

General relativity and representation of solutions

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ABSTRACT

In the general relativity theory, we find the representation of the gravity field equation and solutions. We treat the representation of Schwarzschild solution, Reissner-Nordstrom solution, Kerr-Newman solution, Robertson-Walker solution. Specially, Robertson-Walker solution is a uniqueness.

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1. Introduction

In the general relativity theory, our article's aim is that we find the representation of the gravity field equation and solutions.

First, the gravity potential $g_{\mu\nu}$ is

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu \quad (1)$$

In gravity potential $g_{\mu\nu}$, we introduce tensor $f_{\mu\nu}$ and scalar K .

$$f_{\mu\nu} = Kg_{\mu\nu}, \quad \frac{\partial K}{\partial x^\lambda} = 0$$

$$\begin{aligned} ds^{12} &= f_{\mu\nu} dx^\mu dx^\nu = \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu = \bar{g}_{\mu\nu} \frac{\partial \bar{x}^\mu}{\partial x'^\alpha} \frac{\partial \bar{x}^\nu}{\partial x'^\beta} dx'^\alpha dx'^\beta \\ &= \bar{g}'_{\alpha\beta} d\bar{x}'^\alpha d\bar{x}'^\beta = f'_{\alpha\beta} dx'^\alpha dx'^\beta \end{aligned}$$

$$\bar{g}'_{\alpha\beta} = \bar{g}_{\mu\nu} \frac{\partial \bar{x}^\mu}{\partial \bar{x}'^\alpha} \frac{\partial \bar{x}^\nu}{\partial \bar{x}'^\beta}, \quad f'_{\alpha\beta} = f_{\mu\nu} \frac{\partial x^\mu}{\partial x'^\alpha} \frac{\partial x^\nu}{\partial x'^\beta} \quad (2)$$

In inverse gravity potential $g^{\mu\nu}$,

$$f^{\mu\nu} f_{\mu\nu} = \delta_\mu^\nu = \left(\frac{1}{K} g^{\mu\nu}\right) (Kg_{\mu\nu}), \quad f^{\mu\nu} = \frac{1}{K} g^{\mu\nu} \quad (3)$$

In Christoffel symbol $\Gamma^\rho_{\mu\nu}$,

$$\begin{aligned} \Gamma^{\rho}_{\mu\nu} &= \frac{1}{2} f^{\rho\lambda} \left(\frac{\partial f_{\mu\lambda}}{\partial x^\nu} + \frac{\partial f_{\nu\lambda}}{\partial x^\mu} - \frac{\partial f_{\mu\nu}}{\partial x^\lambda} \right) \\ &= \frac{1}{2} \left(\frac{1}{K} g^{\rho\lambda} \right) \left(K \frac{\partial g_{\mu\lambda}}{\partial x^\nu} + K \frac{\partial g_{\nu\lambda}}{\partial x^\mu} - K \frac{\partial g_{\mu\nu}}{\partial x^\lambda} \right) = \Gamma^{\rho}_{\mu\nu} \end{aligned}$$

$$\bar{\Gamma}^{\rho}_{\mu\nu} = \frac{1}{2} \bar{g}^{\rho\lambda} \left(\frac{\partial \bar{g}_{\mu\lambda}}{\partial \bar{x}^\nu} + \frac{\partial \bar{g}_{\nu\lambda}}{\partial \bar{x}^\mu} - \frac{\partial \bar{g}_{\mu\nu}}{\partial \bar{x}^\lambda} \right) = \frac{1}{\sqrt{K}} \Gamma^{\rho}_{\mu\nu} \quad (4)$$

