Electro-gravity via Geometric Chronon Field and on the Origin of Mass

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Abstract: In 1982, Dr. Sam Vaknin pondered the idea of reconstructing physics based on time as a field. His idea appeared in his doctorate dissertation as an amendment to the Dirac spinor equation. Sam saw the Quantum Field Theory particles and momentum and energy as a result of the language of physics and of the way the human mind perceives reality and not as reality. To the author’s opinion, it is a revolution of the language itself and is not a new interpretation of the existing language. The Special Theory of Relativity was a revolution and so was the General Theory of Relativity but yet these theories did not challenge the use of momentum and energy but rather gave them new relativistic interpretation. Later on, Quantum Mechanics used Energy and Momentum operators and even Dirac’s orthogonal matrices are multiplied by such operators. Quantum Field Theory assumes the existence of particles which are very intuitive and agree with the human visual system. Particles may be merely a human interpretation of events that occur in the human sensory world. This paper elaborates on one specific interpretation of Sam Vaknin’s idea that the author has developed from 2003 up to August 2018. It is a major improvement of previously published papers and it summarizes all of them and includes all the appendices along with new ideas. A key idea in this paper is, that while a preferable coordinate of time violates the principle of general relativity, a scalar field does not, because it does not point to any preferable direction in space time, moreover, such a scalar field need not be unique.

Keywords: Time; General Relativity; Foliations; Reeb class; Godbillon-Vey class; Symplectic Geometry; Muon; Higgs; Fine Structure Constant
Introduction

This paper’s approach to the description of matter is geometrical rather than algebraic. In that aspect, it is more loyal to General Relativity than to Super Algebraic approaches [1] that extend the dimensionality of space-time based on algebras such as Grassman’s Algebra. In that aspect, it presents a much simpler view of what matter is. To better understand the paper, the reader is recommended to be familiar with Geroch’s Splitting Theorem [2], although the Morse functions in this paper are more general as complex functions along integral curves that need not be geodesic or unique and need not always be time-like as they can have critical values.

This paper summarizes two previous papers [3], [4] and contains new results. For background, the reader is referred to the work of Zvi Scarr and Yaakov Friedman [5] as a recommended preliminary material, though uniform covariant acceleration is not identical to a model of force fields as curvature of trajectories of different types of clocks and of gravity, as a response of space-time - by curving itself as an error correction mechanism.

On a Big Bang space-time manifold, it is reasonable to define a Morse function [6] which is a submersion [7] from space-time to the real numbers. The idea is that every event can be connected by an integral curve to an S3 manifold of an infinitely small radius on which clocks can be synchronized and then from all such curves, the Morse function is defined as the length of the maximal proper time curve or curves from the sub-mantifold to the event. A global Morse function as time may not be measured along a single path and if it represents measurement by material clocks, may not be always geodesic because material clocks interact. This approach, however, in some geometries such as the big bang metric offers global time though no global time coordinate and this idea is not new [8]. Cosmologists will argue that such a definition is problematic because there is a difficulty in defining a limit backwards in time to a sub-mantifold on which clocks can be synchronized. To those readers, the author says that also on De-Sitter on open slicing and on Anti De Sitter manifolds, such a Morse function can be defined, although it may not be unique and as we shall see, uniqueness is indeed not required. The definition of a Morse function in space-time, does not even require synchronization of clocks on a sub-mantifold. All it requires is to solve some minimum action integral with several symmetries. The longest proper time curve from each event to a three dimensional sub-manifold is valid as long as the resulting submersion function is a Morse function, i.e. its singularities are non - degenerate. Such a sub-manifold can be represented as a leaf of a foliation [9] of space-time where the Morse function is zero. The classification of non degenerate singularities can be found in the Morse lemma [10]. How can we describe time as a Morse function in order to account for matter ? What is matter ? We try to reconstruct physics bottom-up from its very foundations. We borrow from the old language of physics the idea of time and Minkowsky space-time manifold, though it may be arguable that even space can be deduced from time.

The 1982 argument of Dr. Sam Vaknin was that as observers, we can imagine being out of space but that even language requires time and therefore, time should account for matter and particles [11] as there is no physical phenomenon out of time.

Matter is characterized by force fields. It is a phenomenological approach. In geometric terminology, forces bend trajectories of clocks because a 4-acceleration of a physical clock is perpendicular to it’s 4-velocity. We can say that forces prohibit geodesic motion of material clocks in space-time. Contrary to forces, according to the General Theory of Relativity, motion in a gravitational field is along geodesic curves but in curved space-time. Our objective will be to reach an equation that combines these two types of motion, non-geodetic and geodetic in curved space-time. The non – geodetic motion is not anticipated by the metric alone, which means that we will need an additional structure beside the metric tensor, when we borrow the language of the General Theory of Relativity. This paper is not mathematically difficult to understand but it does offer new perception challenges. Following are important points that the author asks the reader to pay attention to. These are very important points:

1) A foliation which is defined by sub-manifolds perpendicular to a vector field is properly defined as covariant because the property of perpendicular vectors does not depend on choice of coordinates. Foliations can be defined in a non-covariant way but this is not the case in this paper.
2) Time-like curves that describe a physical observable, may not be unique. Such a representation is inherently prone to symmetries.

3) By discussing a “field of time” to account for matter, the idea is to capture non–geodesic motion which is not anticipated by the metric alone and to explain the origin of forces not by the traditional quantum particles exchange or by the classical potentials or vector potentials but by a field that causes non geodesic motion of material clocks that measures time. In another interpretation of such a field, we can say that the extra geometrical information that is needed to represent forces, and thus matter, is stored in the geometry of certain foliations of space time, though due to Lagrangian symmetry, these foliations and the field are not unique.

Throughout this paper $x^\mu$ denotes the contravariant coordinate system with index $\mu$. The comma denotes ordinary derivative, $P_\mu \equiv \frac{dp}{dx^\mu}$, which will often be abbreviated, $P_\mu \equiv P_\mu^\mu$. Likewise, semicolon denotes the ordinary covariant derivative that uses the Christoffel symbols $P_\mu ; \nu \equiv \frac{dp}{dx^\mu} - P_\rho \Gamma_\mu^\rho_\nu$. $g_{\mu \nu}$ is the metric tensor. $g$ is its determinant and $\sqrt{-g}$ is the volume coefficient in integration and the volume element is $\sqrt{-g} d\Omega$. $d\Omega$ can be also written as $d\Omega = dx^0 dx^1 dx^2 dx^3$ in Cartesian coordinates.

The Reeb field and its electromagnetic interpretation

Describing a trajectory of a clock, we can write $\frac{dv^\mu}{d\tau} = a^\mu$ and $\frac{dv^\mu}{d\tau} = 0$. The latter is a result of

$$V^\mu V_\mu = c^2 \Rightarrow 2 V_\mu \frac{dv^\mu}{d\tau} = 2 V_\mu a^\mu = \frac{dc^2}{d\tau} = 0 \quad (1)$$

where $c$ is the speed of light, $V^\mu$ is velocity and in the Special Theory of Relativity it is simply

$$\vec{V} = \left( \frac{c V x, V y, V z}{\sqrt{1 - V^2/c^2}} \right) \quad (2)$$

In coordinate system $(t, x, y, z)$. $\tau$ measures proper time.

We can normalize the velocity 4-vector and get $\vec{V} = \frac{1}{c} \left( V x, V y, V z \right)$. A scaled acceleration $\frac{a^\mu}{c^2}$ measures the acceleration of this unit vector $\frac{v^\mu}{c}$ in relation to an arc length $ct$ so $d\tau = cd\tau$ and we can write

$$\frac{d\vec{v}^\mu}{cd\tau} = \frac{a^\mu}{c^2} \quad (3)$$

Denoting our Morse function $P$, can we derive a vector which will be equivalent to an acceleration of a clock that moves along the integral curves, generated by $P_\mu \equiv \frac{dp}{dx^\mu}$? In this case, $P_\mu P^\mu$ is generally not constant. For example, if it is constant and positive, we may choose a monotonically increasing analytic function of $P$ instead of $P$ such as
\[ \sqrt{P} \text{ and still have a Morse function. To see a non-vanishing } (P^\mu P^\mu) \text{ the reader can refer to Appendix E, see term (80).} \]

We now define the Reeb vector (Reeb, 1948, 1952 [12]) of \( P_\mu \), and we will develop the Reeb vector \( \eta \) as it was originally developed in 1948 in the language of De Rham Cohomology, with adjustment to be derived from the normalized vector.

\[ \omega = \frac{P_\mu}{\sqrt{Z}} dx^\mu, \quad D \omega = \eta^\wedge \omega \]  

(4)

and for the sake of brevity, we write \( Z \equiv P_\lambda P^\lambda \) or if \( P \) is a complex scalar field, \( Z \equiv \sqrt{|P_\lambda P^\lambda + P^*_\lambda P^\lambda|} \) and \( Z_\mu \equiv \frac{dz}{dx^\mu} \).

Note: \( \eta^\wedge D\eta \) is the famous Godbillon - Vey cohomology equivalence class [13].

If we limit the discussion to a real scalar function \( P \), the form \( \eta \) can be easily calculated as \( \eta = \frac{U_\mu}{2} dx^\mu \) such that

\[ U_\mu = \frac{Z_\mu}{Z} - \frac{Z_k P^k}{Z^2} P_\mu \implies \]

\[ \frac{d}{dx^\nu} \frac{P_\mu}{\sqrt{Z}} = \frac{d}{dx^\mu} \frac{P_\nu}{\sqrt{Z}} = \]

\[ \frac{P_{\mu\nu}}{\sqrt{Z}} - \frac{P_\mu Z_\nu}{2Z^2} - \frac{P_\nu Z_\mu}{2Z^2} + \frac{P_\nu Z_\mu}{2Z^2} = \]

\[ \frac{P_\nu Z_\mu}{2Z^2} - \frac{P_\mu Z_\nu}{2Z^2} = \]

\[ \frac{1}{2} \left( \frac{Z_\mu}{Z} P_\nu - \frac{Z_k P^k}{Z^2} P_\mu \frac{P_\nu}{\sqrt{Z}} \right) - \frac{1}{2} \left( \frac{Z_\nu}{Z} P_\mu - \frac{Z_k P^k}{Z^2} P_\nu \frac{P_\mu}{\sqrt{Z}} \right) = \frac{U_\mu}{2} P_\nu - \frac{U_\nu}{2} P_\mu \]

But why to use \( \frac{1}{2} U_\mu = \frac{1}{2} \left( \frac{Z_\mu}{Z} - \frac{Z_k P^k}{Z^2} \frac{P_\nu}{\sqrt{Z}} \right) \) and not simply, \( \frac{Z_\mu}{Z} \)? The reason is that \( \frac{U_\mu}{2} P_\mu = 0 \). We can therefore consider \( \frac{1}{2} U_\mu \) as a substitute for the 4-acceleration \( \frac{\alpha^\mu}{c^2} \) with the very important difference from \( V_\mu V^\mu = c^2 \)(that \( Z_\mu \) does not vanish because \( Z \) is not constant. Also notice that if \( P_\nu, q_\nu \) are parallel at every even, \( P_\nu + q_\nu \) yields the same \( U_\mu \) vector. We may now write the Lie derivative [14] of \( \frac{P_\mu}{\sqrt{Z}} \) with respect to the vector field \( \frac{P_\lambda}{\sqrt{Z}} \),

\[ \text{Lie} \left( \frac{P_\mu}{\sqrt{Z}}, \frac{P_\lambda}{\sqrt{Z}} \right) = \frac{P_\mu}{\sqrt{Z}} \left( \frac{P_\mu}{\sqrt{Z}} \right)_m + \frac{P_\mu}{\sqrt{Z}} \left( \frac{P_\mu}{\sqrt{Z}} \right)_n \]

(6)

In which the second term is positive because the differentiated \( \frac{P_\lambda}{\sqrt{Z}} \) vector has a low index.

The first term becomes,
The second term is,

\[ \left( \frac{p^m}{\sqrt{Z}} \right) \frac{p_m}{\sqrt{Z}} = \frac{p^m p_m}{Z} - \frac{p^m p_m Z_l}{2Z^2} = \frac{p^m p_m}{Z} - \frac{p^m Z_m p_l}{2Z^2} \tag{7} \]

We add (7) and (8) to get (6) and notice that

\[ \frac{p^m p_m}{Z} + \frac{p^m p_m}{Z} = \frac{p^m p_m}{Z} + \frac{p^m p_m}{Z} = \frac{Z_l}{Z} \]

from which (6) becomes

\[ \text{Lie} \left( \frac{p^m}{\sqrt{Z}}, p_l \right) = \frac{Z_l}{Z} - \frac{p^m Z_m p_l}{2Z^2} = \frac{Z_l}{Z} - \frac{p^m Z_m p_l}{2Z^2} = \frac{Z_l}{Z} \]

(9) is an interesting surprise. In (9), we also saw how we can generalize the definition of \( \frac{U_l}{2} \) to the complex case.

We are very close to define a field that describes an acceleration in space-time but we have a problem. The matrix \( A_{\mu \nu} = \frac{U_{\mu} p_{\nu} - U_{\nu} p_{\mu}}{2 \sqrt{Z}} \) is not a regular matrix. It describes rotation and scaling of the vector \( p^\nu \) into \( \frac{U_{\mu}}{2} \) by the following rule,

\[ A_{\mu \nu} \frac{p^\nu}{\sqrt{Z}} = \left( \frac{U_{\mu} p_{\nu} - U_{\nu} p_{\mu}}{2 \sqrt{Z}} \right) \frac{p^\nu}{\sqrt{Z}} = \frac{U_{\mu} p_{\nu} p^\nu}{Z} - \frac{U_{\nu} p_{\mu} p^\nu}{Z} = \frac{U_{\mu} Z}{Z} - \frac{0 p_{\mu}}{Z} = \frac{U_{\mu}}{2} \]

This is where we clearly see why uniqueness of \( \frac{U_{\mu}}{2} \) is not required, \( A_{\mu \nu} \) describes perpendicular rotation and scaling and there are more than just the two vectors, \( \frac{p_{\mu}}{\sqrt{Z}} \) and \( \frac{U_{\nu}}{2} \) that can represent \( A_{\mu \nu} \). The invariance of \( \frac{U_{\mu}}{2} \) if \( P \) is replaced by a monotone smooth function of \( P \) can be found in Appendix C. Luckily, from the matrix

\[ A_{\mu \nu} = \left( \frac{U_{\mu} p_{\nu} - U_{\nu} p_{\mu}}{2 \sqrt{Z}} \right) \]

we can infer the transformation in the plane which is perpendicular to the vectors \( \frac{U_{\mu}}{2} \) and \( \frac{p_{\nu}}{\sqrt{Z}} \) up to a constant scaling factor. To achieve this goal, we need to use the Levi – Civita alternating tensor \( E^{\mu \nu \alpha \beta} \) and not the Levi Civita alternating symbol [15]. If we use the Levi Civita symbol, we will get a tensor density and not a tensor. We then multiply \( E^{\mu \nu \alpha \beta} A_{\alpha \beta} \) and get an anti-symmetric matrix \( B^{\mu \nu} = \frac{1}{2} E^{\mu \nu \alpha \beta} A_{\alpha \beta} \). It is easy to see that

\[ B^{\mu \nu} B_{\mu \nu} = \frac{1}{2} E^{\mu \nu \alpha \beta} A_{\alpha \beta} \frac{1}{2} E_{\mu \nu \alpha \beta} A_{\alpha \beta} = A_{\alpha \beta} A^{\alpha \beta} \]

To remind the reader, the relation between a Levi – Civita symbol \( \delta^{\mu \nu \alpha \beta} \) and the Levi – Civita tensor \( E^{\mu \nu \alpha \beta} \) is brought here,

\[ E^{\mu \nu \alpha \beta} = \frac{\delta^{\mu \nu \alpha \beta}}{\sqrt{-g}} \quad , \quad E_{\mu \nu \alpha \beta} = \sqrt{-g} \delta_{\mu \nu \alpha \beta} \]

where \( g = \text{Det}(g_{\mu \nu}) \). Also please note that \( \delta^{\mu \nu \alpha \beta} \) is an alternating symbol and therefore, if it is contacted twice with the same vector, the result is zero, \( \delta^{\mu \nu \alpha \beta} V_{\alpha} V_{\beta} = 0 \). We are finally able to define an accelerating field in a covariant way. Definition, an acceleration field is:

\[ F_{\mu \nu} = A_{\mu \nu} + \gamma B_{\mu \nu} \]
Such that $\gamma \in U(1)$ and here $U(1) = \{ e^{jx} | x \in \mathbb{R} \}, j = \sqrt{-1}$. The reason for this $\gamma$ is that

$$B_{\mu \nu} = \frac{a_\mu}{c^2} \frac{V_\nu}{\sqrt{V^\alpha \alpha + V^\beta \beta}} - \frac{a_\nu}{c^2} \frac{V_\mu}{\sqrt{V^\alpha \alpha + V^\beta \beta}}$$

for some acceleration vector $\frac{a_\mu}{c^2}$ and we have a degree of freedom $\frac{a_\mu a^*_\mu}{c^2} = \frac{\gamma a_\mu \gamma a^*_\mu}{c^2}$. We also know that $a_\mu V^{*\mu} = 0$ and $a_\mu U^{*\mu} = 0$ and $U_\mu V^{*\mu} = 0$. The degree of freedom in the representation vectors, $a_\mu$ and $V_\mu$ together with the degree of freedom of $\gamma$ is $U(1) \ast SU(2)$. For a velocity $w^\nu = w^{*\nu}$ and real $F_{\mu \nu}$, we derive an acceleration,

$$F_{\mu \nu} \frac{w^\nu}{c} = \frac{a_\mu(w)}{c^2}$$

This rule appears in the Scarr – Friedman interpretation of acceleration [16] and to the author’s opinion, it did not receive enough attention from the physics community. Locally, using a real numbers scalar $q$ and $q_\mu = \frac{dq}{dx^\mu}, B_{\mu \nu}$

can be represented similar to $A_{\mu \nu}$, i.e. $B_{\mu \nu} = \frac{d}{dx^\nu} \frac{q_\mu}{\sqrt{q_\lambda q^\lambda}} - \frac{d}{dx^\mu} \frac{q_\nu}{\sqrt{q_\lambda q^\lambda}}$ and then $F_{\mu \nu} dx^{\mu \wedge} dx^{\nu} = \int \left( \frac{p_\mu}{\sqrt{P^\lambda P_\lambda}} + \frac{q_\mu}{\sqrt{q_\lambda q^\lambda}} \right) dx^{\mu}$ and therefore $F_{\mu \nu} dx^{\mu \wedge} dx^{\nu}$ becomes a Symplectic form. By Darboux theorem [17], there is a neighborhood around an event $e$ in space time, where $F_{\mu \nu}$ is not degenerated, such that in local coordinates $F_{\mu \nu}$ can be represented as

$$\begin{pmatrix}
0 & -1 & 0 & 0 \\
1 & 0 & 0 & 0 \\
0 & 0 & 0 & -1 \\
0 & 0 & 1 & 0
\end{pmatrix}$$

(15)

(13) is anti-symmetrical, $F_{\mu \nu} = -F_{\nu \mu}$ and is a regular matrix and a tensor. A short calculation immediately shows that

$$\frac{1}{4} \left( F_{\mu \nu} F^{*\mu \nu} + F_{\mu \nu} F^{\mu \nu} \right) = \frac{1}{4} \frac{U_\mu U^{*\mu} + U^{*\mu} U_\mu}{2}$$

(16)

and in the real case

$$\frac{1}{4} F_{\mu \nu} F^{\mu \nu} = \frac{U_\mu U^{*\mu}}{4}$$

(17)

(16) has the same format as the classical electromagnetic action in the General Theory of Relativity. The question that arises is, how to relate the older language of energy density to $\frac{U_\mu U^{*\mu}}{4} = \frac{a_\mu a^{*\mu}}{c^4}$, using the real case representation. To answer this question, the reader is referred to Appendix D, that shows that $\frac{U_\mu U^{*\mu}}{4}$ is the squared curvature of the integral curve along $\frac{p_\mu}{\sqrt{Z}}$. That means that as an action, $\frac{U_\mu U^{*\mu}}{4}$ does not need any constants when
used in the Einstein – Grossman action because the Ricci scalar [18] is also a curvature in units \( \frac{1}{\text{Length}^2} \). Just before we develop a new version of Einstein – Grossman equations, we can clarify the inevitable result from (17), Appendix D and [18]. In order to interpret the real numbers version \( \frac{u_{\mu}u^\mu}{4} \) as energy density, we need to multiply it by the inverse of Einstein’s gravity constant, that is,

\[
\frac{c^4}{8\pi K} \frac{u_{\mu}u^\mu}{4} = \rho = \frac{\text{Energy}}{\text{Volume}} \tag{18}
\]

But then in terms of an acceleration vector \( a_\nu \), see (3), \( \frac{u_{\mu}u^\mu}{4} = \frac{a_\lambda a^\lambda}{c^4} \) so

\[
\frac{|a|^2}{8\pi K} = \frac{a_\lambda a^\lambda}{8\pi K} = \rho = \frac{\text{Energy}}{\text{Volume}} \tag{19}
\]

\( |a|^2 \) is the squared norm of acceleration. If we compare that energy density to the classical non-covariant limit of the electrostatic field \( E \) then we have,

\[
\frac{|a|^2}{8\pi K} \approx \frac{1}{2} \varepsilon_0 |E|^2 \Rightarrow \sqrt{4\pi K \varepsilon_0} |E| \approx |a| \tag{20}
\]

where \( K \) is Newton’s gravity constant and \( \varepsilon_0 \) is the permittivity of vacuum. The relation between Minkosky norms and the classical non-covariant limit has another inevitable result,

\[
\left| \frac{u_{\mu}}{c^2} \right|^2 = \left| a_\mu \right| \approx \sqrt{4\pi K \varepsilon_0} |E| \Rightarrow \sqrt{4\pi K \varepsilon_0} \text{Div}(E) \approx \sqrt{4\pi K \varepsilon_0} \frac{\rho}{\varepsilon_0} = \frac{u_{\mu}u^\mu}{4} c^2 \tag{21}
\]

Where in (21), \( \rho \) means charge density. The divergence of the Reeb vector has a classical non-covariant limit which is proportional to the divergence of the electrostatic field and therefore to charge density. (20) means in the old language of physics, that the energy of the electric field is in a very weak acceleration field. We need to understand the interaction between positive and negative charge. By Occam’s razor, this can only mean an alignment of the Reeb vector in the classical limit with the classical electric field and that the time-like component of acceleration is very small. Interacting negative and positive charge therefore, must reduce the integration of \( \frac{u_{\mu}u^\mu}{4} \). So how big is this acceleration field and what exactly does it accelerate? (20) gives us a dauntingly small result, \(~8.61 \text{cm}/\text{sec}^2\), which is less than 0.01 g, if the electrostatic field is an immense 1,000,000 volts over 1 mm. We will see that due to unexpected gravity by electric charge, this acceleration is even smaller, about 4.305 \text{cm}/\text{sec}^2.

