

# Spin<sup>1/2</sup> 'Plane' & Simple

**Abstract:** To fully characterize any spin requires identification of its primary spin axis *and* its plane of rotation. Classical presumptions obscure both for “intrinsic” spin. Here, Euclidean interval-time coordinates literally *lift the veil* of space to reveal it. Probability amplitude is also physically realized.

## Retreat by Conceit?

Several generations of unsurpassed success in Quantum Mechanics has transformed “quantum weirdness” from a set of *mysteries* to be solved, to hollow *cornerstones*, precariously built upon. Foremost among these is quantum spin. Cautions abound:

“...spin is an intrinsic property of a particle, unrelated to any sort of motion in space.”<sup>1</sup>

“Physically, this means it is ill-defined what axis a particle is spinning about”<sup>2</sup>

“...any attempt to visualize it [spin<sup>1/2</sup>] classically will badly miss the point.”<sup>3</sup>

“[Quantum spin] has nothing to do with motion in space...but is somewhat analogous to classical spin”<sup>4</sup>

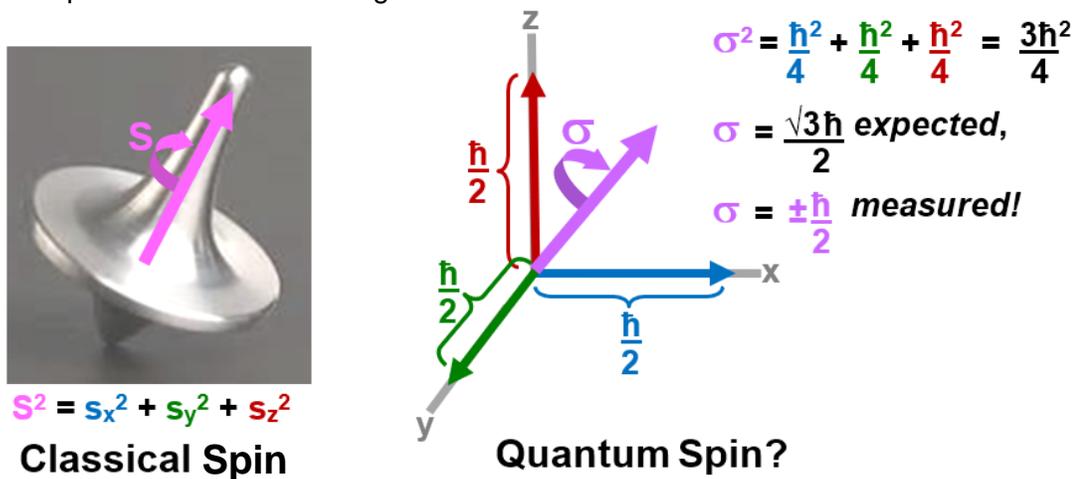
“the spin...of a fundamental fermion...with no classical analog, is...abstract...with no possibility of intuitive visualization.”<sup>5</sup>

Is this resignation to a fundamental reality or complacency, accompanying conceit? To avoid the latter, all feasible models must be exhausted. To be sure, fermion spin is *abstractly* modeled with spinors<sup>6</sup>, but physics isn't physics unless it's about the physical! A real, physical model is presented here.

## Top Performer

Consider some facts about “spin<sup>1/2</sup>” which, for massive fermions, may be conveniently considered in a rest frame. Having *real* angular momentum, with measurable spin components, a *primary spin vector* ( $\sigma$ ) is presumed to exist. The goal is to identify its axis.