Therefore, in the curvature tensor $R^{\rho}_{\mu\nu\lambda}$,

$$\begin{aligned} R^{\rho}_{\mu\nu\lambda} &= \frac{\partial \Gamma^{\rho}_{\mu\nu}}{\partial x^\lambda} - \frac{\partial \Gamma^{\rho}_{\mu\lambda}}{\partial x^\nu} + \Gamma^{\sigma}_{\mu\nu} \Gamma^{\rho}_{\sigma\lambda} - \Gamma^{\sigma}_{\mu\lambda} \Gamma^{\rho}_{\sigma\nu} \\ &= \frac{\partial \Gamma^{\rho}_{\mu\nu}}{\partial x^\lambda} - \frac{\partial \Gamma^{\rho}_{\mu\lambda}}{\partial x^\nu} + \Gamma^{\sigma}_{\mu\nu} \Gamma^{\rho}_{\sigma\lambda} - \Gamma^{\sigma}_{\mu\lambda} \Gamma^{\rho}_{\sigma\nu} = R^{\rho}_{\mu\nu\lambda} \end{aligned}$$

$$\begin{aligned}\bar{R}^\rho{}_{\mu\nu\lambda} &= \frac{\partial \bar{\Gamma}^\rho{}_{\mu\nu}}{\partial \bar{X}^\lambda} - \frac{\partial \bar{\Gamma}^\rho{}_{\mu\lambda}}{\partial \bar{X}^\nu} + \bar{\Gamma}^\sigma{}_{\mu\nu} \bar{\Gamma}^\rho{}_{\sigma\lambda} - \bar{\Gamma}^\sigma{}_{\mu\lambda} \bar{\Gamma}^\rho{}_{\sigma\nu} \\ &= \frac{1}{K} \left(\frac{\partial \Gamma^\rho{}_{\mu\nu}}{\partial X^\lambda} - \frac{\partial \Gamma^\rho{}_{\mu\lambda}}{\partial X^\nu} + \Gamma^\sigma{}_{\mu\nu} \Gamma^\rho{}_{\sigma\lambda} - \Gamma^\sigma{}_{\mu\lambda} \Gamma^\rho{}_{\sigma\nu} \right) = \frac{1}{K} R^\rho{}_{\mu\nu\lambda} \quad (5)\end{aligned}$$

In Ricci tensor $R_{\mu\nu}$,

$$R^\lambda{}_{\mu\nu} = R^{\lambda\rho}{}_{\mu\rho\nu} = R^\rho{}_{\mu\rho\nu} = R_{\mu\nu}, \quad \bar{R}_{\mu\nu} = \bar{R}^\rho{}_{\mu\rho\nu} = \frac{1}{K} R^\rho{}_{\mu\rho\nu} = \frac{1}{K} R_{\mu\nu} \quad (6)$$

In curvature scalar R

$$\begin{aligned}R^\lambda{}_{\lambda} &= f^{\mu\nu} R^\lambda{}_{\mu\nu} = \frac{1}{K} g^{\mu\nu} R_{\mu\nu} = \frac{1}{K} R \\ \bar{R} &= \bar{g}^{\mu\nu} \bar{R}_{\mu\nu} = \frac{1}{K} g^{\mu\nu} R_{\mu\nu} = \frac{1}{K} R \quad (7)\end{aligned}$$

Hence, in the gravity field equation of Einstein,

$$\begin{aligned}R^\lambda{}_{\mu\nu} - \frac{1}{2} f_{\mu\nu} R^\lambda{}_{\lambda} &= R_{\mu\nu} - \frac{1}{2} K g_{\mu\nu} \left(\frac{1}{K} R \right) \\ &= R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = -\frac{8\pi G}{c^4} T_{\mu\nu} \\ \bar{R}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{R} &= \frac{1}{K} \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) \\ &= -\frac{8\pi G}{c^4} \frac{1}{K} T_{\mu\nu} = -\frac{8\pi G}{c^4} \bar{T}_{\mu\nu} \quad (8)\end{aligned}$$

In Newtonian approximation, Energy-momentum tensor $T^\lambda{}_{\mu\nu}$ is

$$\nabla^2 f_{00} = \nabla^2 K g_{00} \approx -\frac{8\pi G}{c^4} K T_{00} = -\frac{8\pi G}{c^4} T^\lambda{}_{\lambda 00} \quad (9)$$

$$\rho c^2 = T_{00}, \quad K \rho c^2 = T^\lambda{}_{\lambda 00}$$

$$\bar{T}_{\mu\nu} = \frac{1}{K} T_{\mu\nu} \quad (10)$$

$$\bar{\nabla}^2 \bar{g}_{00} = \frac{1}{K} \nabla^2 g_{00} \approx -\frac{8\pi G}{c^4} \frac{1}{K} T_{00} = -\frac{8\pi G}{c^4} \bar{T}_{00} \quad (11)$$

$$\rho c^2 = T_{00}, \quad \frac{1}{K} \rho c^2 = \bar{T}_{00}$$

$$T'_{\mu\nu} = KT_{\mu\nu}, \quad \frac{1}{K}T_{\mu\nu} = \bar{T}_{\mu\nu} \quad (12)$$

