Magnetohydrodynamic Equations Solutions

Abstract

The simplest solution is usually the best solution---Albert Einstein

The system of magnetohydrodynamic (MHD) equations have been solved analytically in this paper. The author applied the technique used in solving the Navier-Stokes equations and applied a new law, the law of definite ratio for MHD. This law states that in MHD, the other terms of the system of equations divide the gravity term in a definite ratio, and each term utilizes gravity to function. The sum of the terms of the ratio is always unity. It is shown that without gravity forces on earth, there would be no magnetohydrodynamics on earth as is known. The equations in the system of equations were added to produce a single equation which was then integrated. Ratios were used to split-up this single equation into sub-equations which were readily integrated, and even, the non-linear sub-equations were readily integrated. Twenty-seven sub-equations were integrated. The linear part of the relation obtained from the integration of the linear part of the equation satisfied the linear part of the equation; and the relation from the integration of the non-linear part satisfied the non-linear part of the equation. The solutions revealed the role of each term in magnetohydrodynamics. In particular, the gravity term is the indispensable term in magnetohydrodynamics. The solutions of the MHD equations were compared with the solutions of the N-S equations, and there were similarities and dissimilarities.
Solutions of the Magnetohydrodynamic Equations

This system consists of four equations and one is to solve for \( V_x, V_y, V_z, B_x, B_y, B_z, P(x) \)

\[
\begin{align*}
\text{Magnetohydrodynamic Equations} \\
1. & \quad \frac{\partial V_x}{\partial t} + V_x \frac{\partial V_x}{\partial x} + V_y \frac{\partial V_x}{\partial y} + V_z \frac{\partial V_x}{\partial z} = -\frac{\partial p}{\partial x} + \frac{1}{\mu} (\nabla \times B) \times B + \rho g_x \\
2. & \quad \frac{\partial V_y}{\partial t} + V_x \frac{\partial V_y}{\partial x} + V_y \frac{\partial V_y}{\partial y} + V_z \frac{\partial V_y}{\partial z} = -\frac{\partial p}{\partial y} + \frac{1}{\mu} (\nabla \times B) \times B + \rho g_x \\
3. & \quad \frac{\partial V_z}{\partial t} + V_x \frac{\partial V_z}{\partial x} + V_y \frac{\partial V_z}{\partial y} + V_z \frac{\partial V_z}{\partial z} = -\frac{\partial p}{\partial z} + \frac{1}{\mu} (\nabla \times B) \times B + \rho g_x \\
4. & \quad \nabla \cdot B = 0
\end{align*}
\]

Step 1:
1. If \( \rho \) is constant (for incompressible fluid)
   \[
   \frac{\partial V_x}{\partial t} + V_x \frac{\partial V_x}{\partial x} + V_y \frac{\partial V_x}{\partial y} + V_z \frac{\partial V_x}{\partial z} = 0 \quad \text{<-- continuity equation}
   \]

2. \[\rho \frac{\partial V_x}{\partial t} + \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} = -\frac{\partial p}{\partial x} + \frac{1}{\mu} (\nabla \times B) \times B + \rho g_x\]

3. \[\rho \frac{\partial B}{\partial t} = \nabla \times (V \times B) + \eta \nabla^2 B\]

4. \[\nabla \cdot B = 0\]
Magnetohydrodynamic Equations

Step 2:

After the "vector juggling" one obtains the following system of equations which one will solve.

\[
\begin{align*}
1. \quad & \frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z} = 0 \\
2. \quad & \rho \frac{\partial V_x}{\partial t} + \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} + \frac{\partial p}{\partial x} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} + \frac{1}{\mu} B_z \frac{\partial B_y}{\partial x} + \frac{1}{\mu} B_y \frac{\partial B_x}{\partial y} = \rho g_x \\
3. \quad & \frac{\partial B_x}{\partial t} - V_x \frac{\partial B_y}{\partial y} - V_y \frac{\partial B_x}{\partial x} + B_x \frac{\partial V_y}{\partial y} + B_y \frac{\partial V_x}{\partial x} + \eta \frac{\partial^2 B_x}{\partial x^2} - \frac{\eta \partial^2 B_x}{\partial y^2} - \frac{\eta \partial^2 B_x}{\partial z^2} = 0 \\
4. \quad & \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0
\end{align*}
\]

At a glance, and from the experience gained in solving the Navier-Stokes equations, one can identify equation (2) as the driver equation, since it contains the gravity term, and the gravity term is the subject of the equation. However, since the system of equations is to be solved simultaneously and there is only a single "driver", the gravity term, all the terms in the system of equations will be placed in the driver equation, Equation 2. As suggested by Albert Einstein, Friedrich Nietzsche, and Pablo Picasso, one will think like a child at the next step.