1. The primary spin axis is *not* a classical axis. Such a spin vector would decompose to ordinary spatial spin components and this is not the case (Fig. 1).
2. The primary spin axis is *not* orthogonal to flat space. That would yield zero-magnitude spatial spin components, contrary to those observed.
3. The primary spin axis makes *equal* angles with *every* spatial direction, as revealed by equal size spin components ( $\hbar/2$ ), measured in any direction.<sup>7</sup>
4. As a form of angular momentum, intrinsic spin is a conserved quantity.<sup>8</sup>
5. Spin<sup>1/2</sup> has a 720° horizon. It takes *two* classical rotations to return a fermion to its *exact* original state<sup>7</sup> (accounting for probability amplitude).
6. The signs ( $\pm$ ) of intrinsic spin components correlate probabilistically. Having measured one, the sign of any other component relates to the angle between them.<sup>9</sup>



**Fig. 1.** Left: Classical spin vector ( $S$ ) decomposes to coordinate projections. Right: If spin components  $\hbar/2$  are found on three spatial coordinates, a resultant spin vector ( $\sigma$ ) of  $\sqrt{3}\hbar/2$  is *expected*. Instead,  $\hbar/2$  (or  $-\hbar/2$ ) is *measured* there as well! ( $\hbar$  is the reduced Planck constant  $h/2\pi$ .)

## Putting a Noether Spin on It

Conservation laws are of two types.<sup>10</sup> “Exact” laws are never expected to fail while “approximate” laws hold true within restrictions. For example, *conservation of mass-energy* is exact but divides into two approximate laws: *conservation of mass* and *conservation of energy*, which each hold, barring interchanges (exhibiting  $E = mc^2$ ).

By Noether’s theorem, each conservation law is now associated with a symmetry (i.e. *transformational invariance*).<sup>11</sup> For example, *conservation of mass-energy* is associated with temporal invariance. The mass-energy of an isolated system is unaltered by temporal displacement. It is “time invariant” (Fig. 2).

Two *omissions* in the exact laws should seem conspicuous:

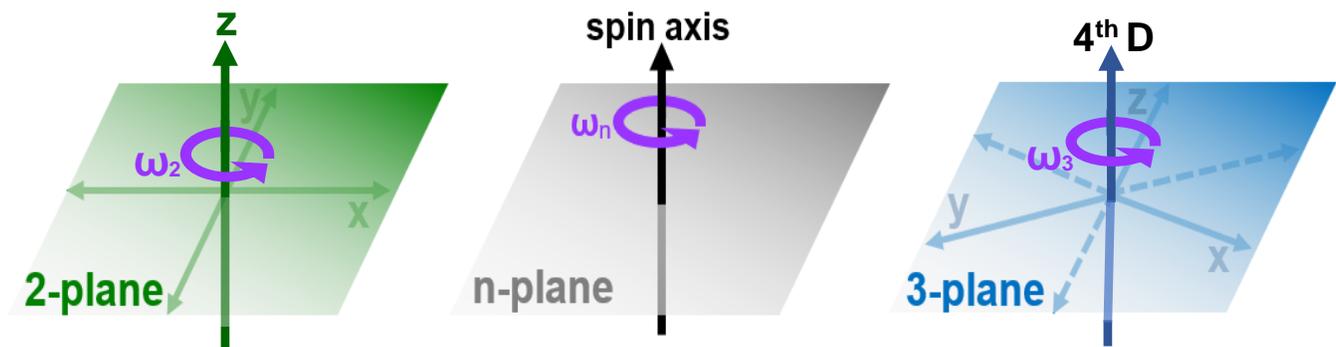
1. Laws exist for both temporal *and* spatial **translation** but only about spatial axes for **rotation**.
2. A *non-spatial*, primary “intrinsic” spin axis ( $\sigma$ ) *should* list with the spatial axes for angular momentum.

Exact Law	Noether Symmetry Invariance	Number of Dimensions	
Conservation of mass-energy	Time invariance	1	translation along time
Conservation of linear momentum	Translation symmetry	3	translation along x, y, z
Conservation of <b>angular momentum</b>	Rotation invariance	3	rotation about x, y, z, $\sigma$
CPT symmetry	Lorentz invariance	1+1+1	charge ( $q \rightarrow -q$ ) + position ( $r \rightarrow -r$ ) + time ( $t \rightarrow -t$ ) inversions
Conservation of electric charge	Gauge invariance	1 $\otimes$ 4	scalar field (1D) in 4D
Conservation of color charge	SU(3) Gauge invariance	3	r, g, b
Conservation of weak isospin	SU(2) <sub>L</sub> Gauge invariance	1	weak charge
Conservation of probability	Probability invariance	1 $\otimes$ 4	$\Sigma$ probabilities = 1 in x,y,z

**Fig. 2** It is inconsistent to purport 4D of translation (rows 1+ 2) and yet pretend only 3D are available for rotation (row 3). Though widely acknowledged to be non-spatial,<sup>1-5</sup> an additional spin dimension should be recognized by adding the primary fermion spin axis,  $\sigma$ .

