

# Is the Universe created from Vacuum Fluctuations?

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## Abstract

It has been speculated that the universe is a zero-energy vacuum fluctuation[1]. I attempt to show how matter might be continually created in the universe from quantum fluctuations using the uncertainty and holographic principles. I derive a unique linear cosmology with a scale  $R = ct$ . I show that the annihilation time of a quantum fluctuation depends exponentially on  $R$  so that there is an effective cosmological horizon at that distance. Previously it has been found that a linear Big-Bang model[2, 3] fits the Hubble expansion data remarkably well[4, 5, 6]. Additionally such a model does not suffer from the classical horizon[7] and flatness[8] problems and so does not require a period of cosmological inflation.

We start from Heisenberg's quantum uncertainty principle

$$\Delta x \Delta p \sim \hbar \tag{1}$$

where  $\Delta x$  is the uncertainty in position and  $\Delta p$  is the uncertainty in momentum.

We assume a fluctuation with temperature  $T$  that is localised in a region of space with linear dimension of order  $R$  so that we have

$$\begin{aligned} \Delta x &\sim R \\ \Delta p &\sim \frac{k_B T}{c}. \end{aligned} \tag{2}$$

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Thus, by substituting (2) into (1), we find the following relationship for the temperature  $T$  given by

$$T \sim \frac{\hbar c}{k_B R}. \quad (3)$$

This temperature is analogous to the Hawking temperature for a black hole of radius  $R$  given by

$$T = \frac{\hbar c}{4\pi k_B R}. \quad (4)$$

Let us assume that there is some type of spherical horizon around us at radius  $R$ .

Following the holographic principle, we hypothesise that the entropy inside the horizon is given by the maximal-entropy black-hole Bekenstein-Hawking formula:

$$S = \frac{k_B A}{4\ell_P^2}. \quad (5)$$

where the area  $A = 4\pi R^2$  and the Planck length  $\ell_P = \sqrt{G\hbar/c^3}$ .

Let us assume the mass-energy inside the horizon obeys the Einstein equation

$$E = Mc^2. \quad (6)$$

Finally let us assume that the system obeys the fundamental thermodynamic relationship

$$dE = TdS \quad (7)$$

where  $E$  is the internal energy,  $T$  is the temperature and  $S$  is the entropy.

By substituting (4), (5), (6) into (7) and integrating we find that

$$R = \frac{2GM}{c^2}. \quad (8)$$

Perhaps unsurprisingly (8) is identical to the Schwarzschild black hole radius formula.

In order to investigate how the space inside the horizon expands we assume Friedmann's equation for a homogeneous and isotropic universe. Since the horizon radius  $R$  is assumed to be the scale of a region of space then we

expect it to expand with the universal scale factor  $a(t)$  such that  $R = a(t)R_0$ . Therefore we can write the Friedmann equation in terms of radius  $R$ . For simplicity we assume a spatially flat universe without an explicit cosmological constant term so that we have

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G\rho}{3}. \quad (9)$$

Now the mass  $M$  inside the spherical volume enclosed by the horizon at radius  $R$  is given by

$$M = \frac{4}{3}\pi R^3\rho. \quad (10)$$

By eliminating the mass density  $\rho$  between (9) and (10), and also the mass  $M$  using (8), we find

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{c^2}{R^2}. \quad (11)$$

Equation (11) has the linear solution

$$R = ct. \quad (12)$$

So far we have assumed that there exists some type of horizon at radius  $R$ . In order to complete this analysis we must investigate the nature of this horizon.

Let us assume that at time  $t_0$  a quantum fluctuation occurs as described by equations (2). In order for the fluctuation to annihilate, a light signal that starts at the origin at time  $t_0$  must reach the horizon at time  $t_1$ .

The comoving distance  $D_0$  travelled by such a light beam is given by

$$\begin{aligned} D_0 &= \int_{t_0}^{t_1} \frac{cdt}{a(t)} \\ &= ct_0 \int_{t_0}^{t_1} \frac{dt}{t} \\ &= R_0 \ln \frac{t_1}{t_0}. \end{aligned} \quad (13)$$

where  $a(t) = R/R_0 = t/t_0$ .

Thus the annihilation time  $t_1$  is given by

$$t_1 = t_0 \exp \frac{D_0}{R_0}. \quad (14)$$

Equation (14) shows that the annihilation time  $t_1$  increases exponentially with  $D_0$  with a length scale  $R_0$ . Thus, for all practical purposes, we do have a cosmological horizon at radius  $R_0$ .

## References

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