Einstein's gravity field equation is

$$\begin{aligned} R'_{\mu\nu} - \frac{1}{2}f_{\mu\nu}R' &= R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -\frac{8\pi G}{c^4}T_{\mu\nu} = -\frac{8\pi G}{c^4}\frac{T'_{\mu\nu}}{K} \\ \bar{R}_{\mu\nu} - \frac{1}{2}\bar{g}_{\mu\nu}\bar{R} &= \frac{1}{K}(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R) = -\frac{1}{K}\frac{8\pi G}{c^4}T_{\mu\nu} = -\frac{8\pi G}{c^4}\bar{T}_{\mu\nu} \end{aligned} \quad (13)$$

Therefore, tensor $f_{\mu\nu}$ satisfy new gravity field equation of Einstein.

$$\begin{aligned} f^{\mu\nu}[R'_{\mu\nu} - \frac{1}{2}f_{\mu\nu}R'] &= \frac{1}{K}g^{\mu\nu}[R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R] = -\frac{8\pi G}{c^4}\frac{1}{K}g^{\mu\nu}T_{\mu\nu} = -\frac{8\pi G}{c^4}\frac{1}{K}T^{\lambda}_{\lambda} \\ \bar{g}^{\mu\nu}[\bar{R}_{\mu\nu} - \frac{1}{2}\bar{g}_{\mu\nu}\bar{R}] &= \frac{1}{K}g^{\mu\nu}[R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R] = -\frac{8\pi G}{c^4}\frac{1}{K}g^{\mu\nu}T_{\mu\nu} = -\frac{8\pi G}{c^4}\frac{1}{K}T^{\lambda}_{\lambda} \\ &= -\frac{8\pi G}{c^4}\bar{g}^{\mu\nu}\bar{T}_{\mu\nu} = -\frac{8\pi G}{c^4}\bar{T}^{\lambda}_{\lambda} \quad \rightarrow \\ -\bar{R} &= -\frac{1}{K}R = -\frac{8\pi G}{c^4}\frac{1}{K}T^{\lambda}_{\lambda} = -\frac{8\pi G}{c^4}\bar{T}^{\lambda}_{\lambda} \\ &= -\frac{8\pi G}{c^4}f^{\mu\nu}\frac{T'_{\mu\nu}}{K} = -\frac{8\pi G}{c^4}\frac{1}{K}T'^{\lambda}_{\lambda} \\ \rightarrow -R' &= -\frac{1}{K}R = -\frac{8\pi G}{c^4}\frac{1}{K}T^{\lambda}_{\lambda} = -\frac{8\pi G}{c^4}\frac{1}{K}T'^{\lambda}_{\lambda}, \\ T'^{\lambda}_{\lambda} &= T^{\lambda}_{\lambda}, \quad \frac{1}{K}T^{\lambda}_{\lambda} = \bar{T}^{\lambda}_{\lambda} \end{aligned} \quad (14)$$

Ricci tensor is

$$\begin{aligned} R'_{\mu\nu} &= R_{\mu\nu} = -\frac{8\pi G}{c^4}(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T^{\lambda}_{\lambda}) = -\frac{8\pi G}{c^4}(\frac{T'_{\mu\nu}}{K} - \frac{1}{2}\frac{f_{\mu\nu}}{K}T'^{\lambda}_{\lambda}) \\ \bar{R}_{\mu\nu} &= \frac{1}{K}R_{\mu\nu} = -\frac{8\pi G}{c^4}\frac{1}{K}(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T^{\lambda}_{\lambda}) = -\frac{8\pi G}{c^4}(\bar{T}_{\mu\nu} - \frac{1}{2}\bar{g}_{\mu\nu}\bar{T}^{\lambda}_{\lambda}) \end{aligned}$$

$$f_{\mu\nu} = Kg_{\mu\nu}, \quad \frac{\partial K}{\partial x^{\lambda}} = 0$$

$$ds^2 = g_{\mu\nu}dx^{\mu}dx^{\nu}, \quad ds'^2 = f_{\mu\nu}dx^{\mu}dx^{\nu} = \bar{g}_{\mu\nu}d\bar{x}^{\mu}d\bar{x}^{\nu} \quad (15)$$

