~B}_{2}, \mathrm{C}_{2}$ ) are chosen randomly.

Orientations of the Alice's part and of the Bob's part will be correlated. This means that the orientation of $A_{1}$ is the same as the orientation of $A_{2}$ and similarly for $B_{1}$ and $B_{2}$ and for $C_{1}$ and $\mathrm{C}_{2}$.

We shall consider also the combinations of orientations of $\mathrm{A}_{1}-\mathrm{B}_{2}, \mathrm{~A}_{1}-\mathrm{C}_{2}, \mathrm{C}_{1}-\mathrm{B}_{2}$. All observables $\mathrm{A}_{1}, \mathrm{~B}_{2}, \mathrm{C}_{1}, \mathrm{C}_{2}$ are simple observables. For the detail description see [2].

We shall assume that the realism of QM holds, i.e. that observables $\mathrm{A}_{1}, \mathrm{~B}_{2}, \mathrm{C}_{1}, \mathrm{C}_{2}$ has the "pre-existing" (i.e. counter-factual) values.

It is possible to formulate the basic properties which have been tested experimentally and which means that the following facts are experimentally proved:

Experimental fact 1 (EF1). The measurements of $A_{1}-A_{2}$ and correspondingly $B_{1}-B_{2}$ and $C_{1}-$ $C_{2}$ gives the $100 \%$ anti-correlation, i.e. that relations $A_{1}=-A_{2}, B_{1}=-B_{2}, C_{1}=-C_{2}$ are true in all rounds.

Experimental fact 2 (EF2). The correlations between $A_{1}$ and $B_{2}, A_{1}$ and $C_{2}$ and $C_{1}$ and $B_{2}$ can be measured. It is obtained that the events $\left[A_{1}=B_{2}\right],\left[A_{1}=C_{2}\right]$ and $\left[C_{1}=B_{2}\right]$ have probabilities approx. 0.75 , i.e. at least 0.7 (in consideration of possible experimental errors). (In fact, the probabilities of events $\left[\mathrm{A}_{1}=-\mathrm{B}_{2}\right],\left[\mathrm{A}_{1}=-\mathrm{C}_{2}\right]$ and $\left[\mathrm{C}_{1}=-\mathrm{B}_{2}\right]$ are measured to be approx. 0.25.). The details of the discussion can be found in [2].)

Theorem 3.1. QM is not realistic (i.e. the counter-factual reasoning cannot be used).

## Proof.

We shall assume that QM is realistic. This means that each simple observable has its preexisting value.

## The first part of the proof.

Using the standard probabilistic formula
$\operatorname{prob}(\mathrm{X} \cup Y)=\operatorname{prob}(\mathrm{X})+\operatorname{prob}(\mathrm{Y})-\operatorname{prob}(\mathrm{X} \cap Y)$, i.e.

$$
\operatorname{prob}(\mathrm{X} \cap \mathrm{Y})=\operatorname{prob}(\mathrm{X})+\operatorname{prob}(\mathrm{Y})-\operatorname{prob}(\mathrm{X} \cup \mathrm{Y})
$$

we obtain that

$$
\operatorname{prob}(\mathrm{X} \cap \mathrm{Y}) \geqq \operatorname{prob}(\mathrm{X})+\operatorname{prob}(\mathrm{Y})-1
$$

Denote $\mathrm{X}=\left[\mathrm{A}_{1}=\mathrm{B}_{2}\right], \mathrm{Y}=\left[\mathrm{A}_{1}=\mathrm{C}_{2}\right], \mathrm{Z}=\left[\mathrm{B}_{1}=\mathrm{C}_{2}\right]=\left[\mathrm{C}_{1}=\mathrm{B}_{2}\right]$ (here we have used EF1). From EF2 we know that probabilities of $\mathrm{X}, \mathrm{Y}, \mathrm{Z}$ are at least 0.7.