We now need to develop the Euler Lagrange equations of the following action, as the minimum action problem in the General Theory of Relativity language,

\[
Z = \left| P_{\mu}P^\mu \right|, Z_\mu = \frac{dZ}{dx^\mu}, U_\lambda = \frac{Z_\lambda}{Z}, U_k = U_k^k, L = U_k U_k \tag{22}
\]

and \( R \) is the Ricci curvature [18]. The integral to be minimized over coordinates domain \( \Omega \) is

\[
\text{Action} = \text{Min} \int_\Omega \left( R - \frac{1}{4} U_k U_k \right) \sqrt{-g} \text{ d}\Omega \tag{23}
\]

Locally, this can be written in Regge’s Causal Dynamic Triangulation \( \sigma(n + 1) \) formalism too [19], where \( n \) is the dimension.

\[
\text{Action} = \text{Min}(\Sigma_h \left( 2\delta_h - \frac{a_h^2}{c^4} \right) V_h) \tag{23.1}
\]

Where \( h \) is a hinge in a Causal Dynamic Triangulation, the missing angle is \( \delta_h = 2\pi - \sum_{h \in \sigma(n)} \theta_h \) and \( \theta_{h \in \sigma(n)} \) is the angle between two faces around \( h \) that share the hinge \( h \) and \( a_h \) is the acceleration through the two dimensional hinge \( h \).
From Appendix A there are two minimum action equations, one by the metric $g_{\mu\nu}$ and one by the Morse function, $P$. The results are, in $\{-1, +1, +1, +1\}$ metric convention,
\[
\frac{1}{4} \left( U_{\mu} U_{\nu} - \frac{1}{2} g_{\mu\nu} U_{\lambda} U^{\lambda} - 2 U^{k; k} \frac{p_{\mu} p_{\nu}}{Z} \right) = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \tag{24}
\]
And
\[
W^{\mu; \mu} = \left( -4 U^{k; k} \frac{p_{\mu} p_{\nu}}{Z} - 2 \frac{Z_{\nu} p_{\nu}}{Z} U^{\mu} \right); \mu = 0 \tag{25}
\]
From (25), the following divergence is zero,
\[
\frac{1}{4} \left( U_{\mu} U_{\nu} - \frac{1}{2} g_{\mu\nu} U_{\lambda} U^{\lambda} - 2 U^{k; k} \frac{p_{\mu} p_{\nu}}{Z} \right); \nu = 0. \text{ A proof can be found in Appendix B.}
\]
Note: By (21), we immediately see a peculiar result, electric charge generates gravity in an unexpected way, by the term
\[
\frac{1}{4} \left( -2 U^{k; k} \frac{p_{\mu} p_{\nu}}{Z} \right) = -\frac{1}{2} U^{k; k} \frac{p_{\mu} p_{\nu}}{Z}, \text{ which is peculiar because } p_{\mu} \text{ may not be aligned with the motion of the “source” of the Reeb vector } \frac{u_{\mu}}{2}. \text{ What is the meaning of “source”? From the theory of foliations, integration of the reduced Reeb vector [12] in leaves of foliations perpendicular to } P_{\mu} \text{ is zero along closed curves [12] and the integration of } \frac{u_{\mu}}{2} \text{ along leaf-wise curves measures the transverse holonomy expansion. In other words, the field } \frac{u_{\mu}}{2}, \text{ when reduced to specific three dimensional leaves, behaves exactly as a classical electric field that has a source, as a negative or as a positive charge. (24) can be generalized for a complex } P, \text{ and } P \text{ can be described as a sum of a Hilbert orthogonal functions, } \psi(1), \psi(2), \ldots, \psi(n), \text{ as } n \to \infty.
\]
\[
P = \lim_{n \to \infty} \psi(1) + \psi(2) + \cdots + \psi(n)
\]
\[
\int_{\Omega_{4}} \psi(k) \psi^* (k) \sqrt{-g} d\Omega_{4} = 1
\]
\[
0 < j < k \implies \int_{\Omega_{4}} \psi(j) \psi^* (k) \sqrt{-g} d\Omega_{4} = 0
\]
\[
\frac{1}{8} \left( U_{\mu} U_{\nu} * \gamma + U * \mu U_{\gamma} - \frac{1}{2} (U * k U_{k} + U^{k} U * k) g_{\mu\gamma} - 2 (U^{k; k} + U * k; k) \frac{(p_{\mu} p_{\nu} + p_{\mu} p_{\nu})}{2Z} \right) - \frac{1}{2} \lambda g_{\mu\nu} =
\]
\[
R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \tag{26}
\]
For some cosmological constant $\lambda$ whose units are $\frac{1}{\text{Length}^2}$. The following constraints can only have one physical meaning, they describe events and not particles.
\[
P = \lim_{n \to \infty} \psi(1) + \psi(2) + \cdots + \psi(n)
\]
\[
\int_{\Omega_{4}} \psi(k) \psi^* (k) \sqrt{-g} d\Omega_{4} = 1
\]
\[
0 < j < k \implies \int_{\Omega_{4}} \psi(j) \psi^* (k) \sqrt{-g} d\Omega_{4} = 0 \tag{27}
\]
More profoundly, the meaning of (14) is of an acceleration due to non–geodesic alignment of these events. A reasonable interpretation of (13), (14) is that material clocks or as we see in (17), neutral “electromagnetic energy”, even with a total charge 0, cannot move along geodesic curves if the events do not align along geodesic curves. In the model in this paper, charge interacts by increasing or decreasing the energy of the weak acceleration field which results in force due to
\[
\frac{1}{4} \left( U_{\mu} U_{\nu} - \frac{1}{2} g_{\mu\nu} U_{\lambda} U^{\lambda} - 2 U^{k; k} \frac{p_{\mu} p_{\nu}}{Z} \right); \nu = 0. \text{ We may have already observed the small acceleration (20) as the Flyby Anomaly [20] above thunderstorms. The Flyby anomaly does not support an energy density } \frac{|q|^2}{4\pi K} \text{ in which charge induced gravity and charge induced acceleration would cancel out. It does support}
values such as (19). A field that is a sum of events, quite similar to (27) was already offered in 1982 by Sam Vaknin [11], later, a more set-theory oriented idea was offered by Sorkin in years 1987 and 2000 [21]. Sam Vaknin’s work was published in modern variations by other researchers [22] around 2001 – 2002. The approach in this paper does not derive from the Dirac spinor equation [23] as Sam Vaknin’s work and as [22] did.

**Unexpected gravity and anti-gravity by electric charge**

From (21) and (24), \( \frac{U^\mu}{2} = \frac{a^\mu}{c^2} \) where \( a^\mu \) is a 4-acceleration representative of the field and \( Z = p^\mu p^\nu \).

\[
\frac{c^4}{8\pi K} \sqrt{\frac{4\pi K}{\varepsilon_0}} \frac{\rho c^2}{c^2} = \frac{\text{Vol}^\nu}{16K\varepsilon_0} = \frac{c^4}{8\pi K} \frac{U^\mu U^\nu}{2} \approx \frac{V^\mu V^\nu}{c^2}
\]

(28)
as in special relativity but with a very important caveat that \( p^\mu \) need not be aligned with any motion \( V^\mu \) as it is merely a gradient of a scalar field and not a velocity \( V^\mu \),

\[
V^\mu = \frac{(c,v_x,v_y,v_z)}{\sqrt{1-v^2/c^2}}
\]

(28.1)

Although \( p^\mu \) is not velocity \( V^\mu \), \( \frac{U^\mu U^\nu}{2} \) behaves as \( \frac{a^\mu a^\nu}{c^2} \) where \( a^\mu \) is acceleration, \( c \), the speed of light, and we saw in (20) \( \sqrt{4\pi K\varepsilon_0}|E| \equiv |a| \) so \( \frac{U^\mu U^\nu}{2} \) is as in the non-covariant limit \( \sqrt{\frac{4\pi K}{\varepsilon_0} c^2} \) where \( \rho \) is charge density and with

\[
\frac{p^\mu p^\nu}{Z} \approx \frac{V^\mu V^\nu}{c^2}
\]

and therefore,

\[
\frac{1}{8\pi K} \frac{U^\mu U^\nu}{2} \approx \frac{1}{8\pi K} \frac{4\pi K}{\varepsilon_0} \frac{\rho_{\text{charge}} V^\mu V^\nu}{c^4}
\]

(28.2)

which generates gravity equivalently to energy density. So

\[
\frac{8\pi K}{c^4} \left( \frac{1}{8\pi K} \frac{4\pi K}{\varepsilon_0} \rho V^\mu V^\nu \right)
\]

(28.3)

That can only mean that the following generates gravity similar to mass but can also be negative:

\[
M = \frac{Q}{\sqrt{16\pi K\varepsilon_0}}
\]

(29)

If we multiply the square root of the Fine Structure Constant by the Planck mass we get a term that attests to the validity of (29),

\[
\frac{\sqrt{\hbar c}}{K} \frac{e^2}{4\pi \varepsilon_0 \hbar c} = \frac{2e}{2\sqrt{4\pi K\varepsilon_0}} = \frac{2e}{\sqrt{16\pi K\varepsilon_0}}
\]

(29.1)

Which means that the electromagnetic portion of mass we get from the Planck mass is twice the gravitational mass by the positron charge \(+e\), an expected term for a pair of positive and negative charged particles whether manifesting positive or negative gravity.

generates gravity as mass does, however, by the note after (25) the motion of the electric charge need not be aligned with \( P^\mu \) which differs from a usual inertial mass. (29) yields an assessment

\[
\pm 1\text{Coulomb} \sqrt{\frac{1}{16\pi K\varepsilon_0}} \approx \pm 5.80213 \times 10^9 K g
\]

(30)

From (30) and from (20) we can approximate the total acceleration \( |a| \) near a charge \( Q \) at radius \( r \) as,
Term (30) requires an experimental evidence which does exist, see Timir Datta’s research [24] and a Brazilian experiment [25]. [24] involves high concentration of separated charge close to a tip of a cone. We do not know yet, whether the positive charge or whether the negative charge act as negative energy, however, evidence of free electrons in the galactic center [26] suggests ionization of the galactic center where these electrons move out from the galactic center. We therefore expect the edges of a spiral galaxy to be more negative than its center. Equation (30) by Miguel Alcubi erre [27], inevitably generates Warp Drive between the edges of the galaxy and its center, providing that a sufficiently large amounts of electric charge are separated. Since galaxies seem to be stable, the Alcubirre Warp Drive must push the galaxies towards the center. We are inevitably led to the conclusion that it is the negative charge that generates negative gravity in the peripheral mass of the galaxy, where the center of the galaxy is positively charged and generates more gravity than expected. By (30) and [27], static charge separation cannot generate a technologically feasible warp drive without dynamic charge oscillation and/or rotation because an Alcubierre warp drive requires $\pm 10^{27}$ Kg $\approx \pm 2 \times 10^{17}$ Coulombs. The acceleration field around a charge must be opposite in sign to the gravity generated by the positive charge and to the anti-gravity generated by the negative charge. In the language of modern cosmology, electrons generate “negative pressure” and protons generate “pressure” on the neutral clocks they weakly accelerate and cause space-time to respond by gravity. Before we move on, it is worth mentioning that the action $\frac{1}{4} U^k U_k$ can be generalized to more than one Reeb vector. This idea is quite straightforward by roots of determinants of Gram matrices of Reeb vectors and is discussed at the end of Appendix F. Taking the third root of a determinant of the Gram matrix of three complex Reeb vectors, has an $SU(3)$ symmetry. There are other ways to consider $SU(3)$ symmetry, see Appendix F for the conditions of one interesting option. To summarize the idea of an acceleration field that acts on mass in a different way than gravity, if $V^\mu = \frac{\nu^\mu}{c}$ is a unit velocity of a clock in the hyper-plane spanned by $\frac{t^\mu}{z}$ and $\frac{p^\nu}{\sqrt{2}}$ then the acceleration $\frac{dV^\mu}{d\tau}$, where $\tau$ is proper time, is also in that plane and the clock reference frame is accelerated by the rule,

$$\frac{dV^\mu}{d\tau} = \frac{dV^\mu}{d\tau} \frac{dx^k}{d\tau} = V^\mu,^k V^k = \left( \frac{P_k U^\mu}{2\sqrt{|z|}} - \frac{U_k P^\mu}{2\sqrt{|z|}} \right) V^k$$

Then multiplying both sides by $\frac{P^\mu}{\sqrt{2}}$ we get $\frac{dV^\mu}{d\tau} \frac{P^\mu}{\sqrt{|z|}} = \left( \frac{P_k U^\mu}{2\sqrt{|z|}} - \frac{U_k P^\mu}{2\sqrt{|z|}} \right) V^k \frac{P^\mu}{\sqrt{|z|}} = -V^k U_k \frac{z}{2}$.

In the same way, assuming $U^A U^A \neq 0$ yields $\frac{dV^\mu}{d\tau} \frac{U^\mu}{\sqrt{|U^A U^A|}} = \left( \frac{P_k U^\mu}{2\sqrt{|z|}} - \frac{U_k P^\mu}{2\sqrt{|z|}} \right) \sqrt{|U^A U^A|} U^A = \sqrt{|U^A U^A|} U^A \frac{P^\mu}{\sqrt{2}}$.

So we see that $V^k$ was transformed into a vector with the pseudo-norm $\frac{|U^A U^A|}{2}$ because,

$$\left\| \frac{V^k}{2} \right\| \frac{P^\mu}{\sqrt{|z|}} + \left\| V^k \frac{P^\mu}{\sqrt{|z|}} \right\| \frac{U^\mu}{\sqrt{|U^A U^A|}} = \left\| V^k \frac{P^\mu}{\sqrt{|U^A U^A|}} \right\|^2 + \left\| V^k \frac{P^\mu}{\sqrt{|U^A U^A|}} \right\|^2 = \frac{|U^A U^A|}{4}$$

Notice that it is possible that $\frac{U^A U^A}{4} \leq 0$ in $(+1, -1, -1, -1)$ convention and $\frac{U^A U^A}{4} \geq 0$ in $(-1, +1, +1, +1)$ convention.
Particle mass ratios by added or subtracted area – Muon to electron mass ratio

This section, unlike the previous ones, relies on scaled area ratios. Some of the work can be seen in the remarkable paper of Lee C. Loveridge [28] and is somewhat speculative unless (34.4) is taken into account or claims of a ~40eV neutrino [29] and of a ~0.0002 eV resonance are taken into account and therefore the reader is honestly advised to take it with a grain of salt especially as this model requires a scaling factor \( \frac{3}{32} \) that can be shown justifiable due to fractional quantization in a Hole Quantum Wire and Steiner Tree optimization in Causal Dynamic Triangulation that will be discussed, however, the results are very interesting and do have a geometrical basis in [12] that worth further investigation - the Reeb vector on the foliation leaves must have drains and sources. It is therefore the authors opinion that the technique used here will be considered by the reader even if the reader does not fully agree with the idea which is presented in this section. The following development has its roots in the lectures of professor Seth Lloyd of the M.I.T [30] and in a paper by Ted Jacobson [31] combined with equation (24). A method to reduce (24) from 4 dimensional Minkowsky space to Riemannian two dimensions will be discussed along with its possible applications to mass ratios between particles and to the fine structure constant. There are two principles that this section is based upon,

A) The holographic spin principle means that infinitesimally, the disk surrounding a particle which is perpendicular to its spin is affected by gravity, not the surrounding sphere. This idea complies with the quantum principle that a spin either points to an observer or away from an observer. For example, if the spin is at angle 60 degrees with the observer, the probability of the spin pointing to at the observer is \( \cos(60) \times \cos(60) = 0.25 \).

In terms of area loss, the infinitesimal radius \( r \) disk area loss if a particle is pointed to by a unit vector \( \frac{p^\mu}{\sqrt{Z}} \) is

\[
\delta \text{Area} = \frac{\pi}{12}\left(R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}\right) \frac{p^\mu p^\nu}{Z} r^4 = \frac{\pi}{24} \left(R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}\right) \frac{p^\mu p^\nu}{Z} r^4
\]

this area’s portion in relation to the sphere surrounding the particle is \( \frac{1}{96} \left(R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}\right) \frac{p^\mu p^\nu}{Z} r^2 \). The holographic spin principle, which stems from physics, dictates that area portion ratios that originate from Einstein’s equation will have a physical meaning only with the factor \( \frac{1}{96} \) or in another form \( b = 1 + \frac{1}{96} \left(R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}\right) \frac{p^\mu p^\nu}{Z} r^2 \) which is by how much a sphere would relatively grow or shrink by adding or subtracting disk delta area due to Einstein’s curvature.

B) Given \( b \) from principle (A), \( (b-1) \) area ratio around a negative charge and \( (1-b) \) area ratio around a positive charge of spin \( \frac{1}{2} \) will have a physical meaning as energy ratios between one charged lepton that decays into another, unless the decay is mediated by a spin 1 boson. The Tau lepton decay into a Muon through a virtual W boson is expected to yield an energy ratio of the form \( \sqrt{b-1} \). Energy ratios between a neutral particle and a lepton should involve terms of the form \( (a - 1)(1 - b) \) for spin \( \frac{1}{2} \) and \( \sqrt{(a-1)(1-b)} \) for spin 1 electro-weak bosons. The reason for this will be clearer once the paper presents an offer for an electro-weak action, see (93) and the square root of a Gram matrix determinant. Principles (A) and (B) will be the main tools in the results (43), (43.8), (43.12) where in (43.12) a maximum allowed ratio between area ratios, of area subtracted around a positive charge and of area added around a negative charge is also assumed. Decay of the W boson into Tau is expected to depend on spin \( \frac{1}{2} \) gravitational area ratios where decay into a W boson is expected to involve a root, e.g. \( \sqrt{b-1} \).
C) A coefficient that regards field configuration is assumed in (43), (43.8), (43.12), (43.17.2). One such coefficient is taken directly from Ettore Majorana’s notebook.