2. Weak gravity field approximation.

Weak gravity field approximation is

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad \bar{g}_{\mu\nu} = \bar{\eta}_{\mu\nu} + \bar{h}_{\mu\nu}$$

$$R_{\mu\nu} = R'_{\mu\nu} = -\frac{8\pi G}{c^4} (T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T^{\lambda}_{\lambda})$$

$$\frac{1}{K} R_{\mu\nu} = \bar{R}_{\mu\nu} = -\frac{8\pi G}{c^4} (\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}^{\lambda}_{\lambda})$$

$$R_{\mu\nu} = -\frac{8\pi G}{c^4} S_{\mu\nu}, \quad \bar{R}_{\mu\nu} = -\frac{8\pi G}{c^4} \bar{S}_{\mu\nu}$$

$$S_{\mu\nu} = T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T^{\lambda}_{\lambda}, \quad \bar{S}_{\mu\nu} = \bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}^{\lambda}_{\lambda} \quad (16)$$

The solution is

$$h_{\mu\nu}(t, \vec{x}) = \frac{4G}{c^2} \int d^4 x' \frac{S_{\mu\nu}(t - |\vec{x} - \vec{x}'|, \vec{x}')}{|\vec{x} - \vec{x}'|}, \quad \int d^3 x T_{00} = M$$

$$\bar{h}_{\mu\nu}(t, \vec{x}) = \frac{4G}{c^2} \int d^4 \bar{x}' \frac{\bar{S}_{\mu\nu}(\bar{t} - |\vec{x} - \vec{x}'|, \vec{x}')}{|\vec{x} - \vec{x}'|},$$

$$\int d^3 \bar{x} \bar{T}_{00} = \int K \sqrt{K} d^3 x \frac{1}{K} T_{00} = \sqrt{K} M = \bar{M}$$

$$\bar{h}_{00}(\vec{x}) \approx \frac{4G}{\bar{r}c^2} \int d^3 \bar{x}' [\bar{T}_{00} - \frac{1}{2} \bar{T}_{00}] = \frac{2\sqrt{K}GM}{\bar{r}c^2},$$

$$\int d^3 \bar{x} \bar{T}_{00} = \int d^3 x K \sqrt{K} \bar{T}_{00} = \int d^3 x \sqrt{K} T_{00} = \bar{M}$$

$$\bar{h}_{ij}(\vec{x}) \approx \frac{4G}{\bar{r}c^2} \int d^3 \bar{x}' [\frac{1}{2} \delta_{ij} \bar{T}_{00}] = \frac{2\sqrt{K}GM}{\bar{r}c^2} \delta_{ij}, \quad T_{\mu\nu} = K \bar{T}_{\mu\nu} \quad (17)$$

The proper distance is

$$-ds^2 = c^2 d\tau^2 = -g_{\mu\nu} dx^\mu dx^\nu \approx (1 - \frac{2GM}{rc^2}) c^2 dt^2 - (1 + \frac{2GM}{rc^2}) \delta_{ij} dx^i dx^j$$

$$-ds'^2 = -K ds^2 = K c^2 d\tau^2 = -K g_{\mu\nu} dx^\mu dx^\nu$$

$$\begin{aligned} &\approx K (1 - \frac{2GM}{rc^2}) c^2 dt^2 - K (1 + \frac{2GM}{rc^2}) \delta_{ij} dx^i dx^j \\ &= (1 - \frac{2\sqrt{K}GM}{\bar{r}c^2}) c^2 d\bar{t}^2 - (1 + \frac{2\sqrt{K}GM}{\bar{r}c^2}) \delta_{ij} d\bar{x}^i d\bar{x}^j \\ &= (1 - \frac{2G\bar{M}}{\bar{r}c^2}) c^2 d\bar{t}^2 - (1 + \frac{2G\bar{M}}{\bar{r}c^2}) \delta_{ij} d\bar{x}^i d\bar{x}^j \end{aligned}$$

$$= \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu$$

$$\sqrt{K}t = \bar{t}, \sqrt{K}x^i = \bar{x}^i, \sqrt{K}x^j = \bar{x}^j, \sqrt{K}r = \bar{r}, \sqrt{KM} = \bar{M} \quad (18)$$