We obtain

$$
\operatorname{prob}(\mathrm{X} \cap \mathrm{Y}) \geqq \operatorname{prob}(\mathrm{X})+\operatorname{prob}(\mathrm{Y})-1 \geqq 0.7+0.7-1=0.4
$$

Then we obtain

$$
\operatorname{prob}(\mathrm{X} \cap \mathrm{Y} \cap \mathrm{Z}) \geqq \operatorname{prob}(\mathrm{X} \cap \mathrm{Y})+\operatorname{prob}(\mathrm{Z})-1 \geqq 0.4+0.7-1=0.1>0
$$

If the experiment is repeated the sufficient number of times then the round where $\mathrm{X} \cap \mathrm{Y} \cap \mathrm{Z}$ happens must exist.

The conclusion: there exists a round such that $\left[\mathrm{A}_{1}=\mathrm{B}_{2}\right],\left[\mathrm{A}_{1}=\mathrm{C}_{2}\right],\left[\mathrm{C}_{1}=\mathrm{B}_{2}\right]$ happen.
The second part of the proof
We shall obtain the contradiction. We shall consider the round where

$$
\mathrm{A}_{1}=\mathrm{B}_{2}, \mathrm{~A}_{1}=\mathrm{C}_{2}, \mathrm{C}_{1}=\mathrm{B}_{2}
$$

which existence was proved in the first part.
If $\mathrm{A}_{1}=1$, then $\mathrm{B}_{2}=1, \mathrm{C}_{2}=1, \mathrm{C}_{1}=1$ and then $\mathrm{C}_{1}=\mathrm{C}_{2}$ is in the contradiction with EF1.
If $A_{1}=-1$, then $B_{2}=-1, C_{2}=-1, C_{1}=-1$ and then $C_{1}=C_{2}$ is in the contradiction with EF1.
This contradiction implies that the assumption of the realism of QM is false. Therefore, we were able to prove that the assumption of realism does not hold. The end of the proof.

We see that the methods, we used for our proof, did not require locality of QM. The first part of the proof was based on the experimental facts (and the elementary probability theory) while the second part was based on the theory of logic. In both parts the realism of QM (i.e. counter-factual reasoning) was the only assumption used.

## 4. Conclusions

We have proved the strong version of Bell's theorem where only the realism of QM (instead of the realism + locality of QM ) is assumed. As a consequence, we obtained the proof of the non-realism of QM.

Some consequences and related considerations of the strong Bell's theorem were discussed, also in [8]. It was shown that no Bell non-locality exists.

We are convinced that the proof of the strong Bell's theorem will change completely the field of Bell's inequalities. Many of the assumptions that were previously used, would need to be fundamentally reconsidered. This reconsideration will be the theme of the future research.

## References

[1] N. David Mermin: "Is the moon there when nobody looks? Reality and the quantum theory" Physics Today, April 1985, pag. 38-47
[2] D.R.Schneider, Bell`s theorem with easy math, http://drchinese.com/David/Bell Theorem Easy Math.htm
[3] J.S. Bell: "On the Einstein Podolsky Rosen paradox" Physics 1 \#3, 195 (1964).
[4] A. Einstein, B. Podolsky, N. Rosen: "Can quantum-mechanical description of physical reality be considered complete?" Physical Review 41, 777 (15 May 1935).
[5] Jiri Soucek, A new model for quantum mechanics and the invalidity of no-go theorems, ISBN 978-613-7-33070-8, LAP LAMBERT Academic Publishing 2018 ; available also at: https://www.academia.edu/36370035/A new model for quantum mechanics and the in validity of no-go theorems
[6] J. Soucek, The principle of anti-superposition in QM and the local solution of the Bell's inequality problem, http://www.nusl.cz/ntk/nusl-177617
[7] J. Soucek, Extended probability theory and quantum mechanics I: non-classical events, partitions, contexts, quadratic probability spaces, arXiv:1008.0295v2[quant-ph]
[8] J. Soucek, The consequences of the non-realism of quantum theory: empty proofs and empty statements,
https://www.academia.edu/42723421/The_consequences_of_the_non_realism_of_quantu m_theory_empty_proofs_and_empty_statement