Pure mathematics without consideration of the holographic spin principle does not readily result in the factor $\frac{1}{96}$. Einstein tensor means added or subtracted area of the sphere in a ball with an infinitesimally small radius $r_0$ which is Minkowsky perpendicular to a unit vector $\frac{P_\mu}{\sqrt{Z}}$ by the equation,

$$-2 \left( R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \right) \frac{P_\mu P_\nu}{Z} = R(3) = 6 \cdot (\text{Dimension} = 3) \lim_{r_0 \to 0} \frac{1 - \frac{A}{4\pi r_0^2}}{r_0^2} = \lim_{r_0 \to 0} \frac{18\delta A}{4\pi r_0^4}$$

$$= - \left( R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \right) \frac{P_\mu P_\nu}{Z} \frac{4}{9\pi} r_0^4 \to \delta A$$

Where $R(3)$ is the scalar curvature in three dimensions.

$$\frac{4}{9\pi} \left( R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \right) \frac{P_\mu P_\nu}{Z} r_0^4 = -\text{Sphere Area}, \text{so by (24), in } (-1, +1, +1, +1) \text{ metric convention,}$$

$$\left( \frac{4}{9\pi} \right) \frac{1}{4} \left( U_k U_k - \frac{1}{2} g_{\mu\nu} U_k U_\lambda - 2 U_k U_k \right) \frac{P_\mu P_\nu}{Z} r_0^4 = \frac{4}{9\pi} \left( R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} \right) \frac{P_\mu P_\nu}{Z} r_0^4 = -\delta A \quad (31)$$

$$U_\mu P^\mu = 0, g_{\mu\nu} \frac{P_\mu P_\nu}{Z} = 1, \frac{P_\mu P_\nu}{Z} = 1 \text{ and therefore (31) yields}$$

$$\left( \frac{4}{9\pi} \right) \left( -\frac{1}{8} U_k U_k - \frac{1}{2} U_k U_k \right) r_0^4 = -\delta A \quad (32)$$

This is not yet the result of contracting the Einstein tensor with a unit vector twice $\left[ 30, 31 \right]$ that we want. The factor 2 can also be found in the outstanding work of Lee C. Loveridge $\left[ 28 \right]$. We now make an assumption that in the subatomic scale, there is a relation between the acceleration $a_\mu$ and the radius $r_0$, which is presented by

$$\sqrt{\frac{|U_\mu U^\mu|}{4}} = \frac{|U_\mu|}{2} = \frac{|a_\mu|}{c^2} = \frac{\zeta}{r_0} \text{ where } \zeta \text{ depends on the field and } c \text{ is the speed of light, also the divergence is calculated along the radius } r_0 \text{ and we also scale } r_0 \text{ by } \frac{3}{32} \text{ to get } \frac{3}{32} \frac{4}{9\pi} = \frac{\pi}{24}. \text{ This scaling is not trivial and it leads to some interesting results. In terms of ordinary Riemannian geometry, the curvature of circle of radius } r_0 \text{ is } \frac{1}{r_0^2}. \text{ We also know that by Gauss law, if area is added around a charge, the intensity of the electric field is reduced and we expect the same rational reduction in the acceleration field. If the area grows by } x \text{ then the field is reduced by } \frac{1}{x} \text{ so, } \frac{\zeta}{r_0} \text{ becomes } \frac{\zeta}{r_0 x}. \text{ Consider that the divergence is calculated along a distance } r_0, \text{ which means that there is a minimal distance along which the field can change from } \frac{\zeta}{r_0 x} \text{ to } 0. \text{ So the divergence term becomes,}$$

$$\left( \frac{\zeta}{r_0 x} - 0 \right) \frac{1}{r_0} = \pm \frac{\zeta}{r_0^2 x^2} \quad (32) \text{ then yields the following,}$$

$$\left( \frac{4}{9\pi} \right) \left( -\frac{1}{8} U_k U_k - \frac{1}{2} U_k U_k \right) r_0^4 = \frac{3}{32} \frac{4}{9\pi} \left( -\frac{1}{2} \frac{\zeta^2}{r_0^2 x^2} \pm \frac{\zeta}{r_0 x} \right) r_0^4 = \frac{\pi}{24} \left( -\frac{1}{2} \frac{\zeta^2}{x^2} \pm \frac{\zeta}{x} \right) r_0^2 = \pm \delta A \quad (33)$$

We now divide this area by the area of a sphere $4\pi r_0^2$. 

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and we also know that \( X = \frac{4\pi r_0^2 \pm \text{Area}}{4\pi r_0^2} \) from which we infer after fixing the ratio \( \frac{\pi}{24} \) for \( \delta \text{Disk}/4\pi r_0^2 \),

\[ 1 + \frac{1}{96} \left( -\frac{1}{2} \frac{x}{x} \pm \frac{x}{x} \right) = x \implies x^2 + \frac{1}{96} \left( -\frac{1}{2} \frac{x^2}{x} \pm \frac{x}{x} \right) = x^3 \quad (34) \]

Note: The real reason for the \( \pm \frac{1}{96} x \) is because an electromagnetic field of a charged particle can be viewed as comprised of two perpendicular components, see (13). An interaction of such a field can be said to be related to the Fine Structure constant Alpha. So the fine structure constant should be actually multiplied by the square root of two in order to relate to an interaction with two perpendicular components of unit length. The following is the outcome of such a consideration:

\[ \frac{1}{\text{Alpha} \sqrt{2}} = \text{Alpha}^{-1} \cdot 2^{\frac{1}{2}} \approx 137.035999149 \cdot \sqrt{0.5} \approx 96.899 \quad (34.1) \]

The nearest integer less than that value is:

\[ \left\lfloor \frac{1}{\text{Alpha} \sqrt{2}} \right\rfloor = 96 \quad (34.2) \]

Now if we look at the following relation between the inverse of the Fine Structure Constant and the number 96:

Then we see that it is not difficult to reach these two equations if we notice that

\[ (1 + \frac{1}{y})^{95} \approx 1.9951740, \quad (1 + \frac{1}{y})^{96} \approx 2.0097335 \text{ for the choice } y \approx 137.0359990368270076 \quad (34.2.1) \]

This also explains the motivation for the choice of 96 as an upper limit of power. The idea behind \( 1 + \frac{1}{y} \) is that it expresses an increase in area ratio by \( \frac{1}{y} \); 2 means that the initial energy is doubled. These two numbers 95, 96 will be used to derive the mass ratio between the Muon and the electron. The idea of using powers will lead us to a derivation of the inverse Fine Structure Constant around 137.035999035747181551 but that can wait for now.

The reader may come to the conclusion that the author first threw a dart and then drew a circle around it and it has to be admitted to be somewhat true. The value \( \frac{1}{96} \) was indeed intended due to (34.2) which is also a result of the interpretation of (13) as two perpendicular acceleration fields as the reason for the energy of the electric field.

Note: The easy way to get \( \frac{1}{96} \) though not intuitive either, is to divide the area loss from a disk perpendicular to \( P_\mu \) which is \( \delta \text{Disk Area} = \pi \frac{1}{12} \left( R_{\mu \nu} - \frac{1}{2} R g_{\mu \nu} \right) \frac{\rho_{\rho\sigma\nu}}{\rho_{\rho\lambda}} r_0^4 = \frac{1}{12} \left( -\frac{1}{8} U_k U^k - \frac{1}{2} U^k U^k \right) r_0^4 \) by the area of a sphere \( 4\pi r_0^2 \) so we have

\[ \frac{\pi}{24} \left( -\frac{1}{2} \frac{\xi^2}{x^2} \pm \frac{\xi}{x} \right) = \frac{1}{96} \left( -\frac{1}{2} \frac{\xi^2}{x^2} \pm \frac{\xi}{x} \right) = \frac{\pm \delta \text{Disk Area}}{\text{Euclidean Sphere Area}} \quad (34.3) \]
Intuitively, we would expect \( \frac{\pm \delta \text{Sphere Area}}{\text{Euclidean Sphere Area}} \) to have a physical meaning when reducing a 4 dimensional Minkowsky geometry to a 2 dimensional Riemannian geometry.

(Fig. 1.) - The M.I.T public lecture by Seth Lloyd on YouTube, notice that (34.3) divides \( \frac{\pi}{24} \) by \( 4\pi \) and not by \( \pi \):

The most surprising argument for 1/96 comes from a simple geometric consideration. Due to (34) consider the following, where \( \zeta = \frac{4}{\pi} \) will be discussed.

\[
c_1^2 + \frac{1}{n} \left( -\frac{1}{2} \left( \frac{4}{\pi} \right)^2 + \frac{4}{\pi} c_1 \right) = c_1^3, \quad c_2^2 + \frac{1}{n} \left( -\frac{1}{2} \left( \frac{4}{\pi} \right)^2 - \frac{4}{\pi} c_2 \right) = c_2^3
\]

Now, look for the following minimum for a natural \( n \):

\[
\min_n \left( (c_1 - 1)(1 - c_2) \right)^{\frac{1}{2}} - n - \frac{2\pi}{n} \Rightarrow n = 96 \quad (34.4)
\]

For \( n=95 \) we get \(-0.015965593\), for \( n=96 \), \( 0.003727368 \) and for \( n=97 \), \( 0.023393663 \). In plain English, \( n=96 \) minimizes the delta between the residual \( (c_1 - 1)(1 - c_2) \) and the angle \( \frac{2\pi}{n} \). For \( n = 96 \) where \( (c_1 - 1)(1 - c_2) \approx 12 \pm 96.06917721 \). These calculations where done, using Excel Datasheet and the reader may prefer Python or C,C++. By professor Ted Jacobson [31], area in the Planck scale is equivalent to energy and therefore if his claims are correct, we should be able to obtain known mass ratios between charged particles, based on area ratios, so \( \frac{\pm \delta \text{Area}}{4\pi r_0^2} \) which is \( X - 1 \) should represent at least one of the known mass ratios in the particles world. The problem here is that we do not have a full solution to either (24) or (26) and we honestly have to make an educated guess about possible values of \( \zeta \) and get different \( X - 1 \) values as mass ratios. We do that by using the coefficient of a normalized ring potential which is reminded in the note book of Ettore Majorana [32] as \( V(r) = \frac{4}{\pi} \frac{Q}{r} \). A phenomenological view of the electron as a ring can be seen in the work of O.F. Schilling [33]. Our most obvious first target is the Muon / Electron mass ratio as nominated to have a gravitational reason. Instead of accepting \( \zeta = \frac{4}{\pi} \) as factual, we can also say that due to the spin of the electron, the acceleration field around the electron is not evenly distributed, to compensate for that, at any given time, \( \zeta \) has to be bigger than 1, otherwise Gauss law would be violated. One simple model is of a field which is maximal at an equator of some sphere around the electron and vanishes at its poles. “Equator” means some maximal length circle in which the radius is typically perpendicular to an axis of rotation. Such a model yields the following value,

\[
\frac{\left| a \right|}{c^2} = \text{Constant} \cdot \frac{1}{r_0} \cos (\theta) \quad (35)
\]

Unlike in (35) a uniform field around a negative charge at a given local time would look as in the followings illustration that represents 3 dimensions, where the dark arrow illustrates an acceleration of a material clock in the
electric field. We can also imagine a neutral particle as if having an alternating field pointing in and out of the particle rather than $2(U^k_{;k} + U^*_k_{;k}) = 0$.

(Fig. 2.0) \[ 2(U^k_{;k} + U^*_k_{;k}) \neq 0. \]

where $\theta$ is the angle from the “equator” and $|a| = \sqrt{|a^\mu a_\mu|}$ is the absolute value of a Minkowsky norm of an acceleration $a^\mu$. Field integration on two hemispheres is then,

\[
\begin{align*}
2 \int_{\theta=0}^{\theta=\pi} \cos(\theta) \cdot 2\pi r_0 \cdot \cos(\theta) \cdot r_0 \cdot d\theta &= \\
&= 4\pi r_0 \left[ \frac{1}{2} \int_{\theta=0}^{\theta=\pi} \cos^2(\theta) d\theta \right] = 4\pi r_0 \cdot \frac{\pi}{4}
\end{align*}
\]

If the field was uniform then the integration would be

\[
2 \int_{\theta=0}^{\theta=\pi} 2\pi r_0 \cdot \cos(\theta) \cdot r_0 \cdot d\theta = 4\pi r_0
\]

And the ratio between (37) and (36) is $\frac{4}{\pi}$, which agrees with Ettore Majorana’s normalized ring potential coefficient [32] as $V(r) = \frac{4}{\pi} \frac{r}{r}$, which means that the field $\frac{1}{r_0}$ in (36) has to grow by a factor of $\frac{4}{\pi}$ in order to sum up as in (37). In the charged particle case, we can see $U_\lambda$ as a vector that points toward or outward of an integral curve in space-time but that the Minkowsky norm of the field is always the same, only the probability that this vector points towards a certain direction in space-time changes. This idea leads to the compensating scaling value $\zeta = \frac{4}{\pi}$. So equation (34) around a negative charge becomes,

\[
\chi^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{4}{\pi} \right)^2 + \frac{4}{\pi} \chi \right) = \chi^3
\]

(38)
The roots of (38) are \( x_1 \approx 1.004836728026 \), \( x_2 \approx 0.760783659050 \), \( x_3 = 0.089280264357 \). The first root is the only
root which is attainable through iterations

\[
x(n + 1) = \left( \frac{192x^2(n) + 2\frac{4}{\pi}x(n) - \left(\frac{4}{\pi}\right)^2}{192} \right)^{\frac{1}{3}}
\]

and it converges starting from 0.1 or from 2 or any other positive number.

But \( x = 1 + \frac{\text{Added or subtracted area}}{4\pi r^2} \), is either bigger or smaller than 1. The ratio

\[
\frac{\text{Added or subtracted area}}{4\pi r^2} = \frac{1}{x-1} \approx 206.751340,
\]

which is very close to the ratio between the mass of the Muon and the mass of the Electron, 206.768277. We followed the M.I.T professor Seth Lloyd offer that addition or subtraction of quantum area, means addition or subtraction of energy and we have reached 206.751340. For the solution around a positive charge

\[
x^2 + \frac{1}{96} \left(-\frac{1}{2} \left(\frac{4}{\pi}\right)^2 - \frac{4}{\pi} x\right) = x^3
\]

we get about 44.63955017596401. These ratios ~1/45 and ~1/207 could mean a decay path for charged leptons where the numerical stability of 1/45 is worse than that of 1/207. A more exact root to (38) yields,

\[
\frac{4\pi r^2}{\delta A} = \frac{1}{x-1} \approx 206.75133988502202
\]

The difference in accuracy in this value by 64 and 128 bits is just the last 2 - 3 digits. If we divide the Muon energy by this value, we get very close to the energy of the electron and the delta in Mega electron volts is: 105.6583745 ±0.0000024 MeV, we add 0.01 of this error to the Muon and 0.01 of the electron energy error and divide

\[
(105.6583745 \pm 0.01*0.0000024) \text{ MeV} / 206.75133988502202\text{ MeV} - 0.510998946131 \text{ (MeV)} = 0.00004187509298 \text{ MeV},
\]

which is 41.875092980 eV and surprisingly fits the Super Nova 1987a ~40 eV Neutrino claim by Cesare Bacci [29]. That energy is small but beyond the energy of any known Neutrino mass. It is an unknown energy. Should it be a new particle, this particle is beyond the Standard Model. The ratio between the electron’s energy and this energy 41.875092980 eV is 0.510998946131 MeV / 0.00004187509298 which is approximately, 12202.93281199539440, almost 12203. For Muon energy 105.6583745 MeV and an electron energy 0.5109989461 MeV, the ratio is 12202.95760492718728.

We can get this value if we check the following polynomials for \( \zeta = 1 - \frac{1}{96} \), see (34):

\[
x^2 + \frac{1}{96} \left(-\frac{1}{2} \left(1 - \frac{1}{96}\right)^2 + \left(1 - \frac{1}{96}\right) x\right) = x^3
\]

Which is 1 + or 1 - the portion of area added around a negative charge or subtracted around a positive charge such that the acceleration field is smaller by a factor of \( \zeta = 1 - \frac{1}{96} \). The idea to use a damping of \( 1 - \frac{1}{96} \) is because of the factor \( \frac{1}{96} \) in (34). This implies that charge quanta can be of the order \( \frac{1}{96} \) of the charge of the electron e. The compelling indirect evidence for the \( \frac{1}{96} \) in (43) can be seen in [34] as it appears that a fractional quantization in a Hole Quantum Wire yields a lowest fraction \( \frac{e}{32} \) and with the direct \( \frac{e}{3} \) fraction from Quarks we get \( \frac{e}{96} = \frac{e}{32} \),

which explains \( \frac{1}{96} \) as the coefficient of electric charge. In resemblance to (38), the two polynomials in (41) with the ± sign have 3 roots each and the big roots are \( x_1 = 1.00520707510980 \) for (+Area) and \( x_2 = 0.98426221868924 \) for (-Area).
\[
\left(\frac{1}{x_1-1}\right)\left(\frac{1}{1-x_2}\right) = 12202.88874066467724 \tag{42}
\]
which with numerical accuracy is even closer to 12202.95760492718728. Other choices except for 96 in \(\zeta = \left(1 - \frac{1}{96}\right)\) are further away from 12202.95760492718728 even for small differences after 3 digits after the floating point. The number \(1 - \frac{1}{96} = \frac{95}{96} = \left(\frac{96}{95}\right)^{-1}\) is the inverse of the Steiner Tree Problem limit. Finding as sub-optimum below \(\frac{96}{95}\) of the minimal length of the Steiner Tree [35] that spans a graph with terminal points – in our case on a sphere - is not solvable in polynomial time. This limit strongly suggests a Dynamic Causal Triangulation [19] approach to space-time in (41), i.e. it describes a field truly in the Planck scale. We have to compromise on a graph that its length fits a larger ball by a factor \(\frac{95}{96}\) which means that an acceleration field that depends on \(\frac{1}{r_0}\) will be smaller, \(\frac{95}{96}\). Recalling (39), (40), \(1 + \frac{1}{96}\left(\frac{1}{2}\left(\frac{4}{\pi}\right)^2 c^{-2} + \frac{4}{\pi} c^{-1}\right) = c\). So the following system of biggest roots as area ratios is

\[
\begin{align*}
1 + \frac{1}{96}\left(\frac{1}{2}\left(1 - \frac{1}{96}\right)^2 a^{-2} + \left(1 - \frac{1}{96}\right) a^{-1}\right) &= a \\
1 + \frac{1}{96}\left(\frac{1}{2}\left(1 - \frac{1}{96}\right)^2 b^{-2} - \left(1 - \frac{1}{96}\right) b^{-1}\right) &= b \\
1 + \frac{1}{96}\left(\frac{1}{2}\left(\frac{4}{\pi}\right)^2 c^{-2} + \frac{4}{\pi} c^{-1}\right) &= c
\end{align*}
\]

\[
\text{MuonMass} \ast (c - 1) = \text{ElectronMass} + \text{ElectronMass} \ast (a - 1)(1 - b)
\]

\[
\text{MuonMass} \frac{(c-1)}{(1+(a-1)(1-b))} = \text{ElectronMass} \tag{43}
\]

which is \(-0.51099894597978\) MeV instead of \(0.51099894613\) MeV. The difference -0.00015122 eV may be related to the electron neutrino. The minus sign means this energy is required in addition to the gravitational energy of the Muon in order to create an electron. A Muon mass of 105.65837455 yields an electron mass 0.510998461 MeV. In Seth Lloyd’s and Ted Jacobson’s terminology, the physical meaning of this finding could be that the energy of the electron is the gravitational energy of a small surface around the Muon. The code in Python that was used to calculate the result of (43) can be found in Appendix G. To summarize, the factor \(\frac{4}{\pi}\) is possibly due to a field distribution, also see [32] and \(\frac{95}{96}\) is possibly due to Causal Sets [19] and the Steiner Tree Problem [35].

If we choose the \(\frac{1}{1-c} \equiv 44.63955018\) solution, we get in (43) a positive charge of 2.366728973 MeV. This is an interesting result about the Up Quark energy as assessed in lattice QCD, 2.36(24) Mev [36] but the charge of an anti-Muon does not match the Up Quark electric charge which could dismiss this option. This fact is used as self criticism on (43). Another research direction is indirect mass ratios, for example, between particle masses and masses from which other particles are derived. For example

\[
\frac{\text{MuonMass}(1+(a-1)(1-b))}{(1-c)} \equiv \left(1 + \frac{1}{12202.88874}\right) \ast 44.63955018 \ast 105.6583745 \text{ MeV} \cong 4.716928822 \text{ MeV}. This energy, about 4.717 GeV ~ 4.72 GeV [37],[38] appears in QCD as the pole energy of the Bottom Quark and is above the Bottom Quark energy of 4.18 GeV. It is worth mentioning that if we divide the energy of the Muon by 12202.88874066467724, see (42), we get about 8.658472331 KeV which complies with excess of photons with such energy in galactic centers [39]. Another
remark is that the energy of the Tauon 1776.82 MeV or 1776.86 MeV divided by 1220.88874066467724 yields about 145.60650 MeV. This value is within one of the spectral lines of the decay of 99Mo and is used among other lines in nuclear medical imaging. Surprisingly, the sharpest image of lymph nodes in the human body can be seen at 145.6 MeV [40], which means that least scattering of photons occurs at that energy. This phenomenon could have a conventional explanation but it could also indicate the conversion of some of these photons into a neutral particle.

Another interesting energy is: $E_{ElectronMass} \times \sqrt{(a - 1)(1 - b)}$ which is about 4625.8194578 eV.

Refuting a “speculation” claim attack – mathematical coincidence analysis

Assessment of such an attack must be based on probability as dependent on relative error, \( \left( \frac{105.658374524}{0.51099894613} \right) \)

\[ 206.75133988502202 \times 0.5109989461305.658374524 \times 12202.93281199539440 - 12202.8887406646772412202.93281199539440 = \]

\[ \frac{1}{3,379,153,062.118533995049156382052} \approx \frac{0.5109989461 - 0.51099894586371}{0.5109989461} \]

\[ \frac{1}{2,162,566044.04535198211670} \]

where a Muon energy of 105.6583745 MeV and an electron energy of 0.5109989461 MeV were considered.