3. The other representation in Schwarzschild solution, Reissner-Nodstrom solution, Kerr-Newman solution and Robertson-Walker solution

Schwarzschild solution (vacuum solution) is

$$R_{\mu\nu} = R^{\mu\nu} = 0$$

$$ds^2 = -c^2 \left(1 - \frac{2GM}{rc^2}\right) dt^2 + \frac{dr^2}{1 - \frac{2GM}{rc^2}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \quad (19)$$

The other representation of Schwarzschild solution is

$$\begin{aligned} ds^{12} &= f_{\mu\nu} dx^\mu dx^\nu = Kg_{\mu\nu} dx^\mu dx^\nu = K \cdot ds^2 \\ &= -c^2 K \left(1 - \frac{2GM}{rc^2}\right) dt^2 + \frac{K dr^2}{1 - \frac{2GM}{rc^2}} + Kr^2 d\theta^2 + Kr^2 \sin^2 \theta d\phi^2 \\ &= -c^2 \left(1 - \frac{2\sqrt{K}GM}{\bar{r}c^2}\right) d\bar{t}^2 + \frac{d\bar{r}^2}{1 - \frac{2\sqrt{K}GM}{\bar{r}c^2}} + \bar{r}^2 d\bar{\theta}^2 + \bar{r}^2 \sin^2 \bar{\theta} d\bar{\phi}^2 \\ &= -c^2 \left(1 - \frac{2G\bar{M}}{\bar{r}c^2}\right) d\bar{t}^2 + \frac{d\bar{r}^2}{1 - \frac{2G\bar{M}}{\bar{r}c^2}} + \bar{r}^2 d\bar{\theta}^2 + \bar{r}^2 \sin^2 \bar{\theta} d\bar{\phi}^2 = \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu \\ \sqrt{K}t &= \bar{t}, \sqrt{K}r = \bar{r}, \theta = \bar{\theta}, \phi = \bar{\phi}, \sqrt{KM} = \bar{M} \end{aligned} \quad (20)$$

Reissner-Nodstrom solution is

$$\begin{aligned} ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\ &= -c^2 \left(1 - \frac{2GM}{rc^2} + \frac{kGQ^2}{r^2 c^4}\right) dt^2 + \frac{dr^2}{1 - \frac{2GM}{rc^2} + \frac{kGQ^2}{r^2 c^4}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \end{aligned} \quad (21)$$

The other representation of Reissner-Nodstrom solution is

$$ds^{12} = f_{\mu\nu} dx^\mu dx^\nu = Kg_{\mu\nu} dx^\mu dx^\nu = K \cdot ds^2$$

$$\begin{aligned}
&= -Kc^2\left(1 - \frac{2GM}{rc^2} + \frac{kGQ^2}{r^2c^4}\right)dt^2 + \frac{Kdr^2}{1 - \frac{2GM}{rc^2} + \frac{kGQ^2}{r^2c^4}} + Kr^2d\theta^2 + Kr^2\sin^2\theta d\phi^2 \\
&= -c^2\left(1 - \frac{2\sqrt{K}GM}{\bar{r}c^2} + \frac{KkGQ^2}{\bar{r}^2c^4}\right)d\bar{t}^2 + \frac{d\bar{r}^2}{1 - \frac{2\sqrt{K}GM}{\bar{r}c^2} + \frac{KkGQ^2}{\bar{r}^2c^4}} + \bar{r}^2d\bar{\theta}^2 + \bar{r}^2\sin^2\bar{\theta}d\bar{\phi}^2 \\
&= -c^2\left(1 - \frac{2G\bar{M}}{\bar{r}c^2} + \frac{kG\bar{Q}^2}{\bar{r}^2c^4}\right)d\bar{t}^2 + \frac{d\bar{r}^2}{1 - \frac{2G\bar{M}}{\bar{r}c^2} + \frac{kG\bar{Q}^2}{\bar{r}^2c^4}} + \bar{r}^2d\bar{\theta}^2 + \bar{r}^2\sin^2\bar{\theta}d\bar{\phi}^2 \\
&= \bar{g}_{\mu\nu}d\bar{x}^\mu d\bar{x}^\nu
\end{aligned}$$

$$\sqrt{K}t = \bar{t}, \sqrt{K}r = \bar{r}, \theta = \bar{\theta}, \phi = \bar{\phi}, \sqrt{KM} = \bar{M}, KQ^2 = \bar{Q}^2 \quad (22)$$

