Introduction

In present article alternative (to Standard Model) hypothesis of structure of electron, proton and neutron is suggested. The others elementary particles (except photon and neutrino) are not stable and they are considered as unsteady soliton-similar formations. In series of experiments indirect confirmations of existence of quarks were obtained, for instance in experiments by scattering of electrons at nuclei, performed at Stanford linear accelerator by R. Hofshtadter, look for instance [1]. At that, experiments by elastic and deeply inelastic scattering gave quite different results: in first case take place pattern of scattering at lengthy object, in second case is pattern of scattering at "point" centers, that is interpreted as confirmations of existence of quarks. However what "point" formations appear only in deeply inelastic scattering don’t may be an evidence of quarks existence, because to above-mentioned fact may be given and another explanations: in moment of birth of new particles, which take place in deeply inelastic scattering, structure of nucleon change, it sharply diminish in volume, but after appearance of new particles nucleon return to initial state. Or process of birth of new particles occur in "point" volume inside nucleon and these energy "point" centers disappear after completion of process particles birth. And fact that experiments by elastic scattering gave pattern of scattering at lengthy object prove inexistence of quarks in nucleus. In theory of Standard (quarkual) Model come into at least 20 parameters artificially introduced from outside, such as "colour" of particles, "aroma" etc., that is its fundamental demerit. Theoretical work, which is present here, has no demerits of Standard Model, it completely describe structure of elementary particles therefore it can help in discovery new ways of making energy, elaboration perfectly new devices for its production and to achieve progress in such fields as nuclear power engineering, nanotechnology, high-powerful lasers, clean energy and others.
Abstract

In paper, which is submitted, electron, proton and neutron are considered as spherical areas, inside which monochromatic electromagnetic wave of corresponding frequency spread along parallels, at that along each parallel exactly half of wave length for electron and proton and exactly one wave length for neutron is kept within, thus this is rotating soliton. This is caused by presence of spatial dispersion and anisotropy of strictly defined type inside the particles. Electric field has only radial component, and magnetic field - only meridional component. By solution of corresponding edge task, functions of distribution of electromagnetic field inside the particles and on their boundary surfaces were obtained. Integration of distribution functions of electromagnetic field through volume of the particles lead to system of algebraic equations, solution of which give all basic parameters of particles: charge, rest energy, mass, radius, magnetic moment and spin.

Keywords:
structure of elementary particles; structure of matter; theory of elementary particles; electron; proton; neutron; nuclei; electromagnetic field; atom; microcosm; elementary particles; fundamental interactions.

1. Rotating monochromatic electromagnetic wave.

Let us write down Maxwell’s equations in spherical coordinates supposing that:

1) there are no losses;

2) only $\hat{E}_r, \hat{H}_\theta, \hat{j}_\phi, \hat{\rho}$ are not equal to zero.

\[
\frac{1}{r} \left( \frac{\partial}{\partial r} (r \hat{H}_\theta) \right) = \hat{j}_\phi; \quad (1)
\]

\[
\frac{1}{r \sin \theta} \frac{\partial \hat{E}_r}{\partial \phi} = -i \omega \mu \hat{H}_\theta; \quad (2)
\]

\[
\frac{1}{r} \frac{\partial \hat{E}_r}{\partial \theta} = i \omega \mu \hat{H}_\phi = 0; \quad (3)
\]

\[
\frac{1}{r \sin \theta} \left( - \frac{\partial \hat{H}_\theta}{\partial \phi} \right) = i \omega \varepsilon \hat{E}_r; \quad (4)
\]

\[
\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \varepsilon \hat{E}_r) = \hat{\rho}; \quad (5)
\]
\[
\frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta (\mu \dot{H}_\theta)) = 0.
\] (6)

Here \( r, \theta, \varphi \) - spherical coordinates of the observation point; \( \dot{E}_r, \dot{H}_\theta \) - components of the electromagnetic field, \( j_\varphi \) - density of electric current, \( \dot{\rho} \) - volume charge density; \( \omega \) - circular frequency of field alteration, \( i \) - imaginary unit, \( \varepsilon \) – dielectric permittivity, \( \mu \) – magnetic permeability.

**Fig. 1**

Field components in the rotating electromagnetic wave.
Substituting the expression for $\dot{H}_\varphi$ from (2) in (4), we obtain:

$$\frac{\partial^2 \dot{E}_r}{\partial \varphi^2} + \varepsilon \mu \omega^2 r^2 \sin^2 \theta \dot{E}_r = 0; \quad (7)$$

$$\frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \dot{E}_r}{\partial \varphi^2} + \omega^2 \varepsilon \mu \dot{E}_r = 0; \quad (7')$$

This is Helmholtz homogeneous equation. Let us designate

$$\sqrt{\varepsilon \mu \omega} \sin \theta = k_i - \quad (7'^*)$$

wave number. General solution of Helmholtz equation:

$$\dot{E}_r = E_0 \, e^{-ik_i \varphi} + E_0 \, e^{ik_i \varphi}. \quad (8)$$

This expression describes two waves, moving to meet one another by circular trajectories, along the parallels. Pointing’s vector in each point is directed at tangent to the corresponding parallel.

Let us consider a wave, moving in positive direction $\varphi$.

$$\dot{E}_r = E_0 \, e^{-ik_i \varphi} \, F(r, \theta); \quad (9)$$

Here

$$k_i \varphi = \sqrt{\varepsilon / \mu} (\omega r \sin \theta) \varphi - \text{wave phase;}$$

$$k_i = \text{dimensionless analog of the wave number. If to introduce a wave number of traditional dimension (} \frac{1}{m}; \text{) }$$

$$\beta = \omega \sqrt{\varepsilon / \mu} = \frac{k_i}{r \sin \theta},$$

the wave phase will be written down as

$$k_i \varphi = \sqrt{\varepsilon / \mu} \omega (r \sin \theta) \varphi = \beta l,$$

where

$$l = (r \sin \theta) \varphi - \text{arc length along the corresponding parallel. In the considered case the wave number is a function of coordinates and frequency. Thus, the wave, which is described, can exist only}$$
at availability of spatial and frequency dispersion. Dispersion equations will be obtained below, apart from the already found expression (7').