This result means less than 5/10,000,000,000 relative error! Before trying \( \zeta = \frac{95}{96} \) the author tried \( \zeta = 1 \). So we can say the probability grows to 1/1,000,000,000 but even 1/10,000,000 is already considered a finding. The significance of \( \zeta = \frac{95}{96} \) is of 4 digits and more after the floating point! For example, \( \zeta = \frac{4}{\pi} - 0.0001 \) in the calculation of \( c - 1 \) in (43) yields 0.5110… MeV and \( \zeta = \frac{4}{\pi} + 0.0001 \) yields 0.510969… both further away from the result 0.510998945979878 MeV. 4 digits sensitivity after the point can be seen in the calculation of \((1 + (a - 1)(1 - b))\) in (43) too, \( \zeta = \frac{95}{96} + 0.0001 \) yields 0.51099894008… and \( \zeta = \frac{95}{96} - 0.0001 \) yields 0.5109989518… both further away from the result 0.510998945979878 MeV. These numerical sensitivities highly disfavor a “speculation” claim attack. The following shows how significant is the choice \( \zeta = \frac{95}{96} = 1 - \frac{1}{96} \). We did not refer to annihilation of two Muons. As one quantum system of two entangled particles a Muon and an anti Muon, before annihilation, should be seen as one energy with zero charge. Then (34) turns into

\[ x^2 + \frac{1}{96} \left( -\frac{1}{2} x^2 \pm 0 \cdot \zeta x \right) = x^3 \Rightarrow x^2 + \frac{1}{96} \left( -\frac{1}{2} x^2 \right) = x^3 \]

This equation yields a biggest root smaller than 1. The question we may ask, due to \( \zeta = \frac{95}{96} \) in (43), is: Is there a reasonable \( \zeta = 1 - \frac{1}{n} \) for which \( \frac{1}{1 - x} \approx 206.768277 \) which is the Muon / Electron mass ratio. The closest fit is

\[ \zeta = \frac{23}{24} \Rightarrow x^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{23}{24} \right)^2 \right) = x^3 \Rightarrow \frac{1}{1 - x} \approx 207.0440 \]

(43.2)
\[ 1 - \frac{1}{n} \approx e^{-\frac{1}{n}} = e^{-\frac{1}{24}}, e = 2.71828 \ldots \] and for \( \zeta = e^{-\frac{1}{n}}, \frac{1}{1-x} \approx 206.670967462857 \). The approximation of \( e^{-\frac{1}{n}} \) is \((1 - \frac{1}{n})^{\frac{n}{24}}\), for which the closest \( \frac{1}{1-x} \) to 206.768277 is achieved when \( n = 96 \). This leads to \( \zeta = \left( 1 - \frac{1}{96} \right)^4 \) and (43.1) becomes

\[ x^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{95}{96} \right)^4 \right) = x^3 \Rightarrow \frac{1}{1-x} \approx 206.762203174798 \] (43.3)

The Muon mass is divided to yield the delta \( 2 \cdot \left( \frac{105.658374524 \text{MeV}}{206.7622} \cdots - 0.510998946131 \right) \approx 2 \cdot 15.024999 \cdots \text{eV} \). Other results are for, \( \frac{23}{24} \cdot 207.044017583727 \ldots \), for \( \zeta = \left( \frac{47}{48} \right)^2 \), 206.854768046788 \ldots, for \( \zeta = \left( \frac{191}{192} \right)^8 \), 206.716421179322 \ldots and for \( e^{-\left( -\frac{1}{24} \right)} \), 206.70967462857 \ldots

The relative error \( \frac{206.7682826432 \cdots - 206.762203174798}{206.762203174798} = 34010.914842 \cdots \) along with the 3,379,153,062.118533995 \cdots \) error is of extremely low probability of being a result of pure mathematical coincidence. Empirically, it is interesting to know if the energies 41.875092980 eV, \(-0.00015122\) eV and \(2\cdot15\) eV have a meaning as neutrino energies. \(0.00015122\) eV \~\, \(0.0002\) eV is likely to be the electron neutrino mass.

**Failed research direction - The Higgs Boson**

Consider an electrically neutral particle and Figure 2, where \( U_x \) can, at any given time, point inward or outward in equal probability. Then a gravitational source of this particle’s energy \( E \) can be interpreted from two polynomials

\[ 1 + \frac{1}{96} \left( -\frac{1}{2} \zeta^2 c_1^{-2} + \zeta c_1^{-1} \right) = c_1 \] and \( 1 + \frac{1}{96} \left( -\frac{1}{2} \zeta^2 c_2^{-2} - \zeta c_2^{-1} \right) = c_2 \) and the area addition ratio around some negative charge \( c_1 - 1 \) and area loss ratio around some positive charge \( 1 - c_2 \). Then we can calculate two energies, \( E_1 = \frac{E}{c_1 - 1} \) and \( E_2 = \frac{E}{1 - c_2} \) where the negative sign stands for anti-gravity. Recall our choice of \( \zeta = \frac{4}{\pi} \),

\[
E_2 - E_1 = \frac{E}{c_1 - 1} + \frac{E}{1 - c_2} \approx E \cdot (206.7513398850 \cdots + 44.639550 \cdots ) \approx E \cdot 251.390890060986031074
\] (43.4)

When the energy of the Higgs Boson, \(~125090 \text{ MeV} \) is divided by 251.390890060986031074 we get,

\(~497.5916190 \text{ MeV} \) which is barely within the assessment \(497.611 \pm 0.013 \text{ MeV} \) [41] of the three known light neutral Kaons. This result is a total surprise and if true, there must be anomalies related to the neutral Kaons beyond the Standard Model, however, the uncertainty of the energy in [41] is yet too high to draw conclusions. If it was true then the Higgs boson mass would be dictated by the neutral Kaon mass and we know that the opposite is true.
The W+ Boson and the Z Boson, the Higgs Boson and the anti-Tau particle

The equation

\[ 1 + \frac{1}{96} \left( \frac{4}{\pi} \right)^2 c^{-2} \leq \frac{4}{\pi} c^{-1} \]  

\( = c \) does not take into account the possibility of a null Reeb field

\( \frac{u_k u^k + u^k u_k}{2} = 0 \) or \( U_k U^k \to 0 \) as the radius \( r \to 0 \) or \( U_k U^k = 0 \), with space-like \( p_k \) and norm \( |p_k p^k| \),

also see [10]. In that case, equations (38), (39), (41) become

\[ 1 + \frac{1}{96} \left( \pm \frac{4}{\pi} c^{-1} \right) = 1 + \frac{c^{-1}}{24\pi} = c \]  

(43.5)

and

\[ 1 + \frac{1}{96} \left( \pm \frac{95}{96} b^{-1} \right) = 1 + \frac{95b^{-1}}{96^2} = b \]  

(43.6)

And the biggest roots of these equations are

\[ c_1 = \frac{1 + \left( \frac{1}{6\pi} \right)^2}{2} \approx 1.0130915 \ldots \]  

\[ c_2 = \frac{1 + \left( \frac{1}{6\pi} \right)^2}{2} \approx 0.986556 \ldots \]  

(43.7)

\[ b_1 = \frac{1 + \left( \frac{95}{96+24} \right)^2}{2} \approx 1.010204037 \ldots \]  

\[ b_2 = \frac{1 + \left( \frac{95}{96+24} \right)^2}{2} = \frac{95}{96} \]  

(43.8)

That number is very close to the mass ratio between the Z boson and the W boson, \( \frac{91.1876 \text{ GeV}}{80.370 \text{ GeV}} \approx 1.134597487 \) but the error is too big to rule out mathematical coincidence: \( \text{Error}^{-1} = \frac{1.1334579487}{1.134470446} \approx 0.9814.16455 \) and with another W Boson assessment \( \frac{91.1876 \text{ GeV}}{80.379 \text{ GeV}} \approx 1.134470446 \) we get \( \text{Error}^{-1} = \approx 10442.60078 \). It is interesting, but insufficient to rule out mathematical coincidence. Nevertheless, (43.7) means that the coefficient \( \frac{95}{96} \) is related to charge-less particles and \( \frac{4}{\pi} \) to electrically charged particles. Similar to (43), if we multiply, \( 1.134361808 \cdot (1 + (c_1 - 1)(1 - c_2)) \) \( \approx 1.134561453 \) then the relative error for \( \frac{91.1876 \text{ GeV}}{80.370 \text{ GeV}} \approx 1.134597487 \) is about \( \frac{1}{31486.95424} \), i.e.

\[ (1 + (c_1 - 1)(1 - c_2)) \left( \frac{(c_1 - 1)(1 - c_2)}{(b_1 - 1)(1 - b_2)} \right)^{1/4} \approx 1.134561453 \]  

(43.7.1)

and for W Boson of 80.3725 GeV the error is about \( \frac{1}{1528961.689} \), however, in (43) it made sense due to the electron neutrino or another type of neutrino. In this case, this \( (1 + (c_1 - 1)(1 - c_2)) \) calibration
requires a more comprehensible theory and is therefore in tension with the assumption of a W boson energy of 80.3725 GeV. The roots in (43.7) should have a meaning when the particle has spin 1 and not spin $\frac{1}{2}$.

**The Higgs Boson**

A slightly different Higgs energy of 125.35 GeV yields a very interesting result with $\frac{1}{\sqrt{(b_1-1)(1-b_2)}} \approx 96.99505572$. 125.35 GeV $\cdot \sqrt{(b_1-1)(1-b_2)} \approx 1.292333914$ GeV. Divide this value by the mass of the neutron, $\frac{125.35 \text{ GeV}}{0.9395654133 \text{ GeV}} \cdot \sqrt{(b_1-1)(1-b_2)} \approx 1.375459224$ this becomes interesting because the mass ratio between the Higgs Boson and the Z Boson is about $\frac{125.35 \text{ GeV}}{91.1876 \text{ GeV}} \approx 1.372774368$. If we check the relative error we get $\text{Error}^{-1} \approx \frac{1.374638657}{1.375459224} \approx 1675.230712$. That is a low error. Not enough yet to rule out mathematical coincidence but indeed low. The energy 1.292333914 GeV could be an indication of a new Boson or a vector Boson of about 1.29 GeV but not an isovector resonance [42] and evidence for its existence should be searched for in particle accelerators. Another research direction was to use the inverted value of $\xi = \frac{4}{\pi}$ i.e. $\xi = \frac{\pi}{4}$ in the negative and positive charge area ratio equations. That yields two new maximal roots $a_1^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{\pi}{4} \right)^2 + \frac{\pi}{4} a_1 \right) = a_1^3$ and $a_2^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{\pi}{4} \right)^2 - \frac{\pi}{4} a_2 \right) = a_2^3$ along with the older ones $b_1^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{4}{\pi} \right)^2 + \frac{4}{\pi} b_1 \right) = b_1^3$ and $b_2^2 + \frac{1}{96} \left( -\frac{1}{2} \left( \frac{4}{\pi} \right)^2 - \frac{4}{\pi} b_2 \right) = b_2^3$.

Quite like the ratio in (43.7), we have, $\sqrt{\frac{(b_1-1)(1-b_2)}{(a_1-1)(1-a_2)}} \approx \sqrt{\frac{201.6240447 \times 86.46523917}{206.7513399 \times 44.63955018}} \approx 1.374383282$ which is close to the following mass ratio between a Higgs Boson of 125.3267 GeV and a Z Boson of 91.1876 GeV which yields, 1.37438314, close to 1.374383282. It is interesting though not sufficiently accurate to draw any conclusion at this stage. The idea behind using charge equations without null Reeb vectors is because the Higgs boson is supposedly responsible for non zero mass. From (43.7.1) and using s instead of b, $\sqrt{(s_1-1)(1-s_2)} \approx 75.3783115$ and 91.1876 GeV $\ast \sqrt{\frac{(b_1-1)(1-b_2)}{(a_1-1)(1-a_2)}} \ast (1 + (s_1 - 1)(1 - s_2)) \approx 125.3487702$ GeV. A similar $(1 + (s_1 - 1)(1 - s_2))$ value was used in (43.7.1).

**The Tau Lepton**

What doesn’t seem right is the use of null Reeb vectors, which may not be even possible in order to reach a relation between particles that have rest mass, W and Z bosons. In this manner if we multiply the W boson mass by $\sqrt[4]{(c_1-1)(1-c_2)}$ and the Z boson mass by $\sqrt[4]{(b_1-1)(1-b_2)}$ we get almost the same value, about 9.26 GeV. Unfortunately there is no such boson and thus, this is a spurious prediction. The value 1.134361808 gives us some hope to find other particles from the W and Z boson masses from the ratios in (43.7). The geometric average of these values is $\frac{1}{\sqrt{(c_1-1)(1-c_2)}} \approx 75.3783115$ and the ordinary average $\frac{(c_1-1)^{-1} + (1-c_2)^{-1}}{2} \approx 75.384949$. Do these values have a physical meaning? They could have if the Anti-Tau
particle can be derived from the gravitational energy of the W Boson. Now returning to (43) and to after (43), the reason for not considering the negative charge is simply because in the real case, where $U^\mu U^\nu = 0$, $p_\mu$ is space-like and that means that negative and positive charge should be considered with opposite signs, so indeed, it is a deep consideration. this portion from the energy of the $W^+$ Boson [43] yields, \( \frac{80370}{44.639550} \text{MeV} \approx 1800.4213681 \text{MeV} \). The delta between this value and the Anti-Tau energy 1776.86 MeV is \( \sim 23.5613681 \text{MeV} \). A higher delta is obtained if the Tau energy is taken to be, 1776.82 MeV, \( \sim 23.6013681 \text{MeV} \). From (43.7) we can see that \( \frac{1776.86}{75.3783115} \approx 23.5725630 \text{MeV} \) which means that a good approximation of the mass of the Tau particle can be obtained from \( \sqrt{(c_1 - 1)(1 - c_2)} \approx 75.3783115^{-1} \)

\[
\frac{80370(1-c) \text{MeV}}{1+\sqrt{(c_1-1)(1-c_2)}} \approx 1776.86 \text{ MeV} \quad \text{or} \quad \frac{80370(1-c) \text{MeV}}{1+\sqrt{(c_1-1)(1-c_2)}} \approx 1776.86 \text{ MeV}
\]

(43.8)

The root, \( \sqrt{(c_1 - 1)(1 - c_2)} \) will be better understood in (93) as a result of taking the root of a determinant of a Gram matrix in (93). And with a more accurate Python code and 80369 MeV $W^+$ Boson we get 1776.82684328632649 MeV and with a W Boson of 80372.88 MeV the Tau energy would be 1776.91 MeV. Unlike in (43) the root should be related to spin 1 and not to spin \( \frac{1}{2} \). There are two big problems unfortunately, first, (43.8) cannot be true because a null Reeb vector would imply that the portion of energy 23.5613681 MeV would not have rest mass. If it is a new type of neutrino, then this particle must have rest mass. Second, the averaging in the denominator of (43.7),(43.8) also implies a zero charge oscillating field. The $W^+$ Boson has a rest mass and is not a composite particle and therefore the 23.5613681 MeV cannot be a part of the $W^+$ Boson. Also, the $W^+$ Boson has rest mass which immediately invalidates any claim that it is related to a null Reeb vector. These considerations unfortunately invalidate the feasibility of (43.5), (43.6), (43.7) as a theory without any modification. Weak evidence for (43.7) can be [44]. Other values for \( \zeta = 1 - \frac{1}{96} = \frac{95}{96} \) in the null Reeb vector equations, from which

\[
\sqrt{(b_1 - 1)(1 - b_2)} \approx 96.9950557^{-1}
\]

(43.9)

This portion from the $Z$ boson, \( \frac{91187.6 \text{MeV}}{96.99505572...} \approx 940.12627 \text{ MeV} \), is close to the neutron energy, 939.5654133 MeV. 939.5654133 MeV * 96.9950557... yields \( \sim 91.131996 \text{ GeV} \) which can be viewed as a an intermediate value $Z'$ which is smaller than the energy of the Z boson 91.1876 GeV. (43.9) offers a different process than (43.4), by which mass ratios of neutral particles can be obtained. The $Z'$ intermediate vector boson energy of 91.133 GeV can be found in [45]. (43.9) is problematic because it requires the use of null Reeb vectors which would imply in this case, that a large portion of the Z boson should be mass-less. Obviously this is not true and that is why (43.9) is not considered as a feasible theory. If it was true then the $Z'$ energy and therefore the Z boson energy would be dictated by the neutron’s mass. Nevertheless, a researcher’s integrity obliges to account for not only successes but also for directions that turned out to be wrong.

**The Tauon and the Muon mass ratio**

**Note:** The more advanced parts of this section require basic knowledge of electrical engineering and especially a good understanding of the trivial subject of Dissipation Factor and Loss Tangent.

The denominator \( 1 + \sqrt{(c_1 - 1)(1 - c_2)} \) in (43.8) and \( 1 + (\alpha - 1)(1 - b) \) in (43) can be used together to yield a nice result that seems to be more than just a mathematical coincidence.

Consider the following:
Such that $(g_1 - 1)^{-\frac{1}{2}} = \frac{1}{2} (1 - g_2)^{-1}$ \hfill (43.10)

With biggest roots $g_1 \approx 1.003629541$ and $g_2 \approx 0.969877163$. $g_1$ means an area portion $275.51693^{-1}$ is added around a negative charge and $33.19740^{-1}$ of the area is subtracted around a positive charge, which reflects a possibly maximal allowed gravitational imbalance between negative and positive charge.

A calculation that uses an electronic datasheet yields,

$$\xi \approx 1.55619853719035, \quad (g_1 - 1)^{-\frac{1}{2}} \approx 16.59870203 \hfill (43.11)$$

which is close to the known mass ratio between the Tauon and the Muon, $\approx 16.817$. Multiplying this value by $1 + \sqrt{(c_1 - 1)(1 - c_2)}$ from (43.8) and dividing by $(1 + (a - 1)(1 - b))$ from (43) yields,

$$((g_1 - 1)^{-\frac{1}{2}})^{1 + \sqrt{(c_1 - 1)(1 - c_2)}} \left(1 + (a - 1)(1 - b)\right) \approx \frac{\text{Tauon 1776.9127923826 MeV}}{\text{Muon 105.6583745 MeV}} \hfill (43.12)$$

Which is $\approx 16.81752914$. So this calculation predicts a Tauon energy of about $1776.9127923826$ MeV from the Muon energy. There are some major caveats before getting excited about this result. Unlike (38), (43) in which $\xi = \frac{4}{\pi}$ came directly from Ettore Majorana’s notebook [32], $\xi \approx 1.55619853719035$ does not seem to have a straightforward geometric meaning. A second caveat is that $1776.9127923826$ MeV is not based on any accurate measurement of the Tau particle’s energy $1776.91$ [46] and is most likely wrong! This assessment is indeed interesting, but it is not as nearly as neat and as accurate as (43). From the [46] Tauon energy result, a prediction of the W and Z masses yields: Tauon energy, $1776.9127923825637936$ MeV, W energy, $80372.8876286350568989$ MeV, Z energy, $91187.9801482948387275$ MeV from $\xi \approx 1.5561985371903484$ and a more accurate ratio than (43.7.1), $1.1345614527330559$. The most remarkable property of the value $16.59870203$ in (43.11) is the following property, $2\left(\frac{\pi}{2} - \xi\right)^{-\frac{1}{2}} \approx 16.55337088$ which means that this expression approximates the value $16.59870203$ from (43.11). Even more interesting is another property, $2\left(\frac{\pi}{2} - \xi\right)^{-1} \approx 137.0070438$ which is close to the inverse of the Fine Structure Constant, $137.035999046$. We will see more about this relation in the following section.

An unexpected relation to the inverse of the Fine Structure Constant

If we check the following with the same $\xi$ from (43.11), $g_2$ from (43.10), we get,

$$\frac{2}{\cos(\xi)} \approx \frac{2}{\cos(1.5561985371903484)} \approx 137.011909869, \hfill (43.13)$$
\[ \tan^{-1}(95^296^2(1 - g_2)^4) \approx 1.5561948778250207190765973767615 \] (43.13.1)

which remarkably approximates \( \xi \approx 1.5561985371903484 \) from (43.10), (43.11).

\[
\text{Error} = \frac{\xi - (95^296^2(1-g_2)^4)}{\xi} \approx 425,263.60132816790517958824157133^{-1} \] (43.13.2)

In terms of electrical engineering Dissipation Factor and Loss Tangent, we can write, 
\[
DF = \frac{95^296^2}{(1-g_2)^4} \approx \tan(\xi) \]

where the numerator is known as the Resistive Power Loss and the denominator as the Reactive Power Oscillation. It is expected that an oscillating charge will generate oscillation in area due gravity changes, however, it is not expected that the area portion that is lost due to gravity will appear as the power of 4. This is a very rare property that connects between trigonometry and the electro-gravity polynomials (43.10). We can get from this relation two insights, the first is that if (43.13) is not a mathematical coincidence, then the inverse Fine Structure constant should come out of a trigonometric function and a numbers relation. The second is that \(95^296^2(1 - g_2)^4\) should be part of this equation. We may think that perhaps scaling of the value of \( \xi \) in a rational way, will yield the exact inverse Fine Structure Constant. So we want to find some \( d \) such that

\[
\frac{2}{\cos\left(\frac{15561985371903484*(1+\frac{1}{d})}{\cos\left(\frac{15561985371903484*(1+\frac{1}{d})}{2}\right)}\right)} \text{ will yield the constant we are looking for. We will soon find such } d,
\]

\[ d \approx 606400.8 \text{ that complies with [47] and we get, } \frac{2}{\cos\left(\frac{15561985371903484*(1+\frac{1}{606400.8})}{2}\right)} \approx 137.0359990462475253. \]

Until now, that is not very interesting because we could not find \( d \) out of any new theory. Well, not very accurate.