Kerr-Newman solution is

$$\begin{aligned}
ds^2 &= g_{\mu\nu}dx^\mu dx^\nu \\
&= -c^2\left(1 - \frac{2c^2GMr - kGQ^2}{c^4\Sigma}\right)dt^2 + 2(2c^2MGr - kGQ^2)\frac{a\sin^2\theta}{c^4\Sigma}cdtd\phi \\
&\quad - \frac{c^4\Sigma}{r^2 - c^22GMr + a^2 + kGQ^2}dr^2 - \Sigma d\theta^2 \\
&\quad - \sin^2\theta[r^2 + a^2 + (2c^2GMr - kGQ^2)\frac{a^2\sin^2\theta}{c^4\Sigma}]d\phi^2
\end{aligned}$$

$$\Sigma = r^2 + a^2 \cos^2\theta \quad (23)$$

The other representation of Kerr-Newman solution is

$$\begin{aligned}
ds'^2 &= f_{\mu\nu}dx^\mu dx^\nu = Kg_{\mu\nu}dx^\mu dx^\nu = K \cdot ds^2 \\
&= -Kc^2\left(1 - \frac{2c^2GMr - kGQ^2}{c^4\Sigma}\right)dt^2 + 2K(2c^2MGr - kGQ^2)\frac{a\sin^2\theta}{c^4\Sigma}cdtd\phi \\
&\quad - \frac{K\Sigma c^4}{r^2 - 2c^2GMr + a^2 + kGQ^2}dr^2 - K\Sigma d\theta^2 \\
&\quad - K\sin^2\theta[r^2 + a^2 + (2c^2GMr - kGQ^2)\frac{a^2\sin^2\theta}{c^4\Sigma}]d\phi^2 \\
&= -c^2\left(1 - \frac{2c^2G\sqrt{KM}\sqrt{K}r - kGKQ^2}{K\Sigma c^4}\right)d\bar{t}^2 + 2(2c^2\sqrt{K}MG\sqrt{K}r - kGKQ^2)\frac{\sqrt{K}a\sin^2\theta}{K\Sigma c^4}cd\sqrt{K}td\bar{\phi}
\end{aligned}$$

$$\begin{aligned}
& - \frac{k\Sigma c^4}{Kr^2 - 2c^2 G\sqrt{KM}\sqrt{Kr} + Ka^2 + kGKQ} d(\sqrt{Kr})^2 - k\Sigma d\theta^2 \\
& - \sin^2 \theta [Kr^2 + Ka^2 + (2c^2 G\sqrt{KM}\sqrt{Kr} - kGKQ^2) \frac{Ka^2 \sin^2 \theta}{k\Sigma c^4}] d\phi^2 \\
= & -c^2 \left(1 - \frac{2c^2 G\bar{M}\bar{r} - kG\bar{Q}^2}{\bar{\Sigma} c^4}\right) d\bar{t}^2 + 2(2c^2 \bar{M}G\bar{r} - kG\bar{Q}^2) \frac{\bar{a} \sin^2 \bar{\theta}}{\bar{\Sigma} c^4} c d\bar{t} d\bar{\phi} \\
& - \frac{\bar{\Sigma} c^4}{\bar{r}^2 - 2c^2 G\bar{M}\bar{r} + \bar{a}^2 + kG\bar{Q}^2} d\bar{r}^2 - \bar{\Sigma} d\bar{\theta}^2 \\
& - \sin^2 \bar{\theta} [\bar{r}^2 + \bar{a}^2 + (2c^2 G\bar{M}\bar{r} - kG\bar{Q}^2) \frac{\bar{a}^2 \sin^2 \bar{\theta}}{c^4 \bar{\Sigma}}] d\bar{\phi}^2 \\
= & \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu \\
\bar{\Sigma} = & k\Sigma = Kr^2 + Ka^2 \cos^2 \theta = \bar{r}^2 + \bar{a}^2 \cos^2 \bar{\theta} \\
\sqrt{K}t = & \bar{t}, \sqrt{K}r = \bar{r}, \theta = \bar{\theta}, \phi = \bar{\phi}, \sqrt{KM} = \bar{M}, KQ^2 = \bar{Q}^2, \sqrt{Ka} = \bar{a}
\end{aligned} \tag{24}$$