From expression (2), taking into account (7") and (9), we have:

\[
H_\phi = \sqrt{\frac{\varepsilon \mu}{\omega^2}} e^{i\omega r \sin \theta} F(r, \theta) = \frac{E_0}{z} e^{i\theta} F(r, \theta) .
\]

(9')

For actual amplitudes:

\[
E_r = E_0 F(r, \theta) \sin k_\phi;
\]

(10)

\[
H_\phi = \frac{E_0}{z} F(r, \theta) \sin k_\phi .
\]

(10')

Here

\[
z = \sqrt{\frac{\mu}{\varepsilon}}
\]

means characteristic impedance.

The last expressions describe an electromagnetic wave, rotating around axis Z in positive direction \( \phi \). Conditions of self-consistency:

1) \( z = \text{const}; \)

2) along each parallel on the circle length, the integer number of half-waves must be kept within.

\[
2\pi r \sin \theta = n \frac{\lambda}{2};
\]

(11)

here \( \lambda = \frac{\nu}{f} \)

wave length, \( \nu \) - phase velocity of wave, \( f \) - frequency, \( n = 1,2,3\ldots \)

Let us consider the case when \( n = 1 \),

\[
2\pi r \sin \theta = \frac{\lambda}{2};
\]

\[
\nu = 2\omega r \sin \theta .
\]

(11')

Along each parallel, exactly half of wave length is kept within.

Phase velocity of wave is the function of frequency and distance up to the axis of rotation.

\[
\nu = \frac{1}{\sqrt{\varepsilon \mu}} = 2\omega r \sin \theta ;
\]
\[ \varepsilon \mu = \frac{1}{4\omega^2 r^2 \sin^2 \theta}; \quad (11^*) \]

\[ z = \frac{\mu}{\varepsilon}; \]

\[ \mu = \varepsilon z^2; \quad (11^*) \]

we are substituting in \((11^*)\):

\[ z = \frac{1}{2\omega r \varepsilon \sin \theta}; \quad (12) \]

\[ \varepsilon = \frac{1}{2\omega rz \sin \theta}. \quad (12') \]

From \((11^*)\) \(\varepsilon = \frac{\mu}{z^2};\)

we are substituting in \((12').\)

\[ \mu = \frac{z}{2\omega r \sin \theta}; \quad (12'^*) \]

\[ z = 2\mu \omega r \sin \theta. \quad (12'^*) \]

Taking into account \((8)\) and \((11^*)\)

\[ k_1 = \frac{\omega r \sin \theta}{2\omega r \sin \theta} = \frac{1}{2}. \]

Then

\[ E_r = E_\phi F(r, \theta) \sin \frac{\phi}{2}; \quad (13) \]

\[ H_\phi = \frac{E_\phi}{z} F(r, \theta) \sin \frac{\phi}{2}. \quad (13'^*) \]

Function \(\sin \frac{\phi}{2}\) is onevalued in angles interval \(0 \leq \phi \leq 2\pi\).

This situation can be interpreted as rotation of spherical coordinate system around axis z in positive direction \(\phi\) with angular velocity \(\frac{d\phi}{dt}\). Let us find it from the condition

\[ \omega t - \frac{\phi}{2} = \text{const.} \]

Having differentiated this expression on t, we receive

\[ \frac{d \phi}{dt} = 2\omega. \]
At the same time the electromagnetic field, about spherical coordinate system, is determined by expressions (13) and (13’).

Further from (3): as \( \dot{H}_\theta = 0 \),

\[
\begin{align*}
\dot{E}_r (\theta) &= \text{const} ; \\
E_r (\theta) &= \text{const} .
\end{align*}
\]

(14)

From equation (6)

\[
\frac{\partial}{\partial \theta} \left( \frac{\sin \theta ( z H_\theta )}{2 \omega r \sin \theta} \right) = 0
\]

follows

\[
\begin{align*}
\dot{H}_\theta (\theta) &= \text{const} ; \\
H_\theta (\theta) &= \text{const} .
\end{align*}
\]

(14’)

To receive field dependence from \( r : E_r (r) ; H_\theta (r) \), let us find solution of three-dimensional Helmholtz equation in spherical coordinates.

\[
\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial E_r}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial E_r}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta \partial \varphi^2} \frac{\partial^2 E_r}{\partial \varphi^2} + k^2 E_r = 0 .
\]

(15)

\( E_r \) does not depend from \( \theta \), look (14), therefore three-dimensional Helmholtz equation transfers into two-dimensional one.

\[
\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial E_r}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta \partial \varphi^2} \frac{\partial^2 E_r}{\partial \varphi^2} + k^2 E_r = 0 .
\]

(15’)

Let us suppose that

\[
k^2 = k_2^2 + k_3^2 ,
\]

now

\[
\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial E_r}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta \partial \varphi^2} \frac{\partial^2 E_r}{\partial \varphi^2} + k_2^2 E_r + k_3^2 E_r = 0 .
\]

(15’)

This equation can be satisfied, if

\[
\begin{align*}
\frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 E_r}{\partial \varphi^2} + k_2^2 E_r &= 0 ; \\
\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial E_r}{\partial r} \right) + k_3^2 E_r &= 0 .
\end{align*}
\]

(16),(17)

Thus, initial Helmholtz equation has split into the system of two equations. We substitute in these equations instead of \( E_r (r, \varphi) = f(r) g(\varphi) \),
(i.e. we are searching the solution as the product of two functions) and divide the first equation by \(f(r)\), and the second - by \(g(\phi)\). We receive

\[
\begin{align*}
\frac{d^2 g}{d \phi^2} + k_2^2 r^2 (\sin^2 \theta) g &= 0; \\
r^2 \frac{d^2 f}{d r^2} + 2r \frac{d f}{d r} + k_3^2 r^2 f &= 0.
\end{align*}
\]  

(18), (19)