First, \( \frac{1}{2}(1 - g_2)^{-4} \approx 607276.5368006824282929 \approx 606400.8 \) and

\[
\frac{2}{\cos\left(\frac{15561985371903484*(1+2(1-g_2)^4))}{2}\right)} \approx 137.0359643018112763
\]

If we test the following values for \( d \approx 606400.8 \) we get; \( \frac{95^4}{d} \approx 134.3181357940161.., \frac{96^4}{d} \approx 140.0635619214.. \)

and the geometric average of these two values is \( \left(\frac{95^4}{d} \frac{96^4}{d}\right)^\frac{1}{2} = \frac{95^296^2}{d} \approx \frac{137.1607689}{d} \) so we may try our luck with each one of the following equations:

\[
\frac{2}{\cos\left(\xi*(1+\frac{1}{d})\right)} = \frac{95^296^2}{d} \quad \text{or} \quad \frac{2}{\cos\left(\xi*(1+\left(\frac{96^2}{95}\right)^\frac{1}{2}d\right)} = \frac{96^4}{d}\]

(43.14)

The idea to use 95 and 96 is from (41), \( 1 - \frac{1}{96} = \frac{95}{96} \). This idea does not lead to the exact inverse Fine Structure Constant but very close to it, 137.03597712551937661 with some numerical error,

\[
\frac{95^296^2}{d} = \frac{2}{\cos\left(\xi*(1+\frac{1}{d})\right)} \approx 0.00000000000028422 \quad (43.15)
\]
Each one of the two equations of (43.14) yields,

\[
\frac{2}{\cos\left(\xi(1+\frac{1}{d})\right)} \cong 137.035977125551937661
\]  

(43.16)

**An exact Inverse Fine Structure Constant**

As a result of the conclusions of (43.13.1), (43.13.2), the exact Fine Structure Constant was found by the following, although some aspects of the following calculation are not resolved yet. We put together (38), (39), (43.10), (43.11), (43.13), \[
\frac{1}{2}(1 - g_2)^{-4} \cong 607276.5368006824282929,
\]

\[
1 + \frac{1}{96}\left(-\frac{1}{2}\left(\frac{4}{\pi}\right)^2 a^{-2} + \frac{4}{\pi} a^{-1}\right) = a
\]

\[
1 + \frac{1}{96}\left(-\frac{1}{2}\left(\frac{4}{\pi}\right)^2 b^{-2} - \frac{4}{\pi} b^{-1}\right) = b
\]

\[
d = \left(\frac{1}{2}(1 - g_2)^{-4}\right)^{1+(a-1)(1-b)} \cong 606401.0372 \cong 606400.8
\]

\[
\frac{2}{\cos(1.5561985371903484*(1+\frac{1}{d}))} \cong 137.0359990368270076
\]  

(43.17)

A more exact term in (43.17) \(d = 606401.0371724886354059\) when multiplied by \(1 - \frac{1}{137.0359990368270076}\) yields 606401.0371724886354059(1 + \(1 + \frac{1}{137.0359990368270076}\)) \(\cong 606,400.80152883067046189980208498\). Now recall that \[
\frac{2}{\cos(1.5561985371903484+(1+\frac{1}{606400.8}))} \cong 137.035999046.
\]

So the previous calculation represents the error in \(d\) in an elegant expression that is easy to remember.

Another result is by finding the variable \(s\) where \(a\) and \(b\) are given in (43.17):

\[
\left(\frac{95^2*96^2}{s}\right)^{1+(a-1)(1-b)} = \frac{2}{\cos\left(\xi\left(1+\frac{1}{s(1+(a-1)(1-b))}\right)\right)} \Rightarrow
\]

(43.17.1)

\[
\left(\frac{95^2*96^2}{s}\right)^{1+(a-1)(1-b)} \cong 137.035999036428876252
\]

Another idea is to solve the following equation where \(s\) is given and \(p\) is a variable:

\[
s = (\frac{1}{2}(1 - g_2)^{-4}) \cong 607276.536800682428292930
\]
Combining (43.17.1) and (43.17.2) we find a numerical attractor at (43.17.2) with

\[
\cos \left( \frac{1}{(1 + \frac{1}{(1 + 1/(p+p)})^2)} \right) = \frac{2}{s^{1+1/(p+p)}} \Rightarrow
\]

\[
\left( \frac{95^2 * 96^2}{s} \right)^{1+1/(p*p)} \approx 137.035999035747181551, p
\]

\[
\approx 96.070666670305840285
\]

Combining (43.17.1) and (43.17.2) we find a numerical attractor at (43.17.2) with

\[
s \approx 607276.5368006824282929301262 \approx \frac{1}{2} (1 - g_2)^{-4}, s \left( \frac{1}{1+1/(p*p)} \right) \approx
\]

\[
606401.06429681263938791, \xi \equiv 1.5561985371903484, \text{ which is surprising. For (43.17) – (43.17.2) see Appendix I. Before we close this discussion, it is nice to mention another}
\]

\[
\text{relation } \left( 1 - \ln \left( \frac{1}{\frac{1}{137.035999035747181551}} \right) \right)^{-1} \approx
\]

\[
275.4045237287 \approx 275.51693 \approx (g_1 - 1)^{-1} \text{ in (43.10). That is not a total surprise because}
\]

\[
(1 - \frac{1}{Z})^{-1} \approx 2z \text{ for big } z.
\]

\[
\text{Tauon to Muon mathematical coincidence analysis}
\]

With \( \xi \equiv 1.5561985371903484 - 0.001 \), from (43.11), the Tau energy is 1775.5044808249303969 MeV and with

\[
\xi \equiv 1.5561985371903484 + 0.001 \text{ the Tau energy is } 1778.3270829769808188 \text{ MeV. The sensitivity to a delta}
\]

\[
0.0001 \text{ is } \xi \equiv 1.5561985371903484 \text{ - 0.0001, Tau energy } 1776.7716927033129650 \text{ MeV and}
\]

\[
\xi \equiv 1.5561985371903484 + 0.0001, \text{ Tau energy } 1777.0539518519776720 \text{ MeV. So for an energy of around 1776.9 MeV or even 1776.86 MeV, a speculation claim attack is disfavored.}
\]

\[
\text{Particle mass ratios by added or subtracted area – Approximation of the Fine Structure constant}
\]

Like the previous section, this section relies on area ratios to represent mass ratios. Unlike (43.17.2) it is brought here for showing another possible approximation of the inverse Fine Structure Constant but not as the process that yields that constant. Some of the work can be seen in the remarkable work of Lee C. Loveridge \([28]\) and is somewhat speculative unless (34.4) or claims of a ~40eV neutrino \([29]\) and other quantum effects \([34]\) are taken into account and therefore the reader is honestly advised to take it with a grain of salt, however, the results are very interesting and do have geometrical basis in \([12]\) that worth further investigation - the Reeb vector on the foliation leaves must have drains and sources. It is therefore the authors opinion that the technique used here will be considered by the reader even if the reader does not fully agree with the idea which is presented in this section. An approximation to the Fine Structure constant was discovered totally by chance when the author calculated the geometric average of area ratios quite similar to the square root of one that appear in the denominator of (43). i.e.

\[
given \text{two polynomials, } 1 + \frac{1}{96} \left( \frac{1}{2} \zeta^2 a^{-2} + \zeta a^{-1} \right) = a \text{ and } 1 + \frac{1}{96} \left( \frac{1}{2} \zeta^2 b^{-2} - \zeta b^{-1} \right) = b, \text{ then we are interested in exploring } \sqrt{(a - 1)(1 - b)}. \text{ The question that was explored was, if } \zeta \text{ is the inverse of the average distance between two points on a sphere S2 then what is } \sqrt{(a - 1)(1 - b)}? \text{ This question is hard to solve because in (31) and in (32), we were not interested in the geometry within the infinitesimal sphere S2 but only}
\]
in area deviations from the surface area \(4\pi r_0^2\), while ignoring the interior of the sphere, which is a three dimensional ball. The greater problem is that in the calculation of the inverse of the average distance, \(\zeta\) is dependent upon \(x\) in the solution of 

\[1 + \frac{1}{96}\left(-\frac{1}{2} \xi^2 x^{-2} + \xi x^{-1}\right) = x\]

because if the area \(4\pi r^2\) becomes \(4\pi r^2 x\) for every radius \(r < r_0\) then the average distance between two points on the sphere, depends on \(rx\). We must therefore calculate \(\zeta = f(x)\) for some function \(f(x)\). In flat geometry, the Riemannian distance between two points on the sphere that have an angle \(\theta\) with the center is

\[S(\theta) = 2r_0 \sin \left(\frac{\theta}{2}\right)\]  \(\text{(44)}\)

Integrating on the sphere and dividing by the area, we get the average distance in flat geometry, \(\frac{1}{4\pi r_0^2} \int_0^\pi S(\theta) 2\pi r_0 \sin(\theta) \, d\theta = \frac{4}{3} r_0\) \(\text{(45)}\)

So if the acceleration field depends on the inverse of that distance, it will depend on \(\frac{3}{4} r_0\) and \(\zeta = \frac{3}{4}\). Plugging this value in (34) and calculating the biggest roots \(a, b\) of the two resulting polynomials,

\[\left(\frac{1}{(1-a)} - \frac{1}{(b-1)}\right)^{1/2} \approx 137.2504256\]  \(\text{(46)}\)

Close to 137.035999173 which is the inverse of the fine structure constant but not good enough.

For the sake of simplicity, we assume \(r_0 = 1\). We need to take into account the effect of dilation and contraction of areas around the electric charge in order to get a closer value. A Riemannian geodesic curve distance \(S(\theta)\) in (44) is replaced by

\[S(\theta, x) = \text{Min} \int_{\alpha=0}^{\alpha=\theta} \left(\sqrt{R(r, x) \cdot r^2 + \left(\frac{d r}{d\alpha}\right)^2}\right) \, d\alpha\]  \(\text{(47)}\)

With boundary conditions \(r(0) = 1, r(\theta) = 1, r_0 = 1\), such that \(x\) is the biggest root of one of the two polynomials (34), \(x = 1 + \frac{\text{Added or subtracted disk area}}{4\pi r^2\text{ instead of }\pi r^2}\), \(x < 1\) or \(x > 1\) and \(x > 0\). \(R(r, x) \cdot r^2 \, d\alpha^2\) is the square length element in a direction perpendicular to the radius \(r\). This is the infinitesimal distance component that is influenced by the ball surface dilation at radius \(r\). \(R(r, x) \cdot r^2 \, d\alpha\) is the infinitesimal square distance component along the radius and is therefore not influenced by surface dilation or contraction by \(x\). The simplest model in which the area dilation is \(x\) at radius \(r_0\) and 1 in the center, which means the curvature is 0 at the center, is \(R(r, x) = 1 + (x - 1) \cdot r\) so \(r = 0 \Rightarrow R(r, x) = 1, \ r = 1 \Rightarrow R(r, x) = x\). See Fig. 3.0

\(\text{(Fig. 3.0)}\) – The reduced 3 dimensional curvature, from Minkowsky 4D to Riemann 3D around a hollowed positive charge whose field depends on average distance between two points on the sphere.

\[R(r, x) = 1 + (x - 1) \cdot r_0 = x\]

\[r_0 = 1\]

\[\text{Center of re charge } R(r, x) = 1\]

(47) becomes,

\[S(\theta, x) = \text{Min} \int_{\alpha=0}^{\alpha=\theta} \sqrt{((x - 1)r + 1) \cdot r^2 + \left(\frac{d r}{d\alpha}\right)^2} \, d\alpha \text{ and } r(0) = r(\theta) = 1\]  \(\text{(48)}\)
Then (45) becomes,

\[
\frac{1}{\xi} = D(x) = \frac{1}{4\pi} \int_0^\pi xS(\theta, x)2\pi \sin(\theta) \, d\theta = \frac{1}{4\pi} \int_0^\pi S(\theta, x)2\pi \sin(\theta) \, d\theta
\]  

(49)

(48) and (49) turned out to be extremely difficult to calculate. The author had to compromise on accuracy, and chose integration along a straight line in the coordinate system as a compromised approximation of (48) but with a systematic error. The term

\[
2 \cos \left( \frac{\theta}{2} \right) \int_{\alpha=0}^{\alpha=\frac{\theta}{2}} \frac{\sqrt{\sin^2(\alpha)+R(r,x)\cos^2(\alpha)}}{\cos^2(\alpha)} \, d\alpha
\]

is reduced to \(2 \sin \left( \frac{\theta}{2} \right)\) as in (44), with \(r_0 = 1\) if \(R(r,x) = 1\) for all \(0 < r \leq 1\). So as an approximation with a systematic error we have,

\[
S(\theta, x) \approx 2 \cos \left( \frac{\theta}{2} \right) \int_{\alpha=0}^{\alpha=\frac{\theta}{2}} \sqrt{1 + (R(r,x) - 1) \cdot \cos^2(\alpha)} \cdot \frac{\cos(\frac{\theta}{2})}{\cos^2(\alpha)} \, d\alpha
\]

and instead of the exact (49) we use,

\[
\frac{1}{\xi} = D(x) \approx \frac{1}{4\pi} \int_{\theta=0}^{\theta=\pi} 2\pi \sin(\theta) \cdot 2 \cos \left( \frac{\theta}{2} \right) \left( \int_{\alpha=0}^{\alpha=\frac{\theta}{2}} \sqrt{1 + (R(r,x) - 1) \cdot \cos^2(\alpha) - \cos(\frac{\theta}{2}) \cdot \cos(\alpha)} \, d\alpha \right) \, d\theta
\]  

(50)

We are now ready to run a computer calculation in C++, with a small systematic error, which is very slow but at least works,

\[
a^2 + \frac{1}{96} \left( -\frac{1}{2} D(a)^{-2} + D(a)^{-1} a \right) = a^3
\]  

(52)

\[
b^2 + \frac{1}{96} \left( -\frac{1}{2} D(b)^{-2} - D(b)^{-1} b \right) = b^3
\]  

(53)

And the approximated \(D(x)\) yields a good result despite the systematic error in (50),(51), instead of using the exact terms (48),(49).

\[
\left( \frac{1}{1-a} \right)^{1/2} \left( \frac{1}{b-1} \right)^{1/2} \approx 137.041002618
\]

(54)

The code that was used to calculate (53) was written in C++ and can be found in Appendix H. Integral (50) is further developed in order to avoid the use of time consuming functions.

**Conclusion**

Using time-like curves which are based on Morse functions, it is possible to describe fields of acceleration that are not predicted by the metric alone. The acceleration fields have a rigid mathematical foundation in the work of Georges Reeb from 1948 and in the theory of foliations. It is possible to say that space-time codim-1 foliations represent geometric information that is not represented by the metric tensor alone. Although a lot of work has to be done in order to show how to reconstruct the results of Quantum Field Theory, the results of this paper should raise a new interest in this work which cannot progress further as a work of one man. Unexpected gravity by electric charge has an immense importance to the development of the human race and it is especially important in order to understand the Dark Matter effect and in order to develop feasible Alcubierre · White or Alcubierre Froning Warp Drive technology. At least part of the Dark Matter effect may not be due to Dark Particles.

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Appendix A: Euler Lagrange minimum action equations

We assume \( \sigma = 8\pi \) (from the previously discussed term, \(-a\mu a^\mu / 8\pi K\) as an energy density).

\[
Z = N^2 = P_\mu P^\mu \quad \text{and} \quad U_\lambda = \frac{Z_\lambda}{Z} - \frac{Z_k P^k P_\lambda}{Z^2} \quad \text{and} \quad L = \frac{1}{4} U^k U_k
\]

\( R = \text{Ricci curvature} \).

\[
\text{Min Action} = \text{Min} \int_\Omega \left( R - \frac{8\pi}{\sigma} L \right) \sqrt{-g} \, d\Omega = \text{Min} \int_\Omega \left( R - \frac{1}{4} U^k U_k \right) \sqrt{-g} \, d\Omega \quad \text{s.t.} \quad \sigma = 8\pi
\]

(55)

The variation of the Ricci scalar is well known. It uses the Platini identity and Stokes theorem to calculate the variation of the Ricci curvature and reaches the Einstein tensor [48], as follows,

\[
\delta R = R_{\mu\nu} \delta g^{\mu\nu} \quad \text{and} \quad \delta \sqrt{-g} = -\frac{1}{2} g_{\mu\nu} \delta g^{\mu\nu} \sqrt{-g} \quad \text{by which we infer} \quad \delta (R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}) = (R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}) \delta g^{\mu\nu}
\]

which will be later added to the variation of \( \left( \frac{1}{2} R - \frac{8\pi}{\sigma} L \right) \sqrt{-g} \) by \( \delta g^{\mu\nu} \).

The following Euler Lagrange equations have to hold,

\[
\left( \frac{\partial}{\partial g^{\mu\nu}} - \frac{d}{dx^m} \frac{\partial}{\partial (g^{\mu\nu}_{m,r})} \right) \left( \frac{1}{2} R - \frac{1}{4} U^k U_k \right) \sqrt{-g} = 0
\]

and

\[
\left( \frac{\partial}{\partial p^m} - \frac{d}{dx^m} \frac{\partial}{\partial (p^m_{r,s})} \right) \left( \frac{1}{2} R - \frac{1}{4} U^k U_k \right) \sqrt{-g} = 0
\]

\[
U^k U_k = \frac{Z_k Z^k}{Z^2} - \frac{(Z_k P^k)^2}{Z^2} \quad \text{which we obtain from the minimum Euler Lagrange equation because}
\]

\[
U_\lambda P^\lambda = \frac{Z_k P^k}{Z} - \frac{Z_k P^k P^\lambda}{Z^2} = 0
\]

In order to calculate the minimum action Euler-Lagrange equations, we will
separately treat the Lagrangians, \( L = \frac{Z \mu Z^\mu}{Z^3} \) and \( L = \left( \frac{Z \mu P^\mu}{Z^3} \right)^2 \) to derive the Euler Lagrange equations of the Lagrangian \( L = \frac{Z \mu Z^\mu}{Z^3} - \left( \frac{Z \mu P^\mu}{Z^3} \right)^2 = U \mu^ \mu \). The Euler Lagrange operator of the Ricci scalar 

\[
\frac{\partial}{\partial g^{\mu \nu}} - \frac{d}{dx^m} \frac{\partial}{\partial (g^{\mu \nu},_m)} + \frac{d^2}{dx^m dx^n} \frac{\partial}{\partial (g^{\mu \nu},_{mn})}.
\]

The reader may skip the following equations up to equation (61). Equations (61), (62) and (63) are however crucial.

\[
L = \left( \frac{P \mu P^\mu}{Z^3} \right)^2 \quad \text{s.t.} \quad Z = P \mu P^\mu \quad \text{and} \quad Z_s \equiv Z_s = \frac{dZ}{dx^s}
\]

\[
\frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu}} - \frac{d}{dx^m} \frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu},_m} = \left( -2 \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu P^m + \Gamma^i_{\mu \nu} P_i P^\mu P^m + \Gamma^i_{\mu \nu} P_i P^m \right) P \mu P^\mu - 2 \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu
\]

\[
\left( -2 \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - 2 \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu - \left( \frac{Z_s P^s}{Z^3} \right) P \mu P^\mu \right) \sqrt{-g} = X (56)
\]

\[
L = \frac{Z^2 Z_s}{Z^2} \quad \text{s.t.} \quad Z = P \mu P^\mu, \quad \text{s.t.} \quad Z = P \mu P^\mu \quad \text{and} \quad Z_s \equiv Z_s = \frac{dZ}{dx^s}
\]

\[
\frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu}} - \frac{d}{dx^m} \frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu},_m} = \left( -2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu P^m + \Gamma^i_{\mu \nu} P_i P^m + \Gamma^i_{\mu \nu} P_i P^m \right) P \mu P^\mu - 2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu
\]

\[
2(\Gamma^i_{\mu \nu} P_i P^m + \Gamma^i_{\mu \nu} P_i P^m) - 2Z \mu P^\mu - 2Z \mu P^\mu - 2Z \mu P^\mu - 2Z \mu P^\mu - 2Z \mu P^\mu = X (57)
\]

We subtract (56) from (57)

\[
Z = P \mu P^\mu, \quad \text{s.t.} \quad Z = P \mu P^\mu \quad \text{and} \quad Z_s \equiv Z_s = \frac{dZ}{dx^s}, \quad U_k = \frac{Z^2}{Z^3}, \quad L = U_k U_k = \frac{Z^2}{Z^3}
\]

\[
\frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu}} - \frac{d}{dx^m} \frac{\partial(L\sqrt{-g})}{\partial g^{\mu \nu},_m} = \left( +2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu + 2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu + 2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu + 2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu + 2 \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu
\]

\[
+ \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu - \left( \frac{Z^2 P^m}{Z^3} \right) P \mu P^\mu
\]
\[
\left(U^\mu U_\nu - \frac{1}{2} U^A U_B g_{\mu\nu} - 2 U^k \frac{\partial p_k}{\partial z}\right) \sqrt{-g} 
\]
(58)