Step 3:

Thinking like a ninth grader, one will apply the following axiom:

\[
\text{If } a = b \text{ and } c = d, \text{ then } a + c = b + d; \text{ and therefore, add the left sides and add the right sides of the above equations. That is, } (1) + (2) + (3) + (4) = \rho g_x
\]

\[
\begin{align*}
\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z} + \rho \frac{\partial V_x}{\partial t} + \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} + \frac{\partial p}{\partial x} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} + \frac{1}{\mu} B_z \frac{\partial B_y}{\partial x} + \frac{1}{\mu} B_y \frac{\partial B_x}{\partial y} &= \rho g_x \\
\frac{1}{\mu} B_y \frac{\partial B_x}{\partial y} - V_x \frac{\partial B_y}{\partial y} - V_y \frac{\partial B_x}{\partial x} + B_x \frac{\partial V_y}{\partial y} + B_y \frac{\partial V_x}{\partial x} + \frac{\eta \partial^2 B_x}{\partial x^2} - \frac{\eta \partial^2 B_x}{\partial y^2} - \frac{\eta \partial^2 B_x}{\partial z^2} &= 0
\end{align*}
\]

(Three lines per equation)

Step 4:

Writing all the linear terms first

\[
\begin{align*}
\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z} + \rho \frac{\partial V_x}{\partial t} + \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} + \frac{\partial p}{\partial x} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} + \frac{1}{\mu} B_z \frac{\partial B_y}{\partial x} + \frac{1}{\mu} B_y \frac{\partial B_x}{\partial y} + \frac{\eta \partial^2 B_x}{\partial x^2} - \frac{\eta \partial^2 B_x}{\partial y^2} - \frac{\eta \partial^2 B_x}{\partial z^2} &= \rho g_x \\
+ \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} + \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_x}{\partial y} + \rho V_z \frac{\partial V_x}{\partial z} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} - B_x \frac{\partial B_x}{\partial x} - B_y \frac{\partial B_x}{\partial y} - B_z \frac{\partial B_x}{\partial z} &= \rho g_x
\end{align*}
\]

(Three lines per equation)

(Since all the terms are now in the same driver equation, let \( \rho g_x \) "drive them" simultaneously.)

Step 5:

Solve the above 28-term equation using the ratio method. (27 ratio terms)

The ratio terms to be used are respectively the following: (Sum of the ratio terms = 1)

\[
\begin{align*}
\beta_1, \beta_2, \beta_3, a, b, c, d, f, m, q, r, s, w_1, w_2, w_3, w_4, w_5, w_6, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5, \lambda_6, \lambda_7, \lambda_8, \lambda_9
\end{align*}
\]

\[
\begin{align*}
1. \quad & \frac{\partial V_x}{\partial x} = \beta_1 \rho g_x \\
2. \quad & \frac{dV_x}{dx} = \beta_1 \rho g_x \\
V_x &= \beta_1 \rho g_x x + C_{16} \\
V_x &= \beta_1 \rho g_x y + C_{17} \\
V_x &= \beta_1 \rho g_x z + C_{18}
\end{align*}
\]

(Three lines per equation)
### Magnetohydrodynamic Equations

#### 5.
\[
\begin{align*}
\frac{\partial p}{\partial x} &= b \rho g_x \\
\frac{dp}{dx} &= b \rho g_x \\
P(x) &= b \rho g_x x + C
\end{align*}
\]

#### 6.
\[
\begin{align*}
\rho \frac{\partial B_x}{\partial t} &= c \rho g_x \\
\frac{\partial B_x}{\partial t} &= c g_x \\
\frac{dB_z}{dt} &= c g_x \\
B_x &= c g_x t + C_1 b
\end{align*}
\]

#### 7.
\[
\begin{align*}
-\eta \frac{\partial^2 B_x}{\partial x^2} &= d \rho g_x \\
\frac{d^2 B_x}{dx^2} &= - \frac{d \rho g_x}{\eta} \\
\frac{dB_x}{dx} &= - \frac{d \rho g_x}{\eta} + C_2 \\
B_x &= - \frac{d \rho g_x x^2}{2 \eta} + C_3 x + C_3
\end{align*}
\]