Robertson-Walker solution is

$$\begin{aligned}
ds^2 &= g_{\mu\nu} dx^\mu dx^\nu \\
= & -c^2 dt^2 + \Omega^2(t) \left[\frac{dr^2}{1 - Kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right]
\end{aligned} \tag{25}$$

The other representation of Robertson-Walker solution is by the other scalar K^1 ,

$$\begin{aligned}
ds'^2 &= f_{\mu\nu} dx^\mu dx^\nu = K^1 g_{\mu\nu} dx^\mu dx^\nu = K^1 ds^2 \\
= & -K^1 c^2 dt^2 + \Omega^2(t) \left[\frac{K^1 dr^2}{1 - k \frac{K^1 r^2}{K^1}} + K^1 r^2 d\theta^2 + K^1 r^2 \sin^2 \theta d\phi^2 \right] \\
= & -c^2 d\bar{t}^2 + \bar{\Omega}^2(\bar{t}) \left[\frac{d\bar{r}^2}{1 - \frac{k\bar{r}^2}{K^1}} + \bar{r}^2 d\bar{\theta}^2 + \bar{r}^2 \sin^2 \bar{\theta} d\bar{\phi}^2 \right] \\
= & -c^2 d\bar{t}^2 + \bar{\Omega}^2(\bar{t}) \left[\frac{d\bar{r}^2}{1 - K^1 \bar{r}^2} + \bar{r}^2 d\bar{\theta}^2 + \bar{r}^2 \sin^2 \bar{\theta} d\bar{\phi}^2 \right] = \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu \\
\sqrt{K^1}t = & \bar{t}, \Omega(t) = \bar{\Omega}(\bar{t}), \\
\sqrt{K^1}r = & \bar{r}, \theta = \bar{\theta}, \phi = \bar{\phi}
\end{aligned}$$

$$k = (0,1,-1), \quad k' = \frac{k}{K'} = \left(0, \frac{1}{K'}, -\frac{1}{K'}\right) \quad (26)$$

Hence, $K' = 1$, In this time, dS^2 is an uniqueness.

4. Conclusion

We find the other representation of solutions in the General relativity theory. In this time, Robertson-Walker solution is an uniqueness.

Reference

- [1]S.Weinberg,Gravitation and Cosmology(John wiley & Sons,Inc,1972)
- [2]W.Rindler, Am.J.Phys.**34**.1174(1966)
- [3]P.Bergman,Introduction to the Theory of Relativity(Dover Pub. Co.,Inc., New York,1976),Chapter V
- [4]C.Misner, K,Thorne and J. Wheeler, Gravitation(W.H.Freedman & Co.,1973)
- [5]S.Hawking and G. Ellis,The Large Scale Structure of Space-Time(Cam-bridge University Press,1973)
- [6]R.Adler,M.Bazin and M.Schiffer,Introduction to General Relativity(McGraw-Hill,Inc.,1965)
- [7]A.Miller, Albert Einstein's Special Theory of Relativity(Addison-Wesley Publishing Co., Inc., 1981)
- [8]W.Rindler, Special Relativity(2nd ed., Oliver and Boyd, Edinburg,1966)
- [9]F. Shojai and A. Shojai,"The equivalence principle and the relative velocity of local inertial frames";arXiv:1505.06691v1[gr-qc]
- [10]G.Birkoff, Relativity and Modern Physics (Harvard University Press,1923),p253
- [11]A.Raychaudhuri, Theoretical Cosmology(Oxpond University Press,1979)
- [12]E. Kasner, Am. J. Math. 43, 217(1921)