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

\[\]

We subtracted the Euler Lagrange operators of \(\frac{(Z^i P^i)^2}{Z^3}\sqrt{-g}\) in (56) from the Euler Lagrange operators of \(\frac{Z^i Z^i}{Z^3} \sqrt{-g}\) in (57) and got (58) and we will subtract (59) from (60) to get two tensor equations of gravity, these will be (61), and (63). Assuming \(\sigma = 8\pi\), where the metric variation equations (55), (56), (57) and (58) yield
\[ Z = N^2 = P^\mu P^\nu, \quad U_\lambda = \frac{Z_\lambda}{Z} - \frac{Z_\lambda P^k P_\lambda}{Z^2}, \quad L = \frac{1}{4} U_i U^i \quad \text{and} \quad Z = P^k P_k \]
\[
\left\{ \begin{array}{l}
8\pi \frac{1}{\sigma} \frac{1}{4} U_\mu U_\nu - \frac{1}{2} U_k U^k g_{\mu \nu} - 2U^k \frac{P_\mu P_\nu}{Z} = R_{\mu \nu} - \frac{1}{2} R g_{\mu \nu} \\
\text{s.t.} \quad R = R_{\mu \nu} g^{\mu \nu} \\
\text{s.t.} \quad Rkj = (\Gamma j^\mu_{\mu^\mu})_{,\mu} - (\Gamma j^\mu_{\mu^\mu})_{,\mu^\mu} + \Gamma j^\mu_{\mu^\mu} \Gamma j^\mu_{\mu^\mu} - \Gamma j^\mu_{\mu^\mu} \Gamma j^\mu_{\mu^\mu}
\end{array} \right. 
\]

(61)

\( R_{\mu \nu} \) is the Ricci tensor and \( R_{\mu \nu} \frac{1}{2} R g_{\mu \nu} \) is the Einstein tensor [48]. In general, by (28) and \( \sigma = 8\pi \), (61) can be written in \((-1, +1, +1, +1)\) metric convention, so \( Z = |P^\mu P^\mu| \) as,

\[
\frac{1}{4} (U_\mu U_\nu - \frac{1}{2} U_k U^k g_{\mu \nu} - 2U^k \frac{P_\mu P_\nu}{Z}) = R_{\mu \nu} - \frac{1}{2} R g_{\mu \nu}
\]

(62)

**Charge-less field:** The term \(-2U^k \frac{P^\mu P^\nu}{Z}\) in (62) can be generalized to:

\[-2((U^k \frac{1}{2} + U^* k \frac{1}{2}) \frac{P^\mu P^* \nu}{Z})/2 \] and can be zero under the following condition:

\[
4(A_{\mu \nu} \frac{P^\nu}{\sqrt{Z}} + A^*_{\mu \nu} \frac{P^\nu}{\sqrt{Z}}) = U_\mu U^* \mu + U^* \mu U^\mu \Rightarrow U^k \frac{1}{2} + U^* k \frac{1}{2} = 0
\]

**Note:** The complimentary matrix \( B_{\mu \nu} = \frac{1}{\sqrt{2}} E^{\mu \nu \beta} A_{\alpha \beta} \), see (11), can be transformed to a real matrix due to the SU(2) x U(1) degrees of freedom and also be imaginary.

From (59), (60) we have,

\[
\frac{d}{dx^\mu} (\frac{\partial}{\partial P_\mu} - \frac{d}{dx^\nu} \frac{\partial}{\partial P_{\mu \nu}})(U_k U^k \sqrt{-g}) = W^\mu \frac{\partial}{\partial P_\mu} - \frac{d}{dx^\nu} \frac{\partial}{\partial P_{\mu \nu}}(U_k U^k \sqrt{-g})
\]

We recall, \( W^\mu = \frac{\partial}{\partial P_\mu} - \frac{d}{dx^\nu} \frac{\partial}{\partial P_{\mu \nu}}(U_k U^k \sqrt{-g}) \)
\[ W^\mu = \]
\[ (-4 \frac{Z^\nu}{Z} P^\nu - 4 \frac{Z_m Z^n}{Z^3} P^\mu + 4 \frac{(Z^\nu P^\nu) P^\mu}{Z^3}) U^\nu P^\mu - 2 \frac{Z_m P^m Z^\mu}{Z^3} + 6 \frac{(Z^m P^m)^2}{Z^3} P^\mu = \]
\[ -4 \frac{Z^\nu}{Z^2} P^\nu - 4 \frac{Z_m Z^n}{Z^3} P^\mu + \]
\[ + 4 \frac{(Z^\nu P^\nu)}{Z^3} P^\mu + 4 \frac{(Z^m P^m)^2}{Z^3} P^\mu \]
\[ - 2 \frac{Z_m P^m}{Z^2} \left( \frac{Z^\mu}{Z} - \frac{Z_m P^m P^\mu}{Z^2} \right) = \]
\[ - 4 \frac{U^k}{Z} V_k + \frac{U^k U^j}{Z} P^\mu - 2 \frac{Z_m P^m}{Z^2} U^\mu = 0 \]
\[ W^\mu \cdot \mu = \left( - 4 U^\nu \cdot \nu \frac{P^\mu}{Z} - 2 \frac{(Z_m P^m)}{Z^2} U^\mu \right) \cdot \mu = 0 \] 

**Appendix B: Proof of conservation**

**Theorem:** Conservation law of the real Reeb vector.

From the vanishing of the divergence of Einstein tensor and (62) in the paper, we have to prove the following in 
\( (-1, +1, +1, +1) \) metric convention:

\[ \frac{1}{4} \left( U_\mu U_\nu - \frac{1}{2} U_k U^k g_{\mu \nu} - 2 U^k \cdot \nu \frac{P^\mu}{Z} \right) \cdot \mu = G_{\mu \nu} \cdot \mu = (R_{\mu \nu} - \frac{1}{2} R g_{\mu \nu}) \cdot \mu = 0 \] 

**Proof:**

From the zero variation by the scalar time field (63)

\[ W^\mu \cdot \mu = \left( - 4 U^\nu \cdot \nu \frac{P^\mu}{Z} - 2 \frac{(Z_m P^m)}{Z^2} U^\mu \right) \cdot \mu = 0 \] 

\[ - 4 \frac{U^\nu}{Z} \cdot \nu \frac{P^\mu}{Z} \cdot \mu = \frac{(Z_m P^m)}{Z^2} U^\mu \] 

\[ - 2 U^k \cdot \nu \frac{P^\mu}{Z} \cdot \mu = \frac{(Z_m P^m)}{Z^2} U^\mu \cdot \mu = P^\nu - \left( 2 U^k \cdot \nu \frac{P^\mu}{Z} \right) P^\nu \cdot \mu = \]

\[ \left( \frac{(Z_m P^m)}{Z^2} U^\mu \right) \cdot \mu P^\nu - U^k \cdot \nu \frac{Z^\nu}{Z} \]

Now let \( t = Z_m P^m \)

\[ \left( \frac{t}{Z^2} U^\mu \right) \cdot \mu P^\nu - U^k \cdot \nu \frac{Z^\nu}{Z} = \left( \frac{t}{Z^2} \right) \cdot \mu U^\mu P^\nu + \frac{t}{Z^2} U^\mu \cdot \mu U^\mu P^\nu - U^k \cdot \nu \frac{Z^\nu}{Z} \]

\[ - U^\mu \cdot \mu U^\nu + \left( \frac{t}{Z^2} \right) \cdot \mu U^\mu P^\nu \]

33
This is because \(-U^v = -\frac{Z^v}{Z} + \frac{t}{Z^2} P^v \Rightarrow -U^\mu;_\mu Z^v + \frac{t}{Z^2} U^\mu;_\mu P^v = -U^\mu;_\mu U^v\). Note that \(-U^v\) is minus twice the real numbered Reeb vector. So,

\[
(-2U^k;_k \frac{P^\mu P^v}{Z})_{;\mu} = -U^\mu;_\mu U^v + \left(\frac{t}{Z^2}\right)_{;\mu} U^\mu P^v
\]

(68)

Returning to the theorem we have to prove and using equation (68), we have to show,

\[
\left( U^\mu U^v - \frac{1}{2} U^k U^k g^{\mu\nu} - 2U^k;_k \frac{P^\mu P^v}{Z}\right)_{;\mu} = U^\mu;_\mu U^v + U^\mu U^v;_\mu - \frac{1}{2} (U^k;_\mu U^k + U^k U^k;_\mu) g^{\mu\nu} -
\]

\[
U^}\mu;_\mu U^v + (\frac{t}{Z^2})_{;\mu} U^\mu P^v =
\]

\[
U^\mu U^v;_\mu - \frac{1}{2} (U^k U^k);_\mu U^\mu P^v = 0
\]

(69)

Notice that

\[
U^\mu U^v;_\mu - \frac{1}{2} U^\nu U^v =
\]

\[
U^\mu \left(\frac{Z_\mu}{Z} ;\mu - \frac{t}{Z^2} \right)_{;\mu} P^i - \left(\frac{t}{Z^2}\right)_{;\mu} P^i =
\]

\[
U^\mu \left(\frac{Z_\mu}{Z} ;\mu - \frac{t}{Z^2} \right)_{;\mu} P^i =
\]

\[
- U^\mu \left(\frac{t}{Z^2}\right)_{;\mu} P^v
\]

(70)

Since \(-\left(\frac{t}{Z^2}\right)_{;k} P^k U^s = 0\) because the Reeb vector is perpendicular to the foliation kernel \(P^k,_{\frac{p^k}{\sqrt{2}},\frac{p^k}{\sqrt{2}}};_k = 0\).

Equation (70) is also a result of \(\ln(Z_{;k})_{;\mu} U^s_{;k} = \ln(Z)_{;k} U^s_{;k} = 0\) and of \(B^i;_\mu U^s_{;k} = P^i_{;k} U^s_{;k}\).

\[
U^\mu U^v;_\mu - \frac{1}{2} (U^k U^k);_\mu U^\mu P^v = -U^\mu \left(\frac{t}{Z^2}\right)_{;\mu} P^v + \left(\frac{t}{Z^2}\right)_{;\mu} U^\mu P^v = 0
\]

(71) and we are done.

**Appendix C: Invariance of the Reeb vector under different functions of P**

Here we wish to explore another degree of freedom in the action operator of the “acceleration field” which results from the Reeb vector, as shown by a representative vector field \(\frac{dP}{dx}\) which is tangent to a non-geodesic integral curve. We wish to show that \(P\) can be replaced with a smooth function \(f(P)\) and that \(U_m\) is invariant under such a transformation. We revisit our acceleration field and write \(U_m = \frac{N^2_{;m}}{N^2} - \frac{N^2_{;m} p^{m}}{N^4} p_m\) s.t. \(Z = N^2 = \frac{p^{k} p_k + p^k p^k}{2}\)

We can omit the comma for the sake of brevity the same way we write \(P_1\) instead of \(P_{;1}\) for \(\frac{dP}{dx}\) and write \(U_m = \frac{N^2_{;m}}{N^2} - \frac{N^2_{;m} p^{m}}{N^4} p_m\). We will prove the invariance of \(U_m\) where \(P\) is real, however, a similar proof is also valid where \(P\) is complex and where \(P\) is replaced with a smooth function of \(P\).
Suppose that we replace \( P \) by \( f(P) \) such that \( f \) is positive and increasing or decreasing, then

\[
f(P) = \frac{df(P)}{dx^i} = \frac{df(P)}{dP} \frac{dP}{dx^i} = f_\mu(P)P^\mu.
\]

Let \( N^2 \equiv P^4 \) then \( \dot{N}^2 \equiv f(P)f'(P)^2 = N^2 f_\mu(P)^2 \)

and \( \frac{\dot{N}^2_k}{N^2} = \frac{N^2_k}{N^2} \frac{2f_{\mu\nu}(P)}{f_\mu(P)} P_k \) but also

\[
\dot{U}_k = \frac{\dot{N}^2_k}{N^2} - \frac{\dot{N}^2_s f_\mu(P)P^\nu f_\nu(P)P_k}{N^2} = \frac{N^2_k}{N^2} \left( \frac{2f_{\mu\nu}(P)}{f_\mu(P)} - \left( \frac{N^2_s}{N^2} + \frac{2f_{\mu\nu}(P)}{f_\mu(P)} \right) \right) \frac{f_\nu(P)P^\nu f_\mu(P)P_k}{N^2 f_\mu(P)^2} = \frac{N^2_k}{N^2} - \frac{N^2_s P^\mu}{N^4} P_k = U_k
\]

(72)

**Appendix D: The curvature of the gradient of \( P \)**

The second power of the curvature of the integral curve by \( P_\mu = \frac{dP}{dx^\mu} \) where \( x^\mu \) denote the coordinates is expressible by

\[
Curv^2 = \frac{d}{dt} \frac{P_\lambda}{\sqrt{P_k P_k}} \frac{d}{dt} \frac{P_\mu}{\sqrt{P_k P_k}} g^{\lambda\mu}
\]

such that \( g^{\lambda\mu} \) is the metric tensor. (73) is an excellent candidate for an action operator. For convenience, we will write \( Norm = \sqrt{P_k P_k} \) and \( \dot{P}_\lambda = \frac{d}{dt} P_\lambda \). For the arc length parameter \( t \). Here is the main trick, as was mentioned about \( Z = Norm^2 \), \( Norm \) may not be constant because \( P_\lambda \) is not the 4-velocity of any particle, (to see an example of a variable \( Norm \), see Appendix E – The time field in the Schwarzschild solution). An arc length parameterization along these curves is equivalent to proper time measured by a particle that moves along the curves, and in the real numbers case, \( P \) can be indeed time. Unlike velocity’s squared norms, \( Z \) is not constant.

Let \( W_\lambda \) denote:

\[
W_\lambda = \frac{d}{dt} \left( \frac{P_\lambda}{\sqrt{P_k P_k}} \right) = \frac{\dot{P}_\lambda}{\text{Norm}} - \frac{P_\lambda}{\text{Norm}^3} P_k \dot{P}_v g^{kv}.
\]

Obviously

\[
W_\lambda P_k g^{jk} = \frac{\dot{P}_\lambda P_k g^{jk}}{\text{Norm}} - \frac{P_\lambda P_k g^{jk}}{\text{Norm}^3} P_k \dot{P}_v g^{kv} = \frac{\dot{P}_\lambda P_k g^{jk}}{\text{Norm}} - \frac{P_\lambda \dot{P}_v g^{kv}}{\text{Norm}} = 0
\]

Thus

\[
Curv^2 = W_\lambda W_\lambda = \frac{\dot{P}_\lambda \dot{P}_v g^{kv}}{\text{Norm}^2} - \frac{P_\lambda \dot{P}_v g^{kv}}{\text{Norm}^3} P_k \dot{P}_v = \frac{\dot{P}_\lambda \dot{P}_v}{\text{Norm}^2} - \left( \frac{P_\lambda \dot{P}_v}{\text{Norm}} \right)^2
\]

Following the curves formed by \( P_\lambda = P_{\lambda^2} = \frac{dP}{dx^\lambda} \), the term \( \frac{dx^\lambda}{dt} = \frac{P_\lambda}{\text{Norm}} \) is the derivative of the normalized curve or normalized “velocity”, using the upper Christoffel symbols, \( P_{\lambda^\mu} = \frac{d}{dx^\nu} P_\lambda - P_{\lambda^s} \Gamma^v_{\lambda^s} \).
\[ \frac{d}{dt} \mathbf{P}_\lambda = \left( \frac{d}{dx^i} \mathbf{P}_\lambda - \frac{\partial}{\partial x^i} \mathbf{P}_\lambda \right) \frac{dx^i}{dt} = \left( \mathbf{P}_{\lambda^i} \right) \frac{\partial}{\partial x^i} \mathbf{P}_\lambda \quad \text{such that} \quad x^r \text{ denotes the local coordinates.} \]

If \( \mathbf{P}_\lambda \) is a conserving field, then \( \mathbf{P}_{\lambda^i} = \mathbf{P}_{\lambda^i} \) and thus \( \mathbf{P}_{\lambda^i} \mathbf{P}^r = \frac{1}{2} \frac{\partial}{\partial x^i} \mathbf{P}_\lambda \), and

\[ \text{Curv}^2 = \frac{\mathbf{P} \mathbf{P}^i \mathbf{P}^i}{\text{Norm}} - \left( \mathbf{P} \mathbf{P}^i \mathbf{P}^i \right) \left( \mathbf{P} \mathbf{P}^j \mathbf{P}^j \mathbf{g}^{jk} \right) - \left( \mathbf{P} \mathbf{P}^i \mathbf{g}^{ij} \right)^2 \]

In the real case, we have achieved the Reeb vector,

\[ U_m = \frac{(\mathbf{P} \mathbf{P}_i)_{,m}}{\mathbf{P} \mathbf{P}_i} - \left( \frac{(\mathbf{P} \mathbf{P}_i)_{,\mu}}{(\mathbf{P} \mathbf{P}_i)^2} \mathbf{P} \mathbf{P}_m - \frac{Z_m}{Z} \right) = \frac{Z - Z \mathbf{P} \mathbf{P}_m}{Z^2} \]

and our candidate for a trajectory curvature action is

\[ \text{Action} = \frac{1}{4} U_m U^m \]

where in the complex case we saw

\[ \text{Action} = \frac{1}{8} (U^*_m U^m + U^*_m U^m) \]

Non-geodesic motion, as a result of interaction with a field, is not a geodesic motion in a gravitational field, i.e. it is not free fall. Moreover, material fields by this interpretation prohibit geodesic motion curves of particles moving at speeds less than the speed of light and by this, reduce the measurement of proper time.

**Appendix E – Time-like field from geodesic curves in the Schwarzschild solution**

**Motivation:** To show a non vanishing \( (\mathbf{P} \mathbf{P}^i)_m \), to make the reader familiar with the idea of maximal proper time from a sub-manifold and to calculate the background scalar time field of the Schwarzschild solution from that sub-manifold. We choose as a sub-manifold, a small 3 dimensional 3-sphere around the “Big Bang” singularity or a synchronized big sphere around the gravity source and far from the source and therefore this example is either limited to a “Big Bang” manifold or to a big sphere. So, we want to connect each event in a Schwarzschild solution to a primordial sub-manifold a fraction of second after the presumed “Big Bang” or to a synchronized big sphere around the gravity source, with the longest possible curve under the assumption that no closed time-like curves occur.

In this limited case, the scalar field is uninteresting as it does not represent interactions with any charged particle or with other force fields and therefore, the Reeb vector should be zero.

We would like to calculate

\[ \frac{U_m U^m}{4} = \frac{1}{4} \left( \left( \mathbf{P}_i \mathbf{P}_j \right)_{,m} \left( \mathbf{P}_i \mathbf{P}_j \right)_{,k} \mathbf{g}^{mk} - \left( \frac{\partial}{\partial x^i} \mathbf{P}_\lambda \right)_{,m} \mathbf{P}^m \right)^2 \]

in Schwarzschild coordinates for a freely falling particle. This theory predicts that where there is no matter, the result must be zero. The speed \( U \) of a falling particle from very far away, as measured by an observer in the gravitational field is

\[ V^2 = \frac{U^2}{c^2} = \frac{R}{r} = \frac{2GM}{r c^2} \]

Where \( R \) is the Schwarzschild radius. If speed \( V \) is normalized in relation to the speed of light then \( V = \frac{U}{c} \). For a far observer, the deltas are denoted by \( dt', dr' \) and,
\[ \dot{r}^2 = \left(\frac{dr}{dt}\right)^2 = \nu^2 \left(1 - \frac{R}{r}\right) \]  

(77)

because \( dr = dr' / \sqrt{1-R/r} \) and \( dt = dt' \sqrt{1-R/r} \).

\[
P = \int_0^t \left(1 - \frac{R}{r}\right) \frac{\dot{r}^2}{1-R} - \frac{1}{1-R} \right) \frac{1}{2} \right) dt = \int_0^t \left(1 - \frac{R}{r}\right) - \frac{\frac{R}{r}}{1-R} \right) \frac{1}{2} \right) dt =
\]

\[
\int_0^t \left(1 - \frac{R}{r}\right)^2 \frac{1}{2} dt = \int_0^t (1 - \frac{R}{r}) dt
\]

which results in,

\[ P_t = \frac{dP}{dt} = (1 - \frac{R}{r}) \]

(78)

Here \( t \) is not a tensor index and it denotes derivative by \( t \)!

On the other hand

\[
P = \int_0^r \left(1 - \frac{R}{r}\right) \frac{1}{2} - \frac{1}{1-R} \right) \frac{1}{2} \right) \]

\[
\int_0^r \sqrt{\frac{r}{R}} dr
\]

Which results in

\[ P_r = \frac{dP}{dr} = \sqrt{\frac{r}{R}} \]

(79)

Here, \( r \) is not a tensor index and it denotes derivative by \( r \)!