Equations (16) and (18) are equivalent to equations (7) и (7’), which were received earlier from Maxwell’s equations, and

\[
k_2 = \omega \sqrt{\varepsilon \mu} = \frac{\omega}{v} = \frac{1}{2r \sin \theta};
\]

\[
k_1 = k_2 r \sin \theta = \frac{1}{2}.
\]

The solution of equation (18) was found earlier, look (13).

\[
g(\phi) = \sin \frac{\phi}{2}.
\]  

(20)

Let us copy (19) as:

\[
r^2 \frac{d^2 f}{d r^2} + 2r \frac{d f}{d r} + k_4^2 f = 0; \tag{19'}
\]

where

\[
k_4 = k_3 r.
\]

(19’) – centrally symmetric Helmholtz equation. Let us suggest,

\[
k_3 = \frac{\omega}{v_r},
\]

where \(v_r\) – phase velocity of electromagnetic wave in radial direction. As in the central symmetric equation angular dependence is absent, it is logical to assume that

\[
v_r = v = 2\omega r \sin \theta
\]

at \(\theta = \frac{\pi}{2}\); i.e.

\[
v_r = 2 \omega r;
\]

\[
k_3 = \frac{\omega}{v_r} = \frac{1}{2r}; \tag{21}
\]

\[
k_4 = k_3 r = \frac{1}{2}. \tag{21'}
\]

Instead of (19’), we are having
\[ r^2 \frac{d^2 f}{d r^2} + 2 r \frac{d f}{d r} + \frac{1}{4} f = 0. \]  \hspace{1cm} (19^*)

This is Euler equation, it has the solution

\[ f = r^{-\frac{1}{2}} (C_1 + C_2 \ln r). \]  \hspace{1cm} (22)

Let us converse expression (22).

\[ f = \frac{(C_1 + C_2 \ln r)}{\sqrt{r}} = \frac{\alpha}{\sqrt{r}} (C_3 + C \ln r) = \frac{\alpha}{\sqrt{r}} (1 + \ln C_s + \ln r^C) = \frac{\alpha}{\sqrt{r}} (1 + \ln C_s r^C). \]  \hspace{1cm} (22')

Here \( C_1 = \sqrt{a} \ C_3; \ C_2 = \sqrt{a} \ C; \ C_3 = 1 + \ln C_s; \ a - value \ of \ radius \ r, \ at \ which \ the \ rotating \ monochromatic \ electromagnetic \ wave \ ceases \ to \ exist, \ and \ E_r = E_0; \ f = 1; \) hence

\[ 1 + \ln C_s a^C = 1; \]
\[ \ln C_s a^C = 0; \]
\[ C_s a^C = 1; \]
\[ C_s = \frac{1}{a^C}. \]  \hspace{1cm} (22'')

In view of this,

\[ f = \frac{\alpha}{\sqrt{r}} (1 + \ln \left(\frac{r}{a}\right)^C) = \frac{\alpha}{\sqrt{r}} (1 + C \ln \frac{r}{a}). \]

Let us designate \( C = p; \) now

\[ f = \frac{\alpha}{\sqrt{r}} (1 + p \ln \frac{r}{a}). \]

Thus, for \( E_r \) we are having

\[ E_r = E_0 \ g(\phi) \ f(r) = E_0 \ \frac{\alpha}{\sqrt{r}} (1 + p \ln \frac{r}{a}) \sin \frac{\phi}{2}. \]  \hspace{1cm} (23)

At \( r \to \infty, \ E_r = 0; \ f = 0. \)

Really

\[ \lim_{r\to\infty} \frac{\ln r}{\sqrt{r}} = \lim_{r\to\infty} \frac{1/r}{1/2 \sqrt{r}} = 0. \]

So that at alteration of \( r \) within the interval from 0 to \( a, \ E_r \) would not change its sign, observance of the following requirement is necessary: \( p \leq 0. \)

At \( r = 0, \ E_r = \infty; \ f = \infty. \)

At \( r = a, \ f = 1; \ E_r = E_0. \)
2. System of equations for electron.

Basing on results of the previous section, let us write down expressions for electromagnetic field inside the electron, assuming that it is concentrated inside the orb of radius $a$.

\[ E_r = E_0 \sqrt{\frac{a}{r}} (1 + p \ln \frac{r}{a}) \sin \frac{\varphi}{2}; \quad (23') \]
\[ H_\theta = \frac{E_0}{z} \sqrt{\frac{a}{r}} (1 + p \ln \frac{r}{a}) \sin \frac{\varphi}{2}. \quad (23^*) \]

Here $a$ is electron radius, $E_0$ - amplitude of electric field intensity at $r = a$; $z = const$ - characteristic impedance inside the electron, $p$ - unknown coefficient and $p \leq 0$.

At that the internal electron medium possesses frequent and spatial dispersion, as well as anisotropy. Dispersion equations have the following appearance.

\[ v_r = v_\varphi = 2 \omega r; \quad (24) \]
\[ v_\varphi = 2 \omega r \sin \theta; \quad (24') \]
\[ z_r = z_\varphi = z_\theta = z = const. \quad (24^*) \]

Here $v_r, v_\varphi, v_\theta$ - phase velocity of rotating monochromatic electromagnetic wave in corresponding direction. In viewed case, the electromagnetic wave is being spread only in the direction $\varphi$, and we shall need expressions $v_r$ and $v_\varphi$ for searching the formulas of dielectric and magnetic permeability, as well as wave numbers of corresponding directions; $z_r, z_\varphi, z_\theta, z$ - characteristic impedances inside the electron; $\varepsilon_\varphi$ и $\mu_\varphi$ were found before, see (12'), (12^*).

\[ \varepsilon_\varphi = \frac{1}{2 \omega r z \sin \theta} = \frac{1}{v_\varphi z}; \]
\[ \mu_\varphi = \frac{z}{2 \omega r \sin \theta} = \frac{z}{v_\varphi}. \]

In view of (24),(24'),(24^*), let us write down expressions for $\varepsilon_r, \varepsilon_\varphi, \mu_r, \mu_\varphi$.