## The ‘Plane’ Truth

To understand intrinsic (or any) spin, one must identify not only the *primary spin axis* but the *plane of rotation*. Classically, knowing one entails the other, but this is not a given in 4D. A spin vector is defined perpendicular to its plane of rotation in a Euclidean space but that is *not* what spacetime provides.



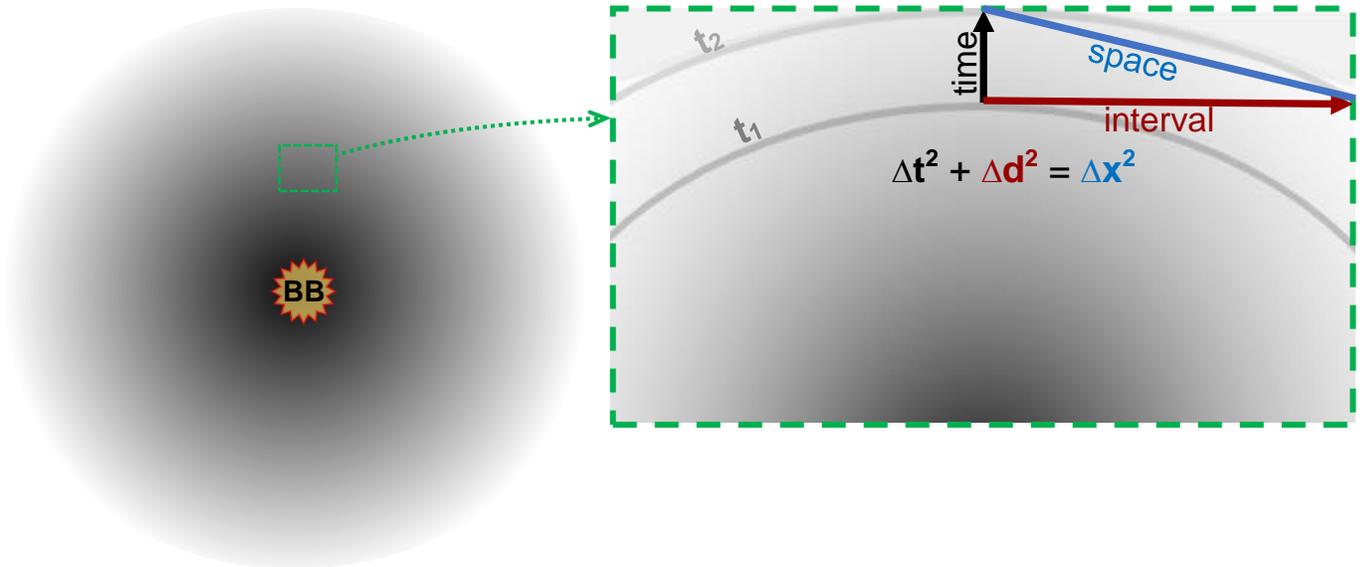
**Fig. 3.** An orthogonal spin axis makes no projections on a flat plane of any dimensionality.

## Put Space in Its Place

To accommodate survival (e.g. bow hunting), our brains are hardwired for space and time, which have become default coordinates. But Minkowski spacetime has a hyperbolic geometry,<sup>12</sup> which yields a distorted view in many respects.