For the square norms of gradients, we use the inverse of the metric tensor,

So, we have \((1 - \frac{R}{r}) \rightarrow (1 - \frac{R}{r})^{-1} \) and \((1 - \frac{R}{r})^{-1} \rightarrow (1 - \frac{R}{r})\)

So, we can write

\[ N^2 = P_t P_t = (1 - \frac{R}{r}) P_t^2 - (1 - \frac{R}{r})^{-1} P_t^2 = (1 - \frac{R}{r})(\frac{r}{R} - 1) = \frac{r}{R} + \frac{R}{r} - 2 \]

\[ N^2 = \frac{r}{R} + \frac{R}{r} - 2 \]

(80)
\[ N^2_\lambda = \frac{dN^2}{dx^\lambda} \] And we can calculate
\[
\frac{N^2_\lambda N^2_\beta}{(N^2)^2} = \frac{(1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right)^2}{\left( \frac{r}{R} + \frac{R}{r} - 2 \right)^2}
\]

We continue to calculate
\[
N^2_\beta P_\beta = (1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right) \sqrt{\frac{r}{R}} \quad \text{and} \quad N^2_\beta P_\beta = (1 - \frac{R}{r}) \left( \frac{1}{R} - \frac{R}{r^2} \right) \sqrt{\frac{r}{R}}
\]

Note that here \( t \) is not a tensor index and it denotes derivative by \( t \) !
\[
(1 - \frac{R}{r}) N^2_r P_r = (1 - \frac{R}{r}) \left( \frac{1}{R} - \frac{R}{r^2} \right) \sqrt{\frac{r}{R}}
\]

Please note, here \( r \) is not a tensor index and it denotes derivative by \( r \) !
\[
N^2_\beta P_\beta = (1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right)^2 \left( \frac{r + R}{r} - 2 \right)
\]

So
\[
\frac{(N^2_\beta P_\beta)^2}{(N^2)^3} = \frac{(1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right)^2}{\left( \frac{r}{R} + \frac{R}{r} - 2 \right)^2}
\]

And finally, from (81) and (85) we have,
\[
\frac{(P_\beta P_\beta)_{\alpha m} (P^\alpha P^\beta)_{\beta k} g^{mk}}{(P_\beta P^\beta)^2} = \frac{(P^\alpha P_\alpha)_{\alpha m} P^{m}}{(P^\alpha P^\alpha)^3}
\]
\[
N^2_\beta N^2_\beta P_\beta \quad \text{and} \quad N^2_\beta N^2_\beta P_\beta = \frac{(1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right)^2}{\left( \frac{r}{R} + \frac{R}{r} - 2 \right)^2} - \frac{(1 - \frac{R}{r})^2 \left( \frac{1}{R} - \frac{R}{r^2} \right)^2}{\left( \frac{r}{R} + \frac{R}{r} - 2 \right)^2} = 0
\]

which shows that indeed the gradient of time measured, by a falling particle until it hits an event in the gravitational field, has zero curvature as expected.

**Appendix F: Conditions for SU(3) symmetry by three complex scalars in the kernel of \( P_\mu \)**

We may want to express the acceleration matrix \( A_{\alpha \beta} \) by three scalar fields that are defined in the foliation \( F \) that is perpendicular to \( \frac{P_i}{\sqrt{Z}} \). This is because \( P_i \) is a geometric object that defines foliations of space-time and can be
conversely defined by the foliations. Another motivation is to show that $SU(3)$ that is seen in Quantum Chromodynamics, may originate from geometry. By a theorem of Frobenius, necessary conditions for 3 vectors $h(j = 1, 2, 3)$ to span the foliation $F$ is that the vectors $h(\xi)$ are Holonomic if their Lie brackets depend on them.

$$[h(i), h(k)] = \sum_{j=1}^{3} c_j h(j)$$

for some coefficients $c_j$. The Lie brackets of each two vectors must depend on the vectors that span $T(F)$. We may write our 3 scalars $a, b, c$ (here $\xi$ is not the speed of light and not the previous coefficients but a scalar field) and their gradients that span the foliation’s tangent space $T(F)$ as follows,

$$h_k(1) = \frac{d}{dx^k} a, h_k(2) = \frac{d}{dx^k} b, h_k(3) = \frac{d}{dx^k} c.$$

We now express $A_{\mu \nu}$ by $h_k(1), h_k(2), h_k(3)$ in a covariant formalism but we need some constraint on $P_{\mu}$.

**Condition:** $P_{\mu}, \lambda P^{\mu \lambda} = P^{\mu \lambda} P_{\mu}, \lambda P^{\mu}$

This condition is not trivial and in general, $P_{\mu}, \lambda P^{\mu \mu} \neq P^{\mu \lambda} P_{\mu}, \lambda P^{\mu}$.

Consider the following matrix:

$$D_{ij} = g_{ij} - a^* i a_j - b^* i b_j - c^* i c_j$$

and $a_i = \frac{\alpha_i}{\sqrt{(\alpha^\lambda \alpha_{* \lambda} + \alpha^* \alpha_{\lambda})/2}}$ for some scalar function $\alpha$ whose gradient $a_i$ is in the foliation perpendicular to $P_\mu$ etc. and in the same manner replace $\beta_i$ by a normalized unit vector $b_i$ and $\gamma_i$ by a normalized vector. Also, we demand orthogonality, $a_i b^* i = a_i c^* i = b_i c^* i = 0$. Obviously

$$D_{ij} a^* i = D_{ij} b^* i = D_{ij} c^* i = 0$$

and $D_{ij} \frac{P^* i}{\sqrt{Z}} = \frac{P^* i}{\sqrt{Z}}$ which then shows that $D_{ij} = \frac{P^* i}{Z} P_j$,

$$Z = (P_\lambda P^{* \lambda} + P^{* \lambda} P_\lambda)/2.$$ That is very interesting, because as was already said, we can represent $P^* i P_j$ by orthogonal fields that span the tangent space of the perpendicular foliation to $P_j$, namely $T(F)$. Consider the following:

$$D_{ij; k} = D_{ik; j} = \left(\frac{P^* i P_j + P^* j P_i}{Z} - \frac{P^* i P_j Z_k}{Z^2} \right) - \left(\frac{P^* i P_j + P^* j P_i}{Z} - \frac{P^* i P_j Z_k}{Z^2} \right) = \frac{P^* i P_j Z_k}{Z^2} - \frac{P^* i P_j Z_k}{Z^2}$$

(88)

Now comes a little trick:
\[(D_{ij} - D_{ik}) \frac{P^i P^j}{Z} = (P^* P^* - P^* P^*) \frac{P^i P^j}{Z} + (P^* P^* - P^* P^*) \frac{P^i P^j}{Z} = \]
\[
\left( \frac{P^* P^j P^i_k}{Z^2} - \frac{P^* P^j P^i_k}{Z} \right) + \left( \frac{P^* P^j P^i_k}{Z} - \frac{P^* P^j P^i_k}{Z} \right) = \]
\[
\left( \frac{Z_j P^* P^j P^i_k}{Z^2} - \frac{Z_j P^* P^j P^i_k}{Z} \right) + \left( \frac{Z_j P^* P^j P^i_k}{Z} - \frac{Z_j P^* P^j P^i_k}{Z} \right) = \]

By (10), it is obvious that the first two terms constitute minus twice the Reeb vector,

\[
\frac{Z_j P^* P^j P^i_k}{Z^2} - \frac{Z_j}{Z} = -2 \frac{U_k}{2}. \]

For the last two terms, we need a special condition

\[
P_{\mu} P_{\mu} = P_{\mu} P_{\mu} \]

although usually \( P_{\mu} P_{\mu} \neq P_{\mu} P_{\mu} \). Then by this condition,

\[
\frac{P^* P^i_k}{Z} - \frac{P^i P^j P^i_k}{Z^2} = \frac{U_k}{2} \]

and therefore

\[
(D_{ij} - D_{ik}) D_{ij} = (D_{ij} - D_{ik}) \frac{P^i P^j}{Z} = - \frac{U_k}{2} \]

Consider our assumption,

\[
D_{ij} = g_{ij} - a^* \cdot a_j - b^* \cdot b_j - c^* \cdot c_j \]

and we have obtained an expression of the Reeb vector by the orthonormal vectors that represent the foliation. The additives of (87) are tensors. This leads us to an open question as follows: Is the condition \( P_{\mu} P_{\mu} = P_{\mu} P_{\mu} \), the minimal condition which is needed for a representation of the Reeb vector by \( a_j, b_j, c_j \) as the sum of tensor terms? In other words, is the condition \( P_{\mu} P_{\mu} = P_{\mu} P_{\mu} \) a necessity for the tensor representation of the acceleration matrix by the foliation scalars, \( a, b, c \)?

**The Electric, Electro-weak and Strong actions**

Our approach to foliations is very different than in any canonical Hamiltonian representation of space-time. The space-time manifold is an observer / reference object. Forces and probabilities are represented in at most 4 functions of space-time that describe forces through non-zero Reeb vectors. So indeterminism and matter should appear in a second layer of reality while space-time as an object is accessible to physical clocks only via these functions. According to this philosophy, there are other ways to achieve a Lagrangian with higher symmetries than U(1) that are shortly discussed. The following action can be extended to U(1) x SU(2) and to SU(3) symmetries by considering more than one Reeb vector.

\[
Z = N^2 = (P \mu \mu + P \mu \mu) / 2 \quad \text{and} \quad U \lambda = \frac{Z \lambda}{Z} - \frac{Z_k P \mu \mu}{Z^2} \quad \text{and} \quad \]
\[
L = \frac{1}{4} \left( U_k U \mu \mu \mu + U \mu \mu \mu U \mu \mu \right) / 2 \]

\[
(90) \]
Since the matrix of a Symplectic form can be described as two rotation and scaling hyper-planes, there is a possibility to locally add another scalar \( P(2) \) and the Reeb vector of its gradient \( P(2)_\mu \)

\[
Z(2) = N^2(2) = (P(2)_\mu ^* P(2) + P(2)^*_\mu P(2)^\mu )/2
\]

And

\[
Z(2)_\lambda = \frac{dZ(2)}{dx^\lambda}
\]

\[
\hat{U}(2)_\lambda = \frac{Z(2)_\lambda}{Z(2)} - \frac{Z(2)_k P(2)^*_k P_\lambda}{Z(2)^2} \quad \text{and} \quad U(2)_\lambda = \hat{U}(2)_\lambda - \hat{U}(2)_\nu \frac{P^*_\nu}{Z} P_\lambda
\]

And such that

\[
U(2)_\lambda U^*_\lambda = P(2)_\lambda P^*_\lambda = P(2)_\lambda U^*_\lambda = P_\lambda U(2)^*_\lambda = 0
\]

(91)

and obviously \( P_\lambda U^*_\lambda = 0 \) and \( P(2)_\lambda U(2)^*_\lambda = 0 \) from the definition of a Reeb vector.

The action is then dictated by the root of the Gram determinant and is added to the previous action,

\[
L = \frac{1}{4} \left( \frac{U^k U^*_k + U^*_k U_k}{2} \right) + \frac{1}{4} \left( \frac{U(2)^k \overline{U(2)}^*_k + \overline{U(2)}^*_k U(2)^k}{2} \right)
\]

(92)

The physical meaning of \( U(2)_\mu \) is another acceleration field in another plane. We will consider (91) as the “Electro-Weak Geometric Chronon Action”.

In the three dimensional space, Minkowsky perpendicular to \( P_\mu \) we can view three holonomic vector fields that span the foliation tangent space as required by the Frobenius theorem. These can be locally described by three gradients,

\( P(3)_\lambda, P(4)_\lambda, P(5)_\lambda \) and accordingly we can discuss their Reeb vectors, \( \hat{S}_\mu, \hat{W}_\mu, \hat{T}_\mu \) and their projection on the foliation perpendicular to \( \frac{P_\lambda}{\sqrt{Z}} \).

\[
S_\lambda = \hat{S}_\lambda - \hat{S}_\nu \frac{P^*_\nu}{Z} P_\lambda, \quad W_\lambda = \hat{W}_\lambda - \hat{W}_\nu \frac{P^*_\nu}{Z} P_\lambda, \quad T_\lambda = \hat{T}_\lambda - \hat{T}_\nu \frac{P^*_\nu}{Z} P_\lambda
\]

This time we can’t require the orthogonality condition which is described in (91) because there are no three Minkowsky – perpendicular hyper planes in space-time.

Now we need the third root of the determinant of the Gram matrix of these new three Reeb vectors and the action becomes a sum of “electric”, “eletro-weak” and “strong” actions in (+1,−1,−1,−1) convention, assuming that the fractions in (93) exist as a limit, e.g. where \( Z \to 0 \).

\[
L = \begin{pmatrix}
\frac{1}{8} U^k U^*_k + U^*_k U_k \\
0
\end{pmatrix} + \begin{pmatrix}
1 \\
0
\end{pmatrix} + \begin{pmatrix}
0 \\
0
\end{pmatrix} + \begin{pmatrix}
\frac{P_k U(2)^*_k + P^*_k U(2)^*_k}{2\sqrt{2Z}} \\
\frac{U(2)^k U^*_k + U(2)^*_k U_k}{8}
\end{pmatrix}
\]

(93)
An open question is if the metric tensor can be represented in one coordinate system as:

\[
\begin{align*}
  g^k_i &= \frac{P(i)_\mu P^*(k)_\mu + P(k)_\mu P^*(i)_\mu \mu}{(P(i)_\lambda P^*(i)_\lambda + P^*(i)_\lambda P(i)_\lambda)(P(k)_\nu P^*(k)_\nu + P^*(i)_\nu P(k)_\nu)^{1/2}}.
\end{align*}
\]

Appendix G: The Python code that was used for some of the mass ratio calculations out of area ratios

```python
import numpy as NP

class ELECTROGRAVITY_CLASS:
    def function_cubic_viete(self, a, b, c, d):
        # Viete's algorithm when all roots are real.

        b2 = NP.longdouble(b * b)
        b3 = NP.longdouble(b2 * b)
        a2 = NP.longdouble(a * a)
        a3 = a2 * a

        p = (3 * a * c - b2) / (3 * a2)
        q = (2 * b3 - 9 * a * b * c + 27 * a2 * d) / (27 * a3)

        offset = b / (3 * a)

        t1 = 2 * NP.sqrt(-p / 3) * NP.cos(NP.arccos(NP.sqrt(-3 / p)) * (3 * q) / (2 * p)) / 3
        t2 = 2 * NP.sqrt(-p / 3) * NP.cos(NP.arccos(NP.sqrt(-3 / p)) * (3 * q) / (2 * p)) / 3 - NP.pi / 3
        t3 = 2 * NP.sqrt(-p / 3) * NP.cos(NP.arccos(NP.sqrt(-3 / p)) * (3 * q) / (2 * p)) / 3 - 2 * NP.pi / 3

        x1 = t1 - offset
        x2 = t2 - offset
        x3 = t3 - offset

        return (x1, x2, x3)
```

MAIN_electrogravity_class = ELECTROGRAVITY_CLASS()

f = 1 - 1/96

x1, x2, x3 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, -f / 96, (f * f) / 192)

x4, x5, x6 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, -f / 96, (f * f) / 192)

f = 4 / NP.pi

x7, x8, x9 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, -f / 96, (f * f) / 192)

print("Anti-gravity: X1,X2,X3 = (%.14lf, %.14lf, %.14lf") % (x1, x2, x3))
print("Gravity: X4,X5,X6 = (%.14lf, %.14lf, %.14lf") % (x4, x5, x6))
print("Anti-gravity: X7,X8,X9 = (%.14lf, %.14lf, %.14lf") % (x7, x8, x9))

print("Muon mass in MeV/C^2 105.658374524")
print("Predicted electron mass in MeV/C^2 %.14lf" % ((105.658374524 * (x7-1)) / (1 + (x1-1)*(1-x4))))

x4, x5, x6 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, f / 96, f*f / 192)

x8 = (1 + NP.sqrt(1 - 1/(NP.pi * 6)))/2

print("Gravity: X4,X5,X6 = (%.14lf, %.14lf, %.14lf") % (x4, x5, x6))
print("Gravity: X8 = %.14lf" % x8)

print("Predicted Tau particle out of the W Boson 80385 MeV/C^2 = %.14lf " % (80385*(1-x4)/(1+0.5*(x7-x10))))

f = (1 - 1/96) * (4 / NP.pi)

x7, x8, x9 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, f / 96, (f * f) / 192)

x10, x11, x12 = \
MAIN_electrogravity_class.function_cubic_viete(1, -1, f / 96, (f * f) / 192)

print("Anti-gravity: X7,X8,X9 = (%.14lf, %.14lf, %.14lf") % (x7, x8, x9))
print("Gravity: X10,X11,X12 = (%.14lf, %.14lf, %.14lf") % (x10, x11, x12))
print("Average 1/(1 - (1/(X7-1) + 1/(1-x10))/2) = %.14lf" % (1/((x7 - x10)/2)))

print("Better prediction: Tau out of the W Boson 80385 MeV/C^2 = %.14lf " \ 
  % (80385*(1-x4)/(1+0.5*(x7-x10))))

x7 = (1 + NP.sqrt(1 + 1/(NP.pi * 6)))/2
x10 = (1 + NP.sqrt(1 - 1/(NP.pi * 6)))/2

print("Another prediction: Tau out of the W Boson 80385 MeV/C^2 = %.14lf " \ 
  % (80369 *(1-x4)/(1+0.5*(x7-x10))))

