

# **Yukawa Potential and Extended Klein-Gordon Equation in Rindler Space-Time**

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## **ABSTRACT**

Yukawa potential satisfy Proka equation or Klein-Gordon equation. If we represent Yukawa potential in Rindler space-time, this Yukawa potential satisfy the extended Klein-Gordon equation in Rindler space-time.

**PACS Number:03.30.+p,03.65**

**Key words: Yukawa potential;**

**Rindler Space-time;**

**Extended Klein-Gordon equation**

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## 1. Introduction

Atom's nucleus force understand by Yukawa potential. We study Yukawa potential in Rindler Space-time.

At first, Yukawa potential  $V$  describes nucleus's combine force in semi-classical method.

$$V = -\frac{g^2}{r} \exp\left(-\frac{m_\pi r c}{\hbar}\right)$$

$$g \text{ is real number, } m_\pi \text{ is the meson's mass} \quad (1)$$

Klein-Gordon equation is satisfied by Yukawa potential  $V$ .

$$-\partial_i \partial^i V + \frac{m^2 c^2}{\hbar^2} V = -\nabla^2 V + \frac{m_\pi^2 c^2}{\hbar^2} V = 0$$

$$V = -\frac{g^2}{r} \exp\left(-\frac{m_\pi r c}{\hbar}\right), j = 1, 2, 3 \quad (2)$$

## 2. Yukawa potential from Extended Klein-Gordon Equation in Rindler-Space-Time

Rindler coordinates are

$$ct = \left(\frac{c^2}{a_0} + \xi^1\right) \sinh\left(\frac{a_0}{c} \xi^0\right), \quad x = \left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}$$

$$y = \xi^2, z = \xi^3 \quad (3)$$

If we write Yukawa potential  $V$  in inertial frame,

$$V = \frac{g^2}{r} \exp\left(\frac{m_\pi r c}{\hbar}\right) \quad (4)$$

If we rewrite Yukawa potential  $V_\xi$  in Rindler space-time,

$$V_\xi = \frac{g^2}{\sqrt{x^2 + y^2 + z^2}} \exp\left(\frac{m_\pi c}{\hbar} \sqrt{x^2 + y^2 + z^2}\right)$$

$$= -\frac{g^2}{\sqrt{\left\{\left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}\right\}^2 + (\xi^2)^2 + (\xi^3)^2}} \exp\left[-\frac{m_\pi c}{\hbar} \sqrt{\left\{\left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}\right\}^2 + (\xi^2)^2 + (\xi^3)^2}\right] \quad (5)$$

This Yukawa potential satisfy the extended Klein-Gordon equation. At first, energy and momentum are in Rindler space-time,

$$E_\xi = i\hbar \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)} \frac{1}{c} \frac{\partial}{\partial \xi^0}, \quad \vec{p}_\xi = i\hbar \vec{\nabla}_\xi \quad (6)$$

Energy-Momentum equation is in Rindler space-time[1],

$$E_{\xi}^2 = \vec{p}_{\xi} \cdot \vec{p}_{\xi} + m^2 c^2 \quad (7)$$

Hence, normal Klein-Gordon equation is in Rindler-spacetime,

$$\begin{aligned} & \frac{m^2 c^2}{\hbar^2} V_{\xi} + \frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_{\xi} - \nabla_{\xi}^2 V_{\xi} \\ &= \frac{m^2 c^2}{\hbar^2} V_{\xi} + \frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_{\xi} - \nabla_{\xi}^2 V_{\xi} \quad (8) \end{aligned}$$

In this time, we focus the gauge  $\Lambda$  equation in Rindler space-time[1],

$$\frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \frac{\partial^2}{(\partial \xi^0)^2} \Lambda - \nabla_{\xi}^2 \Lambda - \frac{\partial \Lambda}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)} = 0 \quad (9)$$

Hence, Eq(8) change extended Klein-Gordon equation in Rindler space-time.

Extended Klein-Gordon Equation is in Rindler space-time,

$$\begin{aligned} & \frac{m^2 c^2}{\hbar^2} V_{\xi} + \frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_{\xi} - \nabla_{\xi}^2 V_{\xi} - \frac{\partial V_{\xi}}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)} \\ &= \frac{m^2 c^2}{\hbar^2} V_{\xi} + \frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_{\xi} - \nabla_{\xi}^2 V_{\xi} - \frac{\partial V_{\xi}}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)} = 0 \quad (10) \end{aligned}$$

Eq(5), Yukawa potential  $V_{\xi}$  satisfy Eq(10), extended Klein-Gordon equation in Rindler space-time.

### 3. Conclusion

We found Yukawa potential mechanism in Rindler Space-time..

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