*“Such maps necessarily distort metric relations and one has to compensate for this distortion.”<sup>13</sup>*

Adapting the balloon analogy of cosmic expansion<sup>14</sup> to contain a central Big Bang event, gives rise to a *curved-space, radial-time* model. Time emanates from the center as a 4D temporal field, enclosed by spatial 3-spheres, representing simultaneities in the rest frame of the cosmos. A radius corresponds to the age of the universe (Fig. 4). All locations on a 3-sphere, find Euclidean coordinates, with time normal to space and intervals tangent to it. A more detailed explanation<sup>14</sup> and illustrations<sup>15-18</sup> are provided.



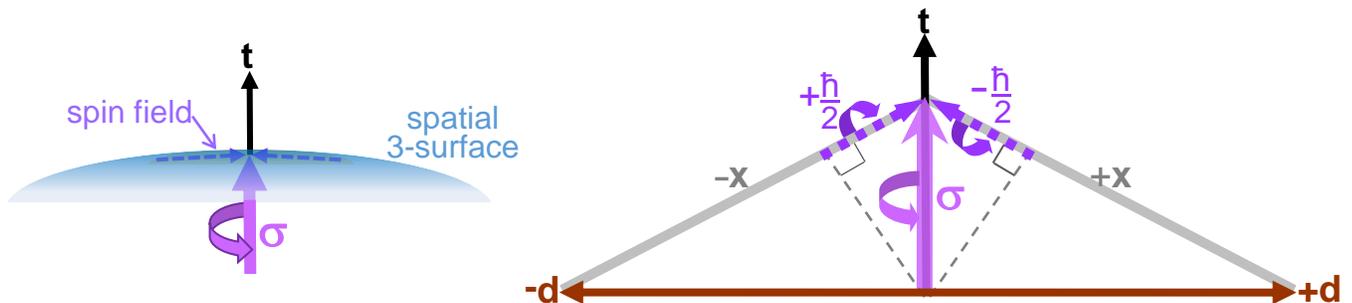
**Fig. 4** Left: A temporal 4-field, centered on the Big Bang (BB) yields a curved-space, radial-time model. Right: The indicated region, between earlier ( $t_1$ ) and later ( $t_2$ ) simultaneities, illustrates Euclidean, interval-time coordinates, allowing for spatial flatness.

Alternatively, consider that the Pythagorean theorem applies uniquely to Euclidean geometry. Adopting a spacelike convention, the interval formula<sup>19</sup>:  $\Delta d^2 = \Delta x^2 - \Delta t^2$  rearranges simply as  $\Delta x^2 = \Delta d^2 + \Delta t^2$ , which implies interval-time coordinates corresponding to the legs of a right triangle.

### Time to Turn Things Around!

With that Euclidean lens,  $\sigma$  can be modeled as *chronaxial spin*, in an *interval 3-plane*. Intrinsic spin is *classical* spin about a *non-classical* axis, time. More generally, it is spin about a particle's worldline which, in its rest frame, is its *timeline*. No longer a coordinate, space instead arcs past  $\sigma$  like an umbrella over its handle (Fig. 5). A field of spin components thus projects equally in *every* spatial direction, consistent with an underlying *curved-space, radial-time* structure.

Relativity makes a 4<sup>th</sup> dimension of spin axes unsurprising. All fermions age (e.g. muons decay) so, time undeniably supports *translation*. There is thus, no basis to deny that time also supports *rotation*. An objection might be that chronaxial spin is effectively instantaneous, easily developing circumferential speeds exceeding universal limit  $c$ . However, a fermion "point particle" of zero radius invokes no such restriction. In fact, instantaneous chronaxial spin provides a perfect source for quantum indeterminism.