\[ \varepsilon_r = \frac{1}{2 \omega r z}; \]
\[ \mu_r = \frac{z}{2 \omega r}. \]
\[
\varepsilon_\theta = \frac{1}{2 \omega rz};
\]
\[
\mu_\theta = \frac{z}{2 \omega r}.
\]

From considerations and formulas adduced, it follows that dielectric and magnetic permeability are tensor values.

\[
\|\varepsilon\| = \begin{bmatrix}
\varepsilon_r & 0 & 0 \\
0 & \varepsilon_\varphi & 0 \\
0 & 0 & \varepsilon_\theta
\end{bmatrix} = \begin{bmatrix}
\frac{1}{2 \omega rz} & 0 & 0 \\
0 & \frac{1}{2 \omega rz \sin \theta} & 0 \\
0 & 0 & \frac{1}{2 \omega rz}
\end{bmatrix}.
\]

\[
\|\mu\| = \begin{bmatrix}
\mu_r & 0 & 0 \\
0 & \mu_\varphi & 0 \\
0 & 0 & \mu_\theta
\end{bmatrix} = \begin{bmatrix}
\frac{z}{2 \omega r} & 0 & 0 \\
0 & \frac{z}{2 \omega r \sin \theta} & 0 \\
0 & 0 & \frac{z}{2 \omega r}
\end{bmatrix}.
\]

Let us find dimensionless wave numbers.

\[
k_\varphi = \frac{\omega}{v_\varphi} r \sin \theta = \frac{\omega r \sin \theta}{2 \omega r \sin \theta} = \frac{1}{2};
\]

\[
k_\theta = \frac{\omega}{v_\theta} r = \frac{\omega r}{2 \omega r} = \frac{1}{2};
\]

\[
k_r = \frac{\omega}{v_r} r = \frac{\omega r}{2 \omega r} = \frac{1}{2}.
\]

Thus

\[
k_\varphi = k_\theta = k_r = \frac{1}{2}.
\]

Let us remind that in the viewed case, the electromagnetic wave is spread only in the direction of \(\varphi\).

At \( r = 0 \) we are having a special point:

\[
E_r = \infty; \quad H_\theta = \infty; \quad \rho = \infty; \quad j_\varphi = \infty; \quad \|\varepsilon\| = \infty; \quad \|\mu\| = \infty.
\]

Despite of this, all basic electron’s parameters - charge \( q \), rest energy \( W \), magnetic moment \( M_r \) - expressed through integrals by volume from the functions specified above, prove to be finite quantities. Look further.
From (5), we find volume charge density inside electron $\rho$.

$$\rho = \text{div}(\varepsilon E_r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left[ \frac{r^2 E_0 \sin \frac{\phi}{2}}{2 \omega rz^2} \left( \left(\frac{a}{r}\right)^2 + p \left(\frac{a}{r}\right)^2 \ln \frac{r}{a} \right) \right] =$$

$$= \frac{E_0 \sqrt{a} \sin \frac{\phi}{2}}{2 \omega rz^2} \left( \left(\frac{1}{2} + p\right) \frac{1}{\sqrt{r}} + \frac{p}{2\sqrt{r}} \ln \frac{r}{a} \right). \quad (25)$$

Integrating $\rho$ on electron’s volume, we shall receive this expression for its charge $q$.

$$q = \int \rho \, dV = \frac{E_0 \sqrt{a}}{2 \omega z} \int_0^{2\pi} \int_0^a \int_0^\pi \frac{\sin \frac{\phi}{2} r^2 \sin \theta}{r^2} \left( \left(\frac{1}{2} + p\right) \frac{1}{\sqrt{r}} + \frac{p}{2\sqrt{r}} \ln \frac{r}{a} \right) \, d\phi \, d\theta \, dr =$$

$$= \frac{4 E_0 \sqrt{a}}{\omega z} \left( \frac{1}{2} + p \right) + p \sqrt{a} \ln a - 2p \sqrt{a} - p \sqrt{a} \ln a - p \sqrt{0} \ln 0 = \frac{4 E_0 a}{\omega z}. \quad (26)$$

On the other hand, from the third integral Maxwell’s equation, it is possible to find electron’s charge as a stream of vector electric induction $D$ through the surface of the orb of radius $a$.

$$q = \oint S E_r \, dS = \frac{2\pi}{\omega z} \int_0^\pi \int_0^a \frac{E_0 \sin \frac{\phi}{2} a^2 \sin \theta}{2 \omega za} \, d\phi \, d\theta = \frac{4 E_0 a}{\omega z}. \quad (26')$$

As we can see, expressions (26) и (26’) are equivalent to each other.

From (1), we obtain expression for current density $j_v$.

$$j_v = \frac{1}{r} \frac{\partial}{\partial r} (r \frac{E_0 \sin \frac{\phi}{2}}{z} \left( \sqrt{\frac{a}{r}} + p \sqrt{\frac{a}{r}} \ln \frac{r}{a} \right)) = \frac{E_0 \sqrt{a} \sin \frac{\phi}{2}}{rz} \left( \left(\frac{1}{2} + p\right) \frac{1}{\sqrt{r}} + \frac{p}{2\sqrt{r}} \ln \frac{r}{a} \right). \quad (27)$$

From expressions (25), (27) it is visible that in the interval of change of $r$ from 0 to $a$, $\rho$ and $j_v$ once change the sign. It can be explained by the fact that in the viewed structure, the substantial role is played by the rotating monochromatic electromagnetic wave, and the space charge density and electric current density – are auxiliary or even fictitious quantities in the sense that inside the particle there is neither any charged substance nor its motion. Inside the electron, it is not the charge that is the source of electric field, but electric field is the source of the charge. In its turn, it is not the electric current that is the
source of magnetic field, but magnetic field is the source of the electric current. Thus, a deduction about vector nature of elementary charge can be made.

