

Theory of Discrete Gravity

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Abstract: *Theory of everything (T.O.E), final theory or ultimate theory is a theoretical framework in the field of physics, which holds an ultimate key to unify all the fundamental forces of nature in a single field. In other words such theory can glue quantum mechanics with general relativity into a single framework. Many theories have been postulated over the decades but the dominant one includes string theory and loop quantum gravity. In this paper I would like to present a new framework which can unify quantum mechanics with general relativity by showing that the change in Riemannian metric or the bend in space time is always an integral multiple of planks constant and since gravity is the result due to bend in space-time, gravity itself is a discrete force.*

Keywords: *Discrete Gravity, General Relativity, Quantum Mechanics, Theory of Everything*

- 1. Introduction:** Physicists have been trying over decades to unify quantum mechanics with general relativity in order to define unified field theory. So far, number of theories has been evolved including String Theory and Loop Quantum Gravity. String theory predicted that the fundamental particles are made up of invisible strings, which vibrate at different frequencies and as a result of its vibration different particles are made like proton, electron etc. Whereas on the other hand Loop Quantum Gravity suggested that the whole universe has a fundamental structure or in other words a granular structure which cannot be broken down any further. But this is in complete violation of Einstein's special relativity. However both theories did have a unique approach for unifying the four fundamental forces but none of them have a complete success so as it can be globally accepted. Theory of discrete gravity is another approach to unify the theory of very big and theory of very small by postulating that the change in Riemannian metric is discrete hence the gravity is also discrete.
- 2. Planks Length:** The planks length is the scale at which the classical idea about space-time and gravity vanishes and quantum effects take over. The value of planks length is roughly 1.6×10^{-35} m and is defined by series by a series of constants i.e. Gravitational Constant (G), Planks constant (h) and speed of light (c) given by

$$l_p = \sqrt{\frac{\hbar G}{c^3}} \quad (1)$$

$$\frac{l_p^2}{\hbar} = \frac{G}{c^3} \quad (2)$$

- 3. General Relativity & Einstein Field Equation:** General Relativity is the geometric theory of gravitation published by Albert Einstein in 1915. General Relativity generalizes special relativity and Newtonian law of universal gravitation, providing a description of gravity in terms of geometric property of space-time. In other words, the curvature of space-time is related to the energy and momentum of any matter present in space-time. The relation is specified by Einstein's field equation given by

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4} T_{\mu\nu} \quad (3)$$

- 4. Discrete Gravity:** Gravity is the result of bend in space-time due to mass of an object present in it as illustrated by equation (3). In general relativity, the whole mathematics is based on minkowski space-time, which forbids space-time to have a discrete structure but what if bend in space-time is discrete rather than the space-time itself. This property of space-time will explain gravity on a quantum scale without violating any rules that general relativity is based on. This hypothesis can be proved mathematically by showing that rate of change of stress-energy-momentum tensor $T'_{\mu\nu}$ with respect to actual tensor $T_{\mu\nu}$ is always an integral multiple of planks constant or h-bar. Let us consider the mass M represented by equation (3)

Now, Let us consider a test mass m revolving around a mass M, then the force between the two masses and Energy of mass M is given by

$$F = \frac{GMm}{r^2} \quad (4)$$

$$E = Mc^2 \quad (5)$$

Also,
$$E = \int F \cdot dr \quad (6)$$

Therefore from equation (4) (5) & (6)

$$Mc^2 = \int \frac{GMm}{r^2} dr \quad (7)$$

Since, planks length is the boundary up to which classical idea exists, so the limit for integral part of equation (7) can be defined from l_p to ∞ . Therefore,

$$Mc^2 = \int_{l_p}^{\infty} \frac{GMm}{r^2} dr$$

$$\begin{aligned}
&= \frac{-GMm}{l_p} \\
&= \frac{-GMm}{l_p} \frac{4\pi r^2}{4\pi r^2} \\
4\pi l_p c^2 M &= - \left[\frac{GMm}{r^2} \right] 4\pi r^2
\end{aligned}$$

The term on the R.H.S represents area integral of force. So above equation can be represented as

$$4\pi l_p c^2 M = - \int F \cdot dA$$

Or,
$$\int F \cdot dA = - 4\pi l_p c^2 M \quad (8)$$

- 5. Divergence Theorem:** According to divergence theorem, the area integral of force is equal to the volume integral of co-variant derivative of force. Mathematically,

$$\int F \cdot dA = \int \nabla F \cdot dV \quad (9)$$

Therefore, from equation (8) and (9)

$$\int \nabla F \cdot dV = - 4\pi l_p c^2 M \quad (10)$$

The mass can be defines as the volume integral of density. So equation (10) can be re-written as

$$\begin{aligned}
\int \nabla F \cdot dv &= -4\pi l_p c^2 \int \rho \cdot dV \\
\nabla F &= -4\pi l_p c^2 \rho \quad (11)
\end{aligned}$$