#### 8.
\[
\begin{align*}
-\eta \frac{\partial^2 B_z}{\partial y^2} &= f \rho g_x \\
\frac{d^2 B_z}{dy^2} &= - \frac{f \rho g_x}{\eta} \\
\frac{dB_z}{dy} &= - \frac{f \rho g_x y}{\eta} + C_4 \\
B_z &= - \frac{f \rho g_x y^2}{2 \eta} + C_4 y + C_5
\end{align*}
\]

#### 9.
\[
\begin{align*}
-\eta \frac{\partial^2 B_z}{\partial z^2} &= m \rho g_x \\
\frac{d^2 B_z}{dz^2} &= - \frac{m \rho g_x}{\eta} \\
\frac{dB_z}{dz} &= - \frac{m \rho g_x z}{\eta} + C_6 \\
B_z &= - \frac{m \rho g_x z^2}{2 \eta} + C_6 x + C_7
\end{align*}
\]

#### 10.
\[
\begin{align*}
\frac{\partial B_x}{\partial x} &= q \rho g_x \\
\frac{dB_x}{dx} &= q \rho g_x \\
B_x &= q \rho g_x x + C_{19}
\end{align*}
\]

#### 11.
\[
\begin{align*}
\frac{\partial B_y}{\partial y} &= r \rho g_x \\
\frac{dB_y}{dy} &= r \rho g_x \\
B_y &= r \rho g_x y + C_{20}
\end{align*}
\]

#### 12.
\[
\begin{align*}
\frac{\partial B_z}{\partial z} &= s \rho g_x \\
\frac{dB_z}{dz} &= s \rho g_x \\
B_z &= s \rho g_x z + C_{21}
\end{align*}
\]

#### 13.
\[
\begin{align*}
\rho V_x \frac{\partial V_x}{\partial x} &= \omega_1 \rho g_x \\
V_x \frac{dV_x}{dx} &= \omega_1 g_x \\
V_x \frac{dV_x}{dx} &= \omega_1 g_x dx \\
V_x &= \frac{1}{2} \omega_1 g_x x \\
V_x &= 2 \omega_1 g_x x \\
V_x &= \pm \sqrt{\frac{1}{2} \omega_1 g_x x} + C_2
\end{align*}
\]

#### 14.
\[
\begin{align*}
\rho V_y \frac{\partial V_y}{\partial y} &= \omega_2 \rho g_x \\
V_y \frac{dV_y}{dy} &= \omega_2 g_x dy \\
V_y \frac{dV_y}{dy} &= \omega_2 g_x y + \psi_y(V_y) \\
V_y &= \frac{\omega_2 g_x y}{V_y} + \frac{\psi_y(V_y)}{V_y} \\
V_y &\neq 0
\end{align*}
\]

#### 15.
\[
\begin{align*}
\rho V_z \frac{\partial V_z}{\partial z} &= \omega_3 \rho g_x \\
V_z \frac{dV_z}{dz} &= \omega_3 g_x \\
V_z \frac{dV_z}{dz} &= \omega_3 g_x dz \\
V_z \frac{dV_z}{dz} &= \omega_3 g_x z + \psi_z(V_z) \\
V_z &= \omega_3 g_x z + \psi_z(V_z) \\
V_z &\neq 0
\end{align*}
\]

#### 16.
\[
\begin{align*}
\frac{\partial B_z}{\partial z} &= - \omega_4 \mu g_x \\
\frac{dB_z}{dz} &= - \omega_4 \mu g_x dz \\
B_z &= - \omega_4 \mu g_x z + \psi_z(B_z) \\
B_z &\neq 0
\end{align*}
\]

#### 17.
\[
\begin{align*}
\frac{\partial B_z}{\partial z} &= \omega_5 \mu g_x \\
\frac{dB_z}{dz} &= \omega_5 \mu g_x dz \\
B_z &= \omega_5 \mu g_x dx \\
B_z &= \sqrt{2} \omega_5 \mu g_x x \\
B_z &= \pm \sqrt{2} \omega_5 \mu g_x x + C
\end{align*}
\]
### Magnetohydrodynamic Equations