Now we shall determine electron’s rest energy as electromagnetic wave energy inside a particle.

\[ W = \int w dV. \]

Here \( w \) - is volume density of electromagnetic wave energy,

\[ w = \frac{\Pi}{v_{\varphi}}, \text{where} \]

\( \Pi \) – Pointing vector,

\[ \Pi = [E, H_{\theta}] \]

\( v_{\varphi} \) - phase velocity of electromagnetic wave in direction of \( \varphi \).

\[ v_{\varphi} = 2 \omega r \sin \theta. \]

\[ W = \int \int \int E_0^2 \sin^2 \frac{\varphi}{2} \frac{\theta}{2 \omega r z \sin \theta} (\frac{a}{r} + 2 p a \ln \frac{r}{a} + p^2 a \ln^2 \frac{r}{a}) r^2 \sin \theta \ d\varphi \ d\theta \ dr = \]

\[ = \frac{\pi^2 E_0^2 a^2}{2 \omega z} (a + 2 p a \ln a - 2 p \ln 0 - 2 p a - 2 p a \ln a + p^2 a (2 - 0 \ln^2 0 + 2*0 \ln 0)) = \]

\[ = \frac{\pi^2 E_0^2 a^2}{2 \omega z} (1 - 2 p + 2 p^2). \]

(28)

\[ \frac{\pi^2 E_0^2 a^2}{2 \omega z} (1 - 2 p + 2 p^2) = \hbar \omega; \]

(28')

here \( \hbar \) is Planck’s constant.

We shall be searching electron’s magnetic moment in the form of a sum.

\[ M = M_m + M_L, \]

where \( M_m \) – is magnetic moment, created by volumetric current; \( M_L \) – magnetic moment, attributed to impulse moment, i.e. to rotation.

\[ M_L = \gamma L, \]

where \( \gamma \) – gyromagnetic ratio; \( L \) – impulse moment of electron.

Basing on Barnett effect, we are making a supposition, that the impulse moment, attributed to rotation, creates additional magnetic moment.
Being aware of the fact that electron’s impulse moment is equal \( \frac{\hbar}{2} \), from (28’) we find expression for \( L \).

\[
L = \frac{\pi^2 E_0^2 a^2}{4 \omega^2 z} (2p^2 - 2p + 1);
\]

\[
M_L = \frac{\gamma \pi^2 E_0^2 a^2}{4 \omega^2 z} (2p^2 - 2p + 1);
\]

or \( M_L = \gamma \frac{\hbar}{2} \).

Let us calculate \( M_m \) as electric current magnetic moment in volume \( V \), relating to axis \( z \) by the formula:

\[
M_m = \frac{1}{2} \int \int \int_{V} r_z j_r \, dV.
\]

See for instance [3], page 111, where \( r_z \) - distance to axis \( z \),

\[
r_z = r \sin \theta.
\]

\[
M_m = \frac{1}{2} \pi \int_0^{2\pi} \int_0^a \int_0^a \frac{r \sin \theta E_0 \sqrt{a} \sin \frac{\theta}{2} \sin \theta}{z} \frac{r^2 \sin \theta}{r^2} (p + \frac{1}{2} + \frac{p'}{2} \ln \frac{r}{a}) \, d\varphi \, d\theta \, dr = \frac{\pi E_0 a^3}{z} (8p + 5), \quad \text{(29)}
\]

\[
M = M_m + M_L = \frac{\pi E_0 a^3}{z} (8p + 5) + \gamma \frac{\hbar}{2}. \quad \text{(29’)}
\]

Or

\[
M = \frac{\pi E_0 a^2}{z} \left[ a(\frac{8p + 5}{25}) + \frac{\gamma \pi E_0}{4 \omega^2} (1 - 2p + 2p^2) \right]. \quad \text{(29’’)}
\]

Thus, we have received the system of algebraic equations for electron.

\[
\begin{align*}
\frac{4 E_0 a}{\omega z} &= -e; \\
\frac{\pi E_0 a^3}{z} \left( \frac{8p + 5}{25} \right) + \frac{\gamma \hbar}{2} &= -1,0011595 \frac{e \hbar}{2m}; \\
\frac{\pi^2 E_0^2 a^2}{2 \omega z} (1 - 2p + 2p^2) &= \hbar \omega;
\end{align*}
\]

Here \( e \) - charge of electron, \( m \) - its mass.

Three equations contain five unknown quantities: \( E_0, a, z, p, \gamma \). Let us add this system with equations, which we shall receive from boundary conditions.
At \( r = a; R = a \):

\[ \varepsilon_r E_0 = \varepsilon_0 E_{\text{wena}}. \]  

(33)

In the exterior area, the same as and in the interior area, electric field intensity possesses only radial component. Here \( R \) - distance from electron’s center to the observation point in the exterior area, \( \varepsilon_0 \) - vacuum dielectric permeability.

Further. \( H_0 = \frac{E_0}{\omega} = H_{\text{wena}}. \)  

(34)

In the exterior area, the same as and in the interior area, magnetic field intensity possesses only meridional component.

It is obvious that

\[ \varepsilon_r \geq \varepsilon_0, \]  

(33’)

then from (33) follows:

\[ E_0 \leq E_{\text{wena}}. \]  

(33*)

On the other hand it is known that the electric field, having passed through dielectric layer, cannot increase, therefore

\[ E_0 \geq E_{\text{wena}}. \]  

(33")

In other words, correlations \((33'),(33*),(33")\) will be simultaneously executed only in one case, if

\[ \varepsilon_r = \varepsilon_0; \]  

(35)

\[ E_{\text{wena}} = E_0. \]  

(36)

Now under Biot-Savart’s law, we are finding magnetic field in the exterior area.

\[ B_{\text{wena}} = \frac{1}{4\pi} \int \frac{[j_{\varphi} R] \mu_0}{R^3} dV. \]

In last expression we substitute \((12*)\) and \((27)\).