- 6. Christoffel Symbol:** The christoffel symbol denoted by Γ are a form of tensor derived from Riemannian metric and is given in form of metric tensor in general relativity represented by

$$\Gamma = \frac{1}{2} g^{ad} \left[\frac{\partial g_{dc}}{\partial x^b} + \frac{\partial g_{db}}{\partial x^c} + \frac{\partial g_{bc}}{\partial x^b} \right] \quad (12)$$

For Newtonian mechanics the above equation or christoffel symbol reduced down to $\frac{1}{2} \frac{\partial g_{00}}{\partial x}$ as the value of other derivatives are very small and can be neglected. Therefore

$$\Gamma = \frac{1}{2} \frac{\partial g_{00}}{\partial x} \quad (13)$$

Equation (13) is the representation of force in Newtonian Mechanics thus the above equation can be treated equivalent to force.

$$F = \frac{1}{2} \frac{\partial g_{00}}{\partial x} \quad (14)$$

Also, $F = -\frac{d\Phi}{dt}$ where Φ is potential (15)

From equation (14) and (15)

$$\begin{aligned} \frac{1}{2} \frac{\partial g_{00}}{\partial x} &= -\frac{d\Phi}{dt} \\ \frac{1}{2} g_{00} + \text{const} &= \Phi \end{aligned} \quad (16)$$

For a three dimensional analysis instead of normal differentiation we need to represent force as covariant derivative of potential. Therefore,

$$\begin{aligned} F &= -\nabla\Phi \\ \nabla F &= -\nabla^2\Phi \end{aligned} \quad (17)$$

Substituting the value of ∇F and Φ from equation (11) and (16) respectively,

$$\begin{aligned} -\nabla^2 \frac{1}{2} g_{00} &= -4\pi l_p c^2 \rho \\ \nabla^2 g_{00} &= 8\pi l_p c^2 \rho \end{aligned} \quad (18)$$

The mass density on the R.H.S of the equation (17) can be represented as a tensor in the form of stress-energy-momentum tensor whose derivative will always be equal to zero as energy can neither be created nor destroyed. Thus,

$$\begin{aligned} 8\pi l_p c^2 \rho &= 8\pi l_p c^2 T_{\mu\nu} \quad \text{and,} \\ \nabla(8\pi l_p c^2 T_{\mu\nu}) &= 0 \end{aligned} \quad (19)$$

Multiplying and dividing by $\frac{c l_p}{\hbar^2}$ on the L.H.S of equation 19, we have

$$\nabla \left(\frac{8\pi l_p^2 c^3}{\hbar^2} \left(\frac{\hbar^2}{c l_p} T_{\mu\nu} \right) \right) = 0 \quad (20)$$

Let $\frac{\hbar^2}{c l_p} T_{\mu\nu}$ be $T'_{\mu\nu}$ So, equation 20 can be written as

$$\nabla \left(\frac{8\pi l_p^2 c^3}{\hbar^2} T'_{\mu\nu} \right) = 0 \quad (21)$$

Also, from general relativity,

$$\nabla \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) = 0 \quad (22)$$

From equation (21) and (22)

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi l_p^2 c^3}{\hbar^2} T'_{\mu\nu} \quad (23)$$

Since we are dealing with four dimensions, that is one dimension of time and three dimension of space we will divide the R.H.S of equation (23) with c^4

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi l_p^2 c^3}{c^4 \hbar^2} T'_{\mu\nu} \quad (24)$$

And, from equation (3) and equation (24)

$$\frac{8\pi l_p^2 c^3}{c^4 \hbar^2} T'_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$\frac{8\pi l_p^2 c^3}{c^4 \hbar} T'_{\mu\nu} = \hbar \frac{8\pi G}{c^4} T_{\mu\nu}$$

Substituting the value of $\frac{l_p^2}{\hbar}$ from equation (2) in above equation

$$\frac{8\pi G}{c^4} T'_{\mu\nu} = \hbar \frac{8\pi G}{c^4} T_{\mu\nu} \quad (25)$$

The tensors $T'_{\mu\nu}$ and $T_{\mu\nu}$ represent different value of Riemannian metric or in General Einstein's metric. Thus equation (25) can be expressed as

$$G'_{\mu\nu} = \hbar G_{\mu\nu}$$

Or,
$$\frac{G'_{\mu\nu}}{G_{\mu\nu}} = \hbar$$

In general,
$$\frac{G'_{\mu\nu}}{G_{\mu\nu}} \geq n\hbar \quad [\text{where } n=1,2,3,\dots] \quad (26)$$

Where $G'_{\mu\nu}$ is the change in Riemannian metric and $G_{\mu\nu}$ is the initial value of Riemannian Metric.