input("Press Enter to exit> ")
Appendix H: The C++ code that was used to approximate the Fine Structure Constant via geometry. There is a better calculation in Appendix I. There is an additional code that was used for the calculation of the exact Fine Structure Constant. The latter code was not included due to its length.

```cpp
#include "stdafx.h"
#include <math.h>
#include <conio.h>

// #define FINE_STRUCTURE_STEP 0.0001
#define FINE_STRUCTURE_LOOP 100000 // Must be 1 / step.
#define FINE_STRUCTURE_STEP ((long double)1.0/FINE_STRUCTURE_LOOP)
#define FINE_STRUCTURE_PI 3.1415926535897932384626433832795
#define FINE_STRUCTURE_ITERATIONS 32

// 0.5 of the distance.
long double FUNCTION_distance(long double d_angle,long double d_area_ratio)
{
    long double d_square_height,d_r,d_scaled,d_l,d_step,d_sum,
            d_l_square,d_delta_ratio;
    int i;
    d_step = sin(d_angle) / FINE_STRUCTURE_LOOP;
    d_square_height = cos(d_angle);
    d_square_height *= d_square_height;
    d_delta_ratio = d_area_ratio - 1;
    // d_scaled = d_square_height * d_area_ratio;
    d_l = 0;
    d_sum = 0;
    for (i = 0; i < FINE_STRUCTURE_LOOP; i++,d_l+=d_step)
    {
        d_l_square = d_l * d_l;
        d_r = sqrt(d_l_square + d_square_height);
        d_scaled = d_square_height * (1 + d_r * d_delta_ratio);
        d_sum += sqrt(d_l_square + d_scaled) / d_r;
    }
    return d_sum * d_step;
}```
// Calculate 0.5 of the distance.
long double FUNCTION_average_distance(long double ad_area_ratio)
{
    long double ad_x, ad_step, ad_sum, ad_distance;
    int i;
    ad_sum = 0;
    ad_step = (long double)FINE_STRUCTURE_PI * FINE_STRUCTURE_STEP;
    for (i = 0, ad_x = 0; i < FINE_STRUCTURE_LOOP; i++, ad_x += ad_step)
    {
        ad_distance = FUNCTION_distance(ad_x * 0.5, ad_area_ratio);
        // 2 * 2Pi / 4Pi = 1.
        // Also area that grows by b and is divided by an area that grows by b is 1.
        ad_sum += sin(ad_x) * ad_distance;
    }
    return ad_sum * ad_step;
}

void FUNCTION_roots(void)
{
    long double r_root1, r_root2, r_f1, r_f2;
    long double r_result;
    int i;
    // Start close to already known attractors & save time.
    r_root1 = 1.0048700774565755;
    r_root2 = 0.98905515302403790;
    r_f1 = 1.3354957147970252;
    r_f2 = 1.3284582439903136;
    for (i = 0; i < FINE_STRUCTURE_ITERATIONS; i++)
    {
        r_f1 = FUNCTION_average_distance(r_root1);
        // ...
    }
}
\[ r_{\text{root1}} = \frac{(192 \cdot r_{\text{root1}} \cdot r_{\text{root1}} + 2 \cdot r_{\text{root1}} \cdot r_{f1} - 1.0/(r_{f1} \cdot r_{f1}))}{192}; \]
\[ r_{\text{root1}} = \sqrt[3]{r_{\text{root1}}}; \]
\[ r_{f2} = \text{FUNCTION_average_distance}(r_{\text{root2}}); \]
\[ r_{\text{root2}} = \frac{(192 \cdot r_{\text{root2}} \cdot r_{\text{root2}} - 2 \cdot r_{\text{root2}} \cdot r_{f2} - 1.0/(r_{f2} \cdot r_{f2}))}{192}; \]
\[ r_{\text{root2}} = \sqrt[3]{r_{\text{root2}}}; \]
\[ r_{\text{result}} = \sqrt{\frac{1}{(r_{\text{root1}} - 1) \cdot (1 - r_{\text{root2}})}); \]
\[ \text{printf}("%.9lf, x1=\%.9lf, x2=\%.9lf, 1/f1=\%.9lf, 1/f2=\%.9lf\n", \]
\[ r_{\text{result}}, r_{\text{root1}}, r_{\text{root2}}, 1/r_{f1}, 1/r_{f2}); \]
\]
int main()
{
    while (_kbhit()) _getch(); // Clear keyboard input.
    \text{FUNCTION_roots();}
    while (_kbhit()) _getch(); // Clear keyboard input.
    \text{puts("Press Enter to exit the console.");}
    getchar();
    return 0;
}

\textbf{Appendix I: The C++ code for the exact inverse Fine Structure Constant}

// Physics_Sum.cpp : Defines the entry point for the console application.
//
#include "stdafx.h"
#include<math.h>
#include<conio.h>

long double PHYSICS_pi;

inline
void PHYSICS_cubic_viete(long double cv_a,
    long double cv_b,
    long double cv_c,
    long double cv_d,
long double cv_x[1])
{
  long double cv_b2, cv_b3, cv_a2, cv_a3, cv_p, cv_q, cv_offset, cv_2sqrt, cv_sqrt, cv_inner, cv_t;
  cv_b2 = cv_b * cv_b;
  cv_b3 = cv_b2 * cv_b;
  cv_a2 = cv_a * cv_a;
  cv_a3 = cv_a2 * cv_a;
  cv_p = ((long double)3.0 * cv_a * cv_c - cv_b2) / ((long double)3.0 * cv_a2);
  cv_q = ((long double)2.0 * cv_b3 - (long double)9.0 * cv_a * cv_b + (long double)27.0 * cv_a2 * cv_d) / ((long double)27.0 * cv_a3);
  cv_offset = cv_b / ((long double)3.0 * cv_a);
  cv_sqrt = sqrtl(-cv_p / (long double)3.0);
  cv_2sqrt = (long double)2.0 * cv_sqrt;
  cv_inner = acosl(((long double)3.0 * cv_q) / ((long double)2.0 * cv_p) / cv_sqrt) / (long double)3.0;
  cv_t = cv_2sqrt * cosl(cv_inner);
  cv_x[0] = cv_t - cv_offset;
}

inline void PHYSICS_full_cubic_viete(long double cv_a, long double cv_b, long double cv_c, long double cv_d, long double cv_x[3])
{
  long double cv_b2, cv_b3, cv_a2, cv_a3, cv_p, cv_q, cv_offset, cv_2sqrt, cv_sqrt, cv_inner, cv_t1, cv_t2, cv_t3;
  cv_b2 = cv_b * cv_b;
  cv_b3 = cv_b2 * cv_b;
  cv_a2 = cv_a * cv_a;
  cv_a3 = cv_a2 * cv_a;
  cv_p = (3 * cv_a * cv_c - cv_b2) / (3 * cv_a2);
  cv_q = (2 * cv_b3 - 9 * cv_a * cv_b * cv_c + 27 * cv_a2 * cv_d) / (27 * cv_a3);
  cv_offset = cv_b / (3 * cv_a);
  cv_sqrt = sqrtl(-cv_p / (long double)3.0);
  cv_2sqrt = (long double)2.0 * cv_sqrt;
  cv_inner =

47
acosl(((long double)3.0 * cv_q) /
   ((long double)2.0 * cv_p) / cv_sqrt) / (long double)3.0;

   cv_t1 = cv_2sqrt * cosl(cv_inner);
   cv_t2 = cv_2sqrt * cosl(cv_inner - PHYSICS_pi / 3);
   cv_t3 = cv_2sqrt * cosl(cv_inner - 2 * PHYSICS_pi / 3);

   cv_x[0] = cv_t1 - cv_offset;
   cv_x[1] = cv_t2 - cv_offset;
   cv_x[2] = cv_t3 - cv_offset;
}

inline
void PHYSICS_results(long double r_f, long double r_x[2])
{
   long double r_y[1];

   PHYSICS_cubic_viete(1, -1, -r_f / 96, (r_f * r_f) / 192, r_y);
   r_x[0] = r_y[0];

   PHYSICS_cubic_viete(1, -1, r_f / 96, (r_f * r_f) / 192, r_y);
   r_x[1] = r_y[0];
}

inline
void PHYSICS_full_results(long double r_f, long double r_x1[3],
   long double r_x2[3])
{
   PHYSICS_full_cubic_viete(1, -1, -r_f / 96, (r_f * r_f) / 192, r_x1);
   PHYSICS_full_cubic_viete(1, -1, r_f / 96, (r_f * r_f) / 192, r_x2);
}

long double PHYSICS_get_MUON_95_96_denominator(void)
{
   long double g_y[2];
   long double g_f;

   g_f = 95.0;
   g_f /= 96.0;

   PHYSICS_cubic_viete(1, -1, -g_f / 96, (g_f * g_f) / 192, g_y);
   PHYSICS_cubic_viete(1, -1, g_f / 96, (g_f * g_f) / 192, g_y + 1);

   return 1 + (g_y[0] - 1) * (1 - g_y[1]);
}

long double PHYSICS_get_MUON_gravity_antigravity_ratios(long double g_y[2])
{
   long double g_f;

   PHYSICS_pi = acosl(-1);

   g_f = 4;
   g_f /= PHYSICS_pi;

   PHYSICS_cubic_viete(1, -1, -g_f / 96, (g_f * g_f) / 192, g_y);
   PHYSICS_cubic_viete(1, -1, g_f / 96, (g_f * g_f) / 192, g_y + 1);

   g_y[0] -= 1;
long double PHYSICS_get_W_Z_ratio(long double g_W[2])
{
    long double g_f, g_r;
    long double g_W[2], g_Z[2];

    PHYSICS_pi = acosl(-1);
    g_f = 4.0;
    g_f /= PHYSICS_pi;
    g_W[0] = (1 + sqrt(1 + g_f / 24)) * 0.5 - 1;
    g_W[1] = 1 - (1 + sqrt(1 - g_f / 24)) * 0.5;
    g_f = 95.0;
    g_f /= 96.0;
    g_Z[0] = (1 + sqrt(1 + g_f / 24)) * 0.5 - 1;
    g_Z[1] = 1 - (1 + sqrt(1 - g_f / 24)) * 0.5;
    g_f = g_W[0] * g_W[1];
    g_W[0] = sqrtl(g_f);
    g_W[1] = sqrtl(g_Z[0] * g_Z[1]);
    g_r = sqrtl(g_W[0] / g_W[1]) * (1 + g_f);
    return g_r;
}

long double PHYSICS_get_derivative(long double g_xi, long double g_d, long double g_power)
{
    long double g_deriv1, g_deriv2, g_J, g_d_to_power, g_cos, g_sin,
    g_inverse_power;

    g_inverse_power = 1 / g_power;
    g_d_to_power = powl(g_d, g_power);
    g_J = g_xi * (1 + 1 / g_d_to_power);
    g_cos = cosl(g_J);
    g_sin = sinl(g_J);

    // g_deriv1 =
    // (sinl(g_J) * 0.5 * g_xi * g_power / (g_d_to_power * g_d) -
    // powl(g_d/ 95 * 95 * 96 * 96, g_inverse_power - 1) * g_inverse_power;
    g_deriv1 = -(2 / (g_cos * g_cos)) * g_sin * g_xi * g_power / (g_d_to_power * g_d);
    g_deriv2 =
        g_inverse_power * powl(95 * 95 * 96 * 96 / g_d, g_inverse_power - 1) *
        (95 * 95 * 96 * 96 / (g_d * g_d));
    return g_deriv1 - g_deriv2;
long double PHYSICS_get_derivative2(long double g_d, long double g_power) {
    long double g_deriv, g_inverse_power;
    g_inverse_power = 1 / g_power;
    g_deriv = g_inverse_power * powl(g_d / (95 * 95 * 96 * 96), g_inverse_power - 1) / 95 * 95 * 96 * 96;
    return g_deriv;
}

long double PHYSICS_get_alpha_iteration(double g_xi) {
    long double g_factor = 95 * 95 * 96 * 96, i;
    long double g_d = 600000, g_r, g_r1, g_r_mid, g_delta = 524288, g_d1, g_d2, g_ret;
    for (i = 0; i < 1000; i++) {
        g_ret = g_factor / g_d;
        g_r = g_ret - 2/cosl(g_xi*(1 + 1 / g_d));
        g_d1 = g_d + g_delta;
        g_r1 = g_factor / g_d1 - 2 / cosl(g_xi*(1 + 1 / g_d1));
        g_delta *= 0.5;
        g_d2 = g_d + g_delta;
        g_r_mid = g_factor / g_d2 - 2 / cosl(g_xi*(1 + 1 / g_d2));
        if (g_r_mid>0) g_d = g_d2;
        printf("1/Alpha: %.18lf, error %.18lf\n", g_ret, g_r);
    }
    printf("1/Alpha: %.18lf, error %.18lf\n", g_ret, g_r);
    return g_r;
}

long double PHYSICS_get_alpha_iteration2(double g_xi) {
    long double g_factor = 96 * 96 * 96 * 96, g_nom;
    int i;
    long double g_d = 600000, g_r, g_r1, g_r_mid, g_delta = 524288, g_d1, g_d2, g_ret;
    g_nom = 1;
    g_nom *= (long double)96 / 95;
    g_nom *= (long double)96 / 95;
    for (i = 0; i < 1000; i++) {
        g_ret = g_factor / g_d;
        g_r = g_ret - 2 / cosl(g_xi*(1 + g_nom / g_d));
        g_d1 = g_d + g_delta;
        g_r1 = g_factor / g_d1 - 2 / cosl(g_xi*(1 + g_nom / g_d1));
    }
g_delta *= 0.5;
g_d2 = g_d + g_delta;
g_r_mid = g_factor / g_d2 - 2 / cosl(g_xi*(1 + g_nom / g_d2));

if (g_r_mid>0) g_d = g_d2;

printf("1/Alpha: %.18lf, error %.18lf\n", g_ret, g_r);
}

printf("1/Alpha: %.18lf, error %.18lf\n", g_ret, g_r);
return g_r;
}

long double PHYSICS_get_alpha_iteration3(long double g_xi,
                                        long double g_45_207)
{
    long double g_factor = 95 * 95 * 96 * 96, i;
    long double g_d = 600000, g_r, g_r1, g_r_mid,
        g_delta = 524288, g_d1, g_d2, g_ret;

    for (i = 0; i < 1000; i++)
    {
        g_ret = powl(g_factor / g_d, 1 + g_45_207);
        g_r =
            g_ret
            - 2 / cosl(g_xi*(1 + 1 / powl(g_d,1.0/(1+g_45_207))));
        g_d1 = g_d + g_delta;
        g_r1 =
            powl(g_factor / g_d1, 1 + g_45_207)
            - 2 / cosl(g_xi*(1 + 1 / powl(g_d1, 1.0 / (1 + g_45_207))));
        g_delta *= 0.5;
        g_d2 = g_d + g_delta;
        g_r_mid =
            powl(g_factor / g_d2, 1 + g_45_207)
            - 2 / cosl(g_xi*(1 + 1 / powl(g_d2, 1.0 / (1 + g_45_207))));
        if (g_r_mid>0) g_d = g_d2;
        // printf("1/Alpha: %.18lf, error %.18lf\n", g_ret, g_r);
    }

    printf("1/Alpha, s iterations: %.18lf, d %.18lf, error %.18lf\n",
            g_ret, powl(g_d, 1.0 / (1 + g_45_207)), g_r);
    return g_r;
}

long double PHYSICS_get_alpha_iteration4(long double g_xi,
                                        long double g_d)
{
    long double g_factor = 95 * 95 * 96 * 96, i;
    long double g_p = 95.9, g_r, g_r1, g_r_mid,
        g_delta = 0.2, g_p1, g_p2, g_ret;

    for (i = 0; i < 1000; i++)
    {
\[ g_{\text{ret}} = \text{powl}(g_{\text{factor}} / g_d, 1 + 1/(g_p * g_p)); \]
\[ g_r = \]
\[ g_{\text{ret}} - 2 / \text{cosl}(g_{\xi}*(1 + 1 / \text{powl}(g_d, 1.0 / (1 + 1 / (g_p * g_p)))); \]
\[ g_{p1} = g_p + g_{\delta}; \]
\[ g_{r1} = \text{powl}(g_{\text{factor}} / g_d, 1 + 1 / (g_{p1} * g_{p1})) - 2 / \text{cosl}(g_{\xi}*(1 + 1 / \text{powl}(g_d, 1.0 / (1 + 1 / (g_{p1} * g_{p1}))))); \]
\[ g_{\delta} *= 0.5; \]
\[ g_{p2} = g_p + g_{\delta}; \]
\[ g_{r_mid} = \text{powl}(g_{\text{factor}} / g_d, 1 + 1 / (g_{p2} * g_{p2})) - 2 / \text{cosl}(g_{\xi}*(1 + 1 / \text{powl}(g_d, 1.0 / (1 + 1 / (g_{p2} * g_{p2}))))); \]
\[ \text{if} (g_{r_mid}>0) g_p = g_{p2}; \]
\[ // \text{printf}("1/\text{Alpha}: \%.18lf, \text{error} \%.18lf\n", g_{\text{ret}}, g_r); \]
\[ \text{printf}("1/\text{Alpha}, p \text{ iterations}: \%.18lf, \text{error} \%.18lf, d \%.18lf, p \%.18lf\n", \]
\[ g_{\text{ret}}, g_r, \text{powl}(g_d, 1.0 / (1 + 1 / (g_{p2} * g_{p2}))), g_{p}); \]
\[ \text{return} g_r; \]

long double PHYSICS\_tau\_by\_iteration(void)\{
\long double tbi\_f, tbi\_f1, tbi\_f\_mid, tbi\_delta, tbi\_r, tbi\_r1, tbi\_r\_mid,
\quad tbi\_x[2], tbi\_95\_96, tbi\_w, tbi\_45, tbi\_207, tbi\_weinberg, tbi\_tau,
\quad tbi\_\tau\_e1, tbi\_\tau\_e2, tbi\_w\_e, tbi\_z\_e, tbi\_geome\_avg;
\long double tbi\_speculation\_attacks[2], tbi\_d, tbi\_d1, tbi\_d\_2,
\quad tbi\_45\_207\_denom, tbi\_err\_r, tbi\_inverse\_fsc;
\int i;
\long double tbi\_f = 1.5;
\long double tbi\_delta = 0.1;
PHYSICS\_\pi = \text{acosl}(-1);
\for (i = 0; i < 1000; i++) \{
\quad PHYSICS\_\text{results}(tbi\_f, tbi\_x);
\quad tbi\_r = \text{sqrt}(1/(tbi\_x[0]-1)) - 0.5 / (1 - tbi\_x[1]);
\quad tbi\_\tau = tbi\_x[1];
\quad tbi\_f1 = tbi\_f + tbi\_delta;
\quad PHYSICS\_\text{results}(tbi\_f1, tbi\_x);
\quad tbi\_r1 = \text{sqrt}(1 / (tbi\_x[0] - 1)) - 0.5 / (1 - tbi\_x[1]);
\quad tbi\_\tau = tbi\_x[1];
\quad tbi\_\delta *= 0.5;
\quad tbi\_\tau\_e1 = tbi\_\tau + tbi\_\delta;
\quad PHYSICS\_\text{results}(tbi\_\tau\_e1, tbi\_x);
\quad tbi\_\tau\_e2 = tbi\_\tau\_e1;
\quad tbi\_\tau\_\text{err} = tbi\_\tau\_e2;
\quad tbi\_\tau\_\text{e} = tbi\_\tau\_\text{err};
\quad tbi\_\text{geome\_avg} = tbi\_\tau\_\text{e};
\quad for (i = 0; i < 2; i++) \{
\quad tbi\_\text{speculation\_attacks}[i] = tbi\_\text{geome\_avg};
\quad tbi\_\text{d} = tbi\_\text{speculation\_attacks}[i];
\quad tbi\_\text{d1} = tbi\_\text{d};
\quad tbi\_\text{d2} = tbi\_\text{d};
\quad tbi\_\text{45\_207\_denom} = tbi\_\text{d};
\quad tbi\_\text{err\_r} = tbi\_\text{d};
\quad tbi\_\text{inverse\_fsc} = tbi\_\text{d};
\quad \}
\quad \}
\quad \}
\quad \}
\quad \}
\quad // PHYSICS\_\text{get\_alpha\_iteration2}(tbi\_f);\}
tbi_tau = 0.5/(1-tbi_tau);

printf("Xi: %.16lf\n", tbi_f);

PHYSICS_results(tbi_f - 0.0001, tbi_x);
tbi_speculation_attacks[0] = sqrt(1 / (tbi_x[0]-1));
PHYSICS_results(tbi_f + 0.0001, tbi_x);
tbi_speculation_attacks[1] = sqrt(1 / (tbi_x[0]-1));
tbi_weinberg = PHYSICS_get_W_Z_ratio(tbi_x);

printf("Z/W, theory: %.16lf\n", tbi_weinberg);

PHYSICS_results(tbi_f - 0.0001, tbi_x);
tbi_207 = tbi_x[0];
tbi_45 = tbi_x[1];

PHYSICS_get_alpha_iteration3(tbi_f, tbi_207 * tbi_45);

// Long double tbi_deriv;

printf("r45, r207 = %.16lf, %.16lf\n", 1/tbi_45, 1/tbi_207);

PHYSICS_get_alpha_iteration4(tbi_f, tbi_r);

// Low absolute derivative means numerical difficulties.
//tbi_deriv =
//PHYSICS_get_derivative2(tbi_r, tbi_45_207_denom);
//tbi_deriv =
//PHYSICS_get_derivative(tbi_f, tbi_r, tbi_45_207_denom);

PHYSICS_get_MUON_95_96_denominator();

printf("d before raised to power -1/(1 + (a-1)(1-b)) : %.20lf\n", tbi_r);

// tbi_d = powl(tbi_r, 1.0/(1.0+1.0/((96+1.0/18)*(96 + 1.0 / 18))));

PHYSICS_get_MUON_95_96_denominator();

printf("d: %.16lf, 1/Alpha of 96^2: %.16lf\n", tbi_d, 2 / cosl(tbi_f * (1 + (long double)1 / tbi_d)));

print("1) >>>>>>>>>>>>>>>>>>> d: %.16lf, 1/Alpha theory: %.16lf\n", tbi_d, tbi_inverse_fsc);

// Low absolute derivative means numerical difficulties.
//tbi_deriv =
//PHYSICS_get_derivative2(tbi_r, tbi_45_207_denom);
//tbi_deriv =
//PHYSICS_get_derivative(tbi_f, tbi_r, tbi_45_207_denom);

PHYSICS_get_MUON_95_96_denominator();

printf("d before raised to power -1/(1 + (a-1)(1-b)) : %.20lf\n", tbi_r);

// tbi_d = powl(tbi_r, 1.0/(1.0+1.0/((96+1.0/18)*(96 + 1.0 / 18))));

PHYSICS_get_MUON_95_96_denominator();

printf("d: %.16lf, 1/Alpha of 96^2: %.16lf\n", tbi_d, 2 / cosl(tbi_f * (1 + (long double)1 / tbi_d)));

print("2) d: %.16lf, 1/Alpha of 96^2: %.16lf\n", tbi_d, 2 / cosl(tbi_f * (1 + (long double)1 / tbi_d)));
tbi_d1 = powl((95 * 95 * 96 * 96) / tbi_r, 1 + tbi_45*tbi_207 / (1+1.0/32768));

tbi_err_r = tbi_inverse_fsc / (tbi_d1 - tbi_inverse_fsc);

printf("3) Err^\-1:%.16lf , r: %.16lf, 1/Alpha approximation: %.16lf\n", tbi_err_r, tbi_r, tbi_d1);

// tbi_d2 = tbi_r * powl(96 * 96 * tbi_x[0] * tbi_x[1], 1 / (96 * 96 * tbi_x[0] * tbi_x[1]));
tbi_d2 = tbi_r * 96 * 96 * tbi_x[0] * tbi_x[1];

tbi_inverse_fsc = 2 / cosl(tbi_f * (1 + (long double)1 / tbi_d2));

printf("4) d2: %.16lf, 1/Alpha approximation: %.16lf\n", tbi_d2, tbi_inverse_fsc);

tbi_d = 1/(1-powl(tbi_r, tbi_45_207_denom) / (96*96*tbi_x[0] * tbi_x[1]*tbi_r));

// tbi_d2 = tbi_r * powl(96 * 96 * tbi_x[0] * tbi_x[1], 1 / (96 * 96 * tbi_x[0] * tbi_x[1]));
tbi_d2 = tbi_r * 96 * 96 * tbi_x[0] * tbi_x[1];

tbi_inverse_fsc = 2 / cosl(tbi_f * (1+(long double)1/606400.8));

printf("d = 606400.8, 1/Alpha %.16lf\n", 2/cosl(tbi_f * (1+(long double)1/606400.8)));

return tbi_f;
References


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