\[ B_{\text{wena}} = \frac{E_0 \sqrt{a}}{4\pi R^2} \frac{2\pi a}{a} \int_0^2 \int_0^\pi \int_0^{2\pi} \sin \varphi \left[ \frac{1}{2} + p \right] \frac{1}{\sqrt{r}} + \frac{p}{2\sqrt{r}} \ln \frac{r}{a} \right] r^2 \sin \theta \ d\varphi \ d\theta \ dr = \]

\[ = \frac{E_0 \sqrt{a}}{2\omega R^2} \left[ \frac{1}{2} + p \right] 2\sqrt{a} - p\sqrt{a} \ln a + p\sqrt{a} \ln a - p\sqrt{a} \ln 0 - 2p\sqrt{a} \right] = \frac{E_0 a}{2\omega R^2}. \]  

(37)
\[ H_{\text{mean}} = \frac{B_{\text{mean}}}{\mu_0} = \frac{E_0}{2 \mu_0 \omega R^2}. \]  
(38)

At \( r = a; \ R = a \)

\[ H_{\text{app}} = \frac{E_0}{z} = H_{\text{mean}}. \]
\[ \frac{E_0}{z} = \frac{E_0}{2 \mu_0 \omega a^2} = \frac{E_0}{2 \mu_0 \omega a}; \]
\[ z = 2 \mu_0 \omega a. \]  
(39)

On the other hand, from \((24^*)\)

\[ z = 2 \mu_0 \omega r. \]

At \( r = a \)

\[ z = 2 \mu_0 \omega a. \]  
(39')

We substitute in \( (39) \).

\[ 2 \mu_0 \omega a = 2 \mu_0 \omega a; \]
\[ \mu_0 = \mu_0. \]  
(40)

Thus, at \( r = a, \)

\[ \varepsilon_r = \varepsilon_\theta = \varepsilon_0; \]
\[ \mu_0 = \mu_r = \mu_0; \]
\[ v_r = v_\theta = 2\omega a = \frac{1}{\sqrt{\mu_0 \varepsilon_0}} = c. \]  
(41)

Here \( c \) - velocity of light, \( \omega = 7,7634421 \times 10^{20} \text{ Hz} \) - Compton circular frequency of electron.

\[ a = \frac{c}{2\omega} = 0,1930796 \times 10^{-12} (m). \]  
(42)

As it is known, atom’s radius approximately equals to \( 10^{-10} \text{ m} \), volume of atom - \( 4,18879 \times 10^{-30} \text{ m}^3 \). We found, that radius of electron equals to \( 1,930796 \times 10^{-13} \text{ m} \), volume of electron –\( 3,0150724 \times 10^{-38} \text{ m}^3 \). That is one electron occupies \( 0,7197955 \times 10^{-8} \) from atom’s volume and, for example, 100 electrons (as in atoms located at the end of the periodic system) occupy \( 0,7197955 \times 10^{-5} \) from atom’s volume.

We substitute \( (42) \) в \( (39) \).

\[ z = \frac{2\mu_0 \omega c}{2\omega} = \frac{\sqrt{\mu_0}}{\varepsilon_0} = 376,73032 (\text{Ohm}). \]  
(43)
Let us solve the system (30), (31), (32), taking into account (42) and (43).

\[ \frac{4 E_0 c}{\omega \sqrt{\frac{\mu_0}{\varepsilon_0} 2\omega}} = -e; \]

\[ E_0 = -\frac{\omega^2 \mu_0 e}{2}. \]  

(30')

\[ \frac{\pi E_0}{8\omega^3 \mu_0^2 \varepsilon_0} \left( \frac{8p + 5}{25} \right) + \frac{\gamma \pi^2 E_0^2 (1 - 2p + 2p^2)}{16\omega^4 \mu_0^{3/2} \varepsilon_0^{1/2}} = -1,0011595 \frac{eh}{2m}. \]  

(31')

\[ \frac{\pi^2 E_0^2}{8\omega^3 \mu_0^{3/2} \varepsilon_0^{1/2}} (1 - 2p + 2p^2) = \hbar \omega. \]  

(32')

We substitute (30') in (32').

\[ p^2 - p + \frac{1}{2} - \frac{16 \hbar \varepsilon_0^{1/2}}{\pi^2 e^2 \mu_0^{1/2}} = 0; \]

\[ p_1 = 4,6747427; \]

\[ p_2 = -3,6747427. \]

\( p \) must be negative, therefore we select

\[ p_2 = p = -3,6747427. \]

We substitute (30') in (31').

\[ - \frac{\pi e}{16\omega \mu_0 \varepsilon_0} \left( \frac{8p + 5}{25} \right) + \frac{\pi^2 \gamma e^2 \mu_0^{1/2}}{64\varepsilon_0^{1/2}} (1 - 2p + 2p^2) = -1,0011595 \frac{eh}{2m}. \]  

(31*)

We substitute \( p \) meaning in (31*) and find \( \gamma \).

\[ \gamma = -0,2434911*10^{12} (\frac{1}{T Staple}). \]

From solution of equation (31), it is visible that two components of magnetic moment of electron \( M_m \) и \( M_L \) are directed to opposite sides and \( M_L > M_m \).

Let us also calculate numerical value of \( E_0 \) by formula (30').

\[ E_0 = -6,0673455*10^{16} (\frac{V}{m}). \]

"Dimensions" of electron for the present are not discovered by experimental way, though precision of measuring is led to \( 10^{-18} \) m. Within the framework of the model considered it may be explained by the next way: electron is not hard particle with this quantity of vector \( E \), which exist inside it, unlike from proton and neutron, quantity of vector \( E \) inside which approximately \( 10^7 \) times as much. Look below.
For positron, the system of equations will take a somewhat different view.

\[
\begin{align*}
\frac{4 E_0 a}{\omega z} &= e; \\
\pi E_0 a^3 \left( \frac{8p + 5}{25} \right) + \frac{\gamma \pi^2 E_0^2 a^2}{4\omega^2 z} (1 - 2p + 2p^2) &= 1,0011595 \frac{e \hbar}{2 m}; \\
\pi^2 E_0^2 a^2 (1 - 2p + 2p^2) &= \hbar \omega;
\end{align*}
\]

(44) \hspace{1cm} (45) \hspace{1cm} (46)

Boundary conditions are the same as for electron. Hence

\[ z = \frac{\mu_0}{\varepsilon_0}, \]
\[ a = \frac{c}{2\omega} = 0,1930796 \times 10^{-12} (m). \]