7. Conclusion: Due to the change in the mass of a body, the Riemannian metric changes accordingly. When there is a change in mass by a factor of $\frac{\hbar^2}{cl_p}$ or by Mihir factor (denoted

by M) the Riemannian metric changes discretely. In general, the change in mass by an integral multiple of M leads to change in space-time discretely. This can be clearly observed from equation (26). And since, gravity is the result of bend in Space-time and the bend of space-time is discrete in nature, this must conclude that the gravitational force itself is discrete in nature.

8. References:

- [1] R. P. Feynman, *Space-time approach to non-relativistic quantum mechanics*, *Rev. Mod. Phys.* 20, 367-387 (1948)
- [2] Arnowitt R, Deser S and Misner C W 1962 *The dynamics of general relativity*, in *Gravitation: An introduction to current research* ed Witten L (John Wiley, New York)
- [3] Wheeler J A 1962 *Geometrodynamics*, (Academic Press, New York)
- [4] Wheeler J A 1964 *Geometrodynamics and the issue of the final state* *Relativity, Groups and Topology* eds DeWitt C M and DeWitt B S (Gordon and Breach, New York)
- [5] Komar A 1970 *Quantization program for general relativity*, in *Relativity* Carmeli M, Fickler S. I. and Witten L (eds) (Plenum, New York)
- [6] Ashtekar A and Geroch R 1974 *Quantum theory of gravitation*, *Rep. Prog. Phys.* 37 1211-1256
- [7] Weinberg S 1972 *Gravitation and Cosmology* (John Wiley, New York)
- [8] DeWitt B S 1972 *Covariant quantum geometrodynamics*, in *Magic Without Magic: John Archibald Wheeler* ed Klauder J R (W. H. Freeman, San Francisco)
- [9] Isham C. J. 1975 *An introduction to quantum gravity*, in *Quantum Gravity*, An Oxford Symposium Isham C J, Penrose R and Sciama D W (Clarendon Press, Oxford)
- [10] Duff M 1975 *Covariant quantization in Quantum Gravity*, An Oxford Symposium Isham C J, Penrose R and Sciama D W (Clarendon Press, Oxford)
- [11] Penrose R 1975 *Twistor theory, its aims and achievements* *Quantum Gravity*, An Oxford Symposium Isham C J, Penrose R and Sciama D W (Clarendon Press, Oxford)
- [12] Israel W and Hawking S W eds 1980 *General Relativity*, An Einstein Centenary Survey (Cambridge UP, Cambridge)
- [13] Bergmann P G and Komar A 1980 *The phase space formulation of general relativity and approaches toward its canonical quantization* *General Relativity and Gravitation* vol 1, On Hundred Years after the Birth of Albert Einstein, Held A ed (Plenum, New York)
- [14] Wolf H (ed) 1980 *Some Strangeness in Proportion* (Addison Wesley, Reading)
- [15] Hawking S W 1980 *Is End In Sight for Theoretical Physics?: An Inaugural Address* (Cambridge UP, Cambridge)
- [16] Kuchař K 1981 *Canonical methods of quantization*, in *Quantum Gravity 2*, A Second Oxford Symposium Isham C J, Penrose R and Sciama D W (Clarendon Press, Oxford)
- [17] Isham C J 1981 *Quantum gravity—An overview*, in *Quantum Gravity 2*, A Second Oxford Symposium Isham C J, Penrose R and Sciama D W (Clarendon Press, Oxford)
- [18] Ko M, Ludvigsen M, Newman E T and Tod P 1981 *The theory of H space* *Phys. Rep.* 71 51–139
- [19] Ashtekar A 1984 *Asymptotic quantization* (Bibliopolis, Naples); also available at <http://cgpg.gravity.psu.edu/research/asymquant-book.pdf>
- [20] Greene M B, Schwarz J H and Witten E 1987 *Superstring theory*, volumes 1 and 2 (Cambridge UP, Cambridge)
- [21] R. Penrose and W. Rindler, *Spinors and space-times*, Vol 2, (Cambridge University Press, Cambridge 1988)
- [22] Ashtekar A 1991 *Lectures on non-perturbative canonical gravity*, Notes prepared in collaboration with R. S. Tate (World Scientific, Singapore)
- [23] Ashtekar A *Mathematical problems of non-perturbative quantum general relativity in Gravitation and Quantizations: Proceedings of the 1992 Les Houches summer school* eds Julia B and Zinn-Justin J (Elsevier, Amsterdam); also available as *gr-qc/9302024*
- [24] Williams R W and Tucker P M 1992 *Regge calculus: a brief review and bibliography*, *Class. Quant. Grav.* 9 1409-1422
- [25] Connes A 1994 *Non-commutative geometry*, (Academic Press, New York)