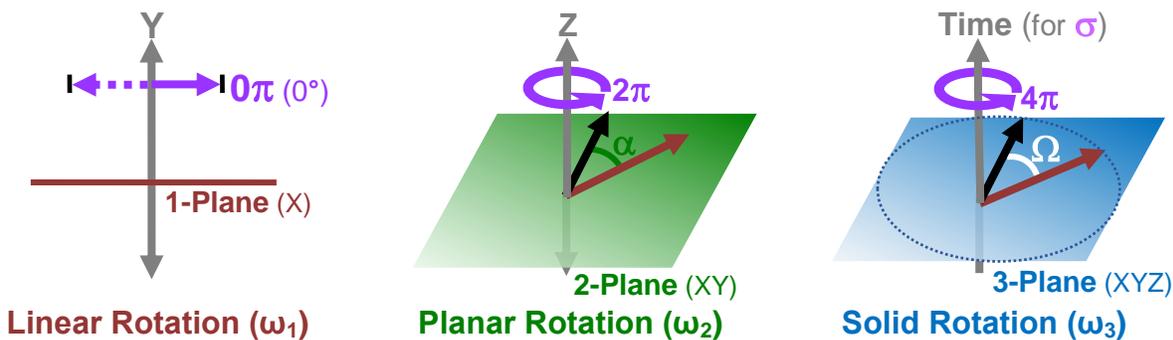


**Fig. 5.** Left: Arching past time, space exhibits a symmetric field of spin projections in all directions from the primary *chronaxial* spin vector ( $\sigma$ ). Right: For clarity, a 2D slice in Euclidean interval-time coordinates shows spatial arc ( $\pm x$ ) locally flat and highly inclined.  $\sigma$  projects symmetric  $\pm \hbar/2$  components on space.

## Solid Reasoning

Just as *conservation of mass-energy* divides to approximate laws for mass and energy, *conservation of angular momentum* divides into three *approximate* laws (barring interchanges), distinguished by the dimensionality of their angular velocities ( $\omega_1, \omega_2, \omega_3$ ). Time is always excluded from an *n-plane* of rotation since its fundamental unidirectionality denies the needed oscillatory freedom.

1. Conservation of **Linear**-Angular Momentum - vibration, simple harmonic motion.  $\omega_1 = 0\pi f$
2. Conservation of **Planar**-Angular Momentum - classical spin & orbits.  $\omega_2 = 2\pi f$
3. Conservation of **Solid**-Angular Momentum - quantum spin & *orbitals* (both chronaxial).  $\omega_3 = 4\pi f$



**Fig. 6.** Dimensional Spin Progression: Each rotation occurs in a flat *n-plane* about an orthogonal axis. Angular velocity ( $\omega_n$ ) relates to the approximately-conserved, angular momentum of each. Solid angle  $\Omega$  appears circular in a 3-plane (above right) and within a sphere (below).

Planar rotation entails  $2\pi$  radians. Going up a dimension, chronaxial spin may be depicted in a 3-plane about a timeline, where a sphere's volume is flatly exposed (Fig. 6). This entails a *solid angle* of  $4\pi$  steradians (sr).<sup>20</sup>

## Easy as Pi

"A half quantum" is an oxymoron because a "quantum" is "*the minimum amount of any physical entity involved in an interaction.*"<sup>21</sup> Yet "spin $\frac{1}{2}$ " implies such a halving, arising from inadequate *reduction* of the Planck constant ( $h$ ).

*"In applications where it is natural to use the angular frequency (i.e. ...in terms of radians per second...) it is often useful to absorb a factor of  $2\pi$  into the Planck constant...called the reduced Planck constant ...equal to the Planck constant divided by  $2\pi$ , and is denoted ' $\hbar$ ' (pronounced 'h-bar')"*<sup>22</sup>

Division by  $2\pi$  is fine for *classical* rotation, but there is *no basis* to apply this to quantum spin. Solid-angular, *chronaxial* spin must instead be reduced by  $4\pi$ . Applied to fermions, spin is not " $\frac{1}{2}$ " but quite *whole* at  $h/4\pi$  (i.e.  $\hbar/2$ ), exactly as measured (Fig. 7). Further, QED rightly boasts 12 digits of precision for the electron magnetic moment, but mysteriously remains off by a factor of two!