The system of equations (44), (45), (46) with exactness to a sign, has the same solutions, as the system (30), (31), (32).

\[ E_0 e^+ = - E_0 e = 6,0673455 \times 10^{16} \left( \frac{V}{m} \right); \]
\[ \gamma_{e^+} = - \gamma_e = 0,2434911 \times 10^{12} \left( \frac{1}{T* S} \right); \]
\[ p_{e^+} = - p_e = -3,6747427. \]


By applying reasoning and mathematical calculations of the previous section in relation to proton, we shall receive the relevant system of equations.

\[
\begin{align*}
\frac{4 E_0 a}{\omega z} &= e; \\
\pi E_0 a^3 \left( \frac{8p + 5}{25} \right) + \frac{\gamma \pi^2 E_0^2 a^2}{4\omega^2 z} (1 - 2p + 2p^2) &= -2,7928475 \frac{e \hbar}{2 m}; \\
\pi^2 E_0^2 a^2 (1 - 2p + 2p^2) &= \hbar \omega;
\end{align*}
\]

(47) \hspace{1cm} (48) \hspace{1cm} (49)

Here corresponding letters mean parameters of proton.

Boundary conditions: at \( r = a \)

\[ \varepsilon_r = \varepsilon_a = \varepsilon_0; \]
\[ \mu_\theta = \mu_r = \mu_0; \]
hence

\[ z = \frac{\mu_0}{\varepsilon_0}; \]

\[ a = \frac{c}{2\omega} = 1.0515447 \times 10^{-16} \text{m}. \]

Here \( \omega = 1.425486 \times 10^{24} \text{Hz} \) - Compton circular frequency of proton.

Solving the system (47), (48), (49), we shall receive:

\[ E_0 = 2.0455794 \times 10^{23} \left( \frac{V}{m} \right); \]
\[ p = -3.6747427; \]
\[ \gamma = -2.3081218 \times 10^8 \left( \frac{1}{T \cdot s} \right). \]

From the solution of equation (48) it is visible that two components of proton’s magnetic moment \( M_m \) и \( M_L \) have identical direction, and \( M_L \gg M_m \).

Let us write down the system of equations for antiproton.

\[ 4E_0 a \frac{\omega}{z} = -e; \]  
(50)

\[ \frac{\pi E_0 a^3}{z} \left( \frac{8p + 5}{25} \right) + \gamma \frac{\pi^2 E_0^2 a^2 (1 - 2p + 2p^2)}{4\omega^2 z} = 2.7928475 \frac{e \hbar}{2m}; \]  
(51)

\[ \frac{\pi^2 E_0^2 a^2}{2\omega z} (1 - 2p + 2p^2) = \hbar \omega. \]  
(52)

Boundary conditions: at \( r = a \)

\[ \varepsilon_r = \varepsilon_\theta = \varepsilon_0; \]
\[ \mu_r = \mu_\theta = \mu_0; \]

hence

\[ z = \frac{\mu_0}{\varepsilon_0}; \]

\[ a = \frac{c}{2\omega} = 1.0515447 \times 10^{-16} (m). \]

System of equations (50), (51), (52) with exactness to a sign has the same solutions, as system (47), (48), (49).
\[ E_0^\gamma = - E_0^\alpha = - 2.0455794 \times 10^{23} \left( \frac{V}{m} \right); \]
\[ \gamma_p = - \gamma_p = 2.3081218 \times 10^8 \left( \frac{1}{T \ast S} \right); \]
\[ p_p = p_p = p_e = - 3.6747427. \]


\[ E_r = E_0 \sqrt{\frac{a}{r}} (1 + p \ln \frac{r}{a}) \sin \varphi; \quad (53) \]
\[ H_\theta = \frac{E_0}{z} \sqrt{\frac{a}{r}} (1 + p \ln \frac{r}{a}) \sin \varphi. \quad (53') \]

Along each parallel, exactly one wavelength is kept within. In this case:

\[ v_\varphi = \omega r \sin \theta; \]
\[ \varepsilon_\varphi = \frac{1}{\omega r z \sin \theta}; \quad (54) \]
\[ \mu_\varphi = \frac{z}{\omega r \sin \theta}. \]
\[ v_r = 2 \omega r; \]
\[ \varepsilon_r = \frac{1}{2 \omega r z}; \quad (54') \]
\[ \mu_r = \frac{z}{2 \omega r}. \]
\[ v_\theta = v_r = 2 \omega r; \]
\[ \varepsilon_\theta = \varepsilon_r = \frac{1}{2 \omega r z}; \quad (54'') \]
\[ \mu_\theta = \mu_r = \frac{z}{2 \omega r}. \]

In other words, anisotropy is taking place, \( \varepsilon \) and \( \mu \) are tensor quantities.

\[ \| \varepsilon \| = \begin{bmatrix} \varepsilon_r & 0 & 0 \\ 0 & \varepsilon_\theta & 0 \\ 0 & 0 & \varepsilon_\theta \end{bmatrix} = \begin{bmatrix} 1 \\ \frac{1}{2 \omega r z} \\ 0 \\ \frac{1}{\omega r z \sin \theta} \\ 0 \\ \frac{1}{2 \omega r z} \end{bmatrix}. \]

\[ \| \mu \| = \begin{bmatrix} \mu_r & 0 & 0 \\ 0 & \mu_\theta & 0 \\ 0 & 0 & \mu_\theta \end{bmatrix} = \begin{bmatrix} \frac{z}{2 \omega r} \\ 0 \\ \frac{z}{\omega r \sin \theta} \\ 0 \\ \frac{z}{2 \omega r} \end{bmatrix}. \]
Here and further, corresponding letters mean parameters of neutron.

Let us find rest energy of neutron.

\[ W = \int_V w \, dV = \int_V \left[ E_r, H_\theta \right] dV = \]

\[ = \frac{2\pi}{\omega z} \frac{E_0^2 a}{r^2 \sin^2 \phi} \left( 1 + 2p \ln \frac{r}{a} + p^2 \ln^2 \frac{r}{a} \right) r^2 \sin \theta \, d\phi \, d\theta \, dr = \]

\[ = \frac{\pi^2 E_0^2 a^2}{\omega z} \left[ a + 2p a \ln a - 2p0 \ln 0 - 2p a - 2p a \ln a + p^2 a(\ln^2 1 - 2 \ln 1 + 2) - p^2 a(0 \ln^2 0 - 2*0 \ln 0) \right] = \]

\[ = \frac{\pi^2 E_0^2 a^2 (1 - 2p + 2p^2)}{\omega z}. \quad (55) \]

Further. Charge of neutron is equal to zero.

\[ q = \oint_S \varepsilon, E_r \, dS = 0. \]

Really,

\[ \int_0^{2\pi} \int_0^\omega \frac{E_0 \sin \phi}{2\omega z} a^2 \sin \theta \, d\phi \, d\theta = 0. \]

It is obvious that

\[ \int_0^{\pi} \int_0^\omega E_0 a \sin \phi \sin \theta \, d\phi \, d\theta = -\int_0^{2\pi} \int_0^\omega E_0 a \sin \phi \sin \theta \, d\phi \, d\theta \neq 0. \]

It is logical to assume that

\[ \int_0^{\pi} \int_0^\omega E_0 a \sin \phi \sin \theta \, d\phi \, d\theta = -\int_0^{2\pi} \int_0^\omega E_0 a \sin \phi \sin \theta \, d\phi \, d\theta = e. \]

Then

\[ \frac{2 E_0 a}{\omega z} = e. \quad (56) \]

Magnetic moment for neutron will be searched as the sum:

\[ M = M_m + M_L, \]

where \( M_m \) - magnetic moment created by volume current; \( M_L \) - magnetic moment, attributed to impulse moment, i.e. to rotation.

\[ M_m = \frac{1}{2} \int_0^{2\pi} \int_0^\omega \frac{E_0 \sqrt{a} r \sin \theta \sin \phi}{z r^{3/2}} (p + \frac{1}{2} + \frac{1}{2} p \ln \frac{r}{a}) r^2 \sin \theta \, d\phi \, d\theta \, dr = 0; \]
Now we shall write down the system of equations for neutron.

Boundary conditions: at $r = a$

From (54) and (57), that follows that

$$\varepsilon_r = \varepsilon_0 = \varepsilon_0;$$

$$\mu_r = H_0 = H_0;$$

$$\omega = \omega_0 = a.$$

From (54) and (57), that follows that

$$z = \left[ \frac{\mu_0}{\varepsilon_0} \right].$$

So

$$v_0 = \frac{w_0 \sin \theta}{2} = \left[ \frac{\mu_0}{\varepsilon_0} \right].$$

and from (54) and (57), that follows that

$$\mu = H_0;$$

hence

$$z = \left[ \frac{\mu_0}{\varepsilon_0} \right].$$

Here $\omega = 1.4274508 \times 10^{-24}$ Hz - Compton circular frequency of neutron.
Let us solve system (56),(55'),(57').

\[ E_0 = \frac{e \omega^2 \sqrt{\mu_0}}{c} \sqrt{\varepsilon_0}. \tag{56'} \]

\[ E_0 = E_0^n = 4,1024444 * 10^{23} \left( \frac{V}{m} \right). \]

We substitute (56') in (55').

\[ p^2 + p + 0.5 - \frac{2 \hbar}{\pi^2 e^2 \sqrt{\mu_0 / \varepsilon_0}} = 0; \]

\[ p_1 = 1,8999321; \]

\[ p_2 = -0,8999321; \]

\( p \) must be negative, therefore we select

\[ p_2 = p_n = -0,8999321. \]

From (57') we find \( \gamma \).

\[ \gamma = -1,8324711 * 10^8 \left( \frac{1}{T * s} \right). \]

Let us write down the system of equations for antineutron.

\[
\begin{aligned}
\frac{2 E_0 a}{\omega z} &= e; \\
\frac{\pi^2 E_0^2 a^2}{\omega z} (1 - 2p + 2p^2) &= \hbar \alpha; \\
\frac{\hbar}{2} &= 0,96623707 * 10^{-26}.
\end{aligned}
\]

Boundary conditions are the same, as at neutron, hence

\[ z = \sqrt{\frac{\mu_0}{\varepsilon_0}}; \]

\[ a = \frac{c}{2 \omega} = 1,0500973 * 10^{-16} (m). \]

The last system with exactness to a sign has the same solutions, as system (56),(55'),(57').

\[ E_0^n = E_0^n = 4,1024444 * 10^{23} \left( \frac{V}{m} \right); \]

\[ p_n = p_n = -0,8999321; \]

\[ \gamma_n = -\gamma_n = 1,8324711 * 10^8 \left( \frac{1}{T * s} \right). \]
Conclusion

Within the framework of the model, which is considered, electron, proton and neutron represent a monochromatic electromagnetic wave of corresponding frequency spread along parallels inside the spherical area, i.e. a wave, rotating around some axis. At that along each parallel, exactly half of wave length for electron and proton and exactly one wave length for neutron, is kept within, thus this is rotating soliton. This is caused by presence of spatial dispersion and anisotropy of a strictly defined type inside the particles. In electron vector E is directed to centre of particle, that correspond to negative charge, and in proton vector E is directed from centre of particle, that correspond to positive charge. Thus, by natural way, all basic parameters of particles are obtained: charge, rest energy, mass, radius, magnetic moment and spin, that is confirmed by mathematical expressions, which are discovered.

Literature

26.


34. I. P. Ivanov. Last days of Standard Model?

P. S. Further researches on the basis of results, which were obtained, intend solution of following tasks:


2. Elaboration of physic-mathematical model of atomic nuclei structure for all chemical elements.

    It is my firm belief that solution of this tasks will assist to achieve great leap in following fields: discovery new ways of making energy; elaboration perfectly new devices for its production; nuclear power engineering; nanotechnology, high-powerful lasers, clean energy and others.