*"... one famous triumph of the Quantum Electrodynamics theory is the accurate prediction of the electron g-factor. The magnetic moment of an electron is approximately twice what it should be in classical mechanics. The factor of two [ $g_s$ ] implies that the electron appears to be twice as effective in producing a magnetic moment as the corresponding classical charged body. ...a correction term [ $a_e$ ]... takes account of ...interaction...with the magnetic field"*<sup>23</sup>

$$\mu_s = -\frac{g_s \mu_B S}{\hbar} = -\frac{g_s \mu_B S}{\frac{h}{2\pi}} \quad \leftarrow 4\pi$$

**Fig. 7** Denominator  $\hbar$  is only *half* reduced, as  $h/2\pi$ . Correcting with *solid-angular* range  $4\pi$  is equivalent to having a factor of 2 in the numerator. Thus,  $g_s$  does not mysteriously need to be "*twice*" the classical  $g$ -factor  $g_L$ .<sup>24</sup> The anomalous magnetic moment ( $a_e$ )<sup>25</sup> is then accommodated at *half* the conventional value in:  $g_s = 1 + a_e = 1.001159652181643$ .

Both  $\mathbf{S}$  (electron spin angular momentum) and  $\mu_B$  (Bohr magneton)<sup>23</sup> incorporate  $\hbar/2$  which, in that form, is fully reduced (i.e.  $h/4\pi$ ). Sufficiently reducing denominator  $\hbar$  as well makes *fudge factor*  $g_s$  obsolete.

## Probable Cause

Two related mysteries of fermion spin remain.

1. While any two spin components have equal magnitude, their signs ( $\pm$ ) vary, correlating probabilistically with the angle separating them. An essential, but so far *abstract*, “probability amplitude” ( $a$ ) is strangely considered the *square root* of that probability ( $P$ ).

*“The probability of an event is represented by the square of an arrow [probability amplitude].”<sup>26</sup>*

*“The [probability] amplitude arrows are fundamental to the description of the world given by quantum theory. No satisfactory reason has been given for why they are needed.”<sup>27</sup>*

*“There have been many attempts to derive the Born rule from the other assumptions of quantum mechanics, with inconclusive results. ... probability is equal to the amplitude-squared”<sup>28</sup>*

*“These [probability amplitudes] are extremely abstract, and it is not at all obvious what their physical significance is.”<sup>29</sup>*

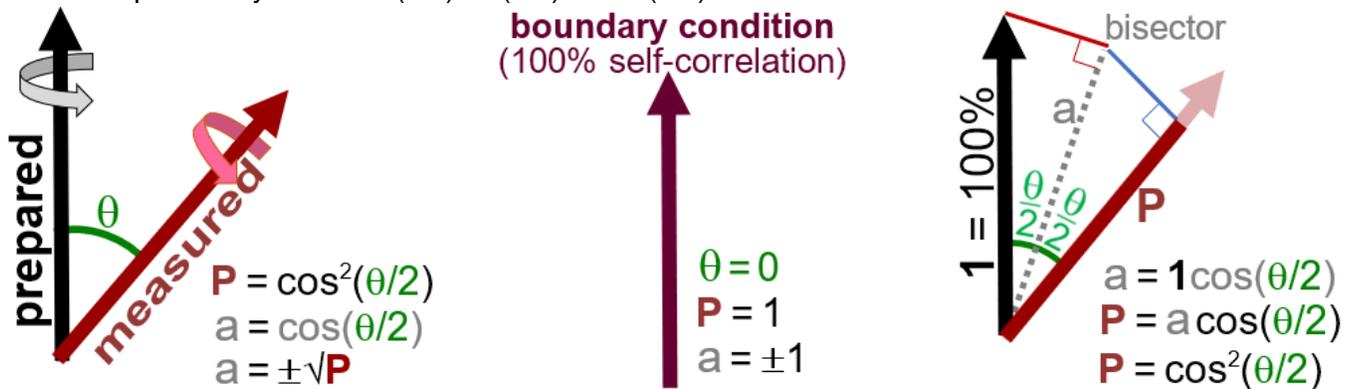
2. The probability amplitude sees a  $720^\circ$  horizon.

*“[The] physical effects of the difference between the rotation of a spin- $\frac{1}{2}$  particle by  $360^\circ$  as compared with  $720^\circ$  have been experimentally observed in classic experiments in neutron interferometry.”<sup>30</sup>*

One might guess that  $720^\circ$  relates to the  $4\pi$ , noted earlier for solid angles. But solid angles range to  $4\pi$  *steradians* (square radians of *area*), while  $720^\circ$  refers to  $4\pi$  *radians* (of arc *length*).

Experimentally, the sign of a prepared spin will correlate with that of a subsequently-measured spin, at angle  $\theta$ , with probability ( $P$ ) such that:  $P = \cos^2(\theta/2)$ . If amplitude ( $a$ ) has a real representation, it will be confined in a boundary condition from which to generalize (Fig. 8).

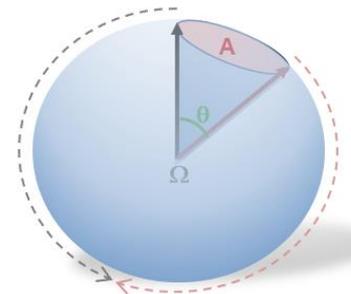
For example, amplitude  $a$  *must* coincide with the prepared spin when the subsequently measured spin has the same axis (i.e.  $\theta = 0$ ). More generally, as the *half-angle* specification hints,  $a$  is recognized on the angle *bisector*. Instead of amplitude as a *square root* of probability, it should be viewed as a probability in its own right, the projection of 100% self-correlation on the bisector, i.e.  $a = 1\cos(\theta/2)$ . That value is in turn, projected onto the subsequently-measured component, which results in the observed correlation probability:  $P = 1\cos(\theta/2)\cos(\theta/2) = \cos^2(\theta/2)$ .



**Fig. 8.** Left: The half-angle correlation of prepared and subsequently measured spin components entails a  $720^\circ$  range for a yet, *unidentified* “probability amplitude” ( $a$ ). Center: Amplitude  $a$  is pinned down in the boundary condition of **100%** self-correlation. Right:  $a$  is spin probability **1**, projected on the bisector. It is the spin correlation probability for the *unmeasured* bisector. Probability **P** is in turn, the projection of  $a$  on the subsequently-measured axis.

Being on the angle bisector, probability amplitude  $a$  has the mysterious property of existing, while *never* directly measurable. To do so would make it the subsequently-measured component, which immediately redefines the bisector. The amplitude is thus always out of reach, as is the half-way point of Zeno’s dichotomy paradox.<sup>31</sup>

With equal magnitude spin components as sides, the triangle they describe is isosceles. Its altitude is the probability amplitude. Both change sign when angle  $\theta$  crosses  $180^\circ$  (becomes convex). At  $360^\circ$  the value is  $-1$ , which continues back to  $+1$  at  $720^\circ$ .



**Fig. 9** An equator of solid angle  $\Omega$  is described by  $\theta$ , for which altitude turns negative after  $180^\circ$ .

## Simple as Riding a Bi-Cycle

Having previously dealt with bosons,<sup>14</sup> interval-time coordinates here provide a Euclidean lens which clearly reveals fermion spin by axis and plane of rotation. With that,  $4\pi$  seems unavoidable, whether  $4\pi$  steradians of solid angular rotation, the combined ranges ( $2\pi$  each) of two spherical coordinates or the  $4\pi$  radians of probability amplitude. It seems awkward at first, but with practice, it becomes second nature, like riding a bike.



Bohr stated, "...however far [quantum] phenomena transcend the scope of classical physical explanation, the account of all evidence must be expressed in classical terms."<sup>32</sup>

Quantum spin is **classical** spin about a **non-classical** axis. It is **chronaxial spin**.

Wheeler said, "Behind it all is surely an idea so simple, so beautiful, that when we grasp it - in a decade, a century, or a millennium - we will all say to each other, how could it have been otherwise? How could we have been so stupid for so long?"<sup>33</sup>

We don't have to be "stupid" about quantum spin anymore.



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