#### 18.

\[
\begin{align*}
B_y \frac{\partial B_y}{\partial x} &= \omega_6 \mu_0 g_x \\
B_y \frac{\partial B_y}{\partial y} &= \omega_6 \mu_0 g_x \\
B_y \frac{\partial B_y}{\partial z} &= \omega_6 \mu_0 g_x \\
B_y \frac{\partial B_y}{\partial t} &= \omega_6 \mu_0 g_x \\
B_y^2 &= \omega_6 \mu_0 g_x x \\
\frac{D B_y}{D t} &= 2\omega_6 \mu_0 g_x x + C \\
B_y &= \pm \sqrt{2\omega_6 \mu_0 g_x x + C} \\
\end{align*}
\]

#### 19.

\[
\begin{align*}
-\frac{1}{\mu} \frac{\partial B_y}{\partial y} &= \lambda_1 \rho g_x \\
B_y \frac{\partial B_y}{\partial y} &= -\lambda_1 \mu g_x \\
B_y \frac{\partial B_y}{\partial y} &= -\lambda_1 \mu g_x \\
B_y B_y &= -\lambda_1 \mu g_x + \psi_y(B_y) \\
B_y &= -\frac{\lambda_1 \mu g_x y + \psi_y(B_y)}{B_y} \\
B_y &= 0 \\
\end{align*}
\]

#### 20.

\[
\begin{align*}
-\frac{V_y}{\mu} \frac{\partial B_y}{\partial y} &= \lambda_2 \rho g_x \\
V_y \frac{\partial B_y}{\partial y} &= -\lambda_2 \rho g_x \\
V_y \frac{\partial B_y}{\partial y} &= -\lambda_2 \rho g_x \\
V_y B_y &= -\lambda_2 \rho g_x y + \psi_y(V_y) \\
B_y &= -\frac{\lambda_2 \rho g_x y + \psi_y(V_y)}{V_y} \\
V_y &= 0 \\
\end{align*}
\]

#### 21.

\[
\begin{align*}
-\frac{B_y}{\mu} \frac{\partial V_y}{\partial y} &= \lambda_3 \rho g_x \\
B_y \frac{\partial V_y}{\partial y} &= -\lambda_3 \rho g_x \\
B_y \frac{\partial V_y}{\partial y} &= -\lambda_3 \rho g_x \\
B_y V_y &= -\lambda_3 \rho g_x y + \psi_y(B_y) \\
V_y &= -\frac{\lambda_3 \rho g_x y + \psi_y(B_y)}{B_y} \\
B_y &= 0 \\
\end{align*}
\]

#### 22.

\[
\begin{align*}
-\frac{V_y}{\mu} \frac{\partial B_y}{\partial y} &= \lambda_4 \rho g_x \\
V_y \frac{\partial B_y}{\partial y} &= \lambda_4 \rho g_x \\
V_y \frac{\partial B_y}{\partial y} &= \lambda_4 \rho g_x \\
V_y B_y &= \lambda_4 \rho g_x y + \psi_y(V_y) \\
B_y &= \frac{\lambda_4 \rho g_x y + \psi_y(V_y)}{V_y} \\
V_y &= 0 \\
\end{align*}
\]

#### 23.

\[
\begin{align*}
-\frac{V_x}{\mu} \frac{\partial B_y}{\partial y} &= \lambda_5 \rho g_x \\
V_x \frac{\partial B_y}{\partial y} &= \lambda_5 \rho g_x \\
V_x \frac{\partial B_y}{\partial y} &= \lambda_5 \rho g_x \\
V_x B_y &= \lambda_5 \rho g_x y + \psi_y(V_y) \\
B_y &= \frac{\lambda_5 \rho g_x y + \psi_y(V_y)}{V_y} \\
V_y &= 0 \\
\end{align*}
\]

#### 24.

\[
\begin{align*}
V_z \frac{\partial B_x}{\partial z} &= \lambda_6 \rho g_x \\
V_z \frac{\partial B_x}{\partial z} &= \lambda_6 \rho g_x \\
V_z \frac{\partial B_x}{\partial z} &= \lambda_6 \rho g_x \\
V_z B_z &= \lambda_6 \rho g_s z + \psi_z(V_z) \\
B_z &= \frac{\lambda_6 \rho g_s z + \psi_z(V_z)}{V_z} \\
V_z &= 0 \\
\end{align*}
\]

#### 25.

\[
\begin{align*}
B_x \frac{\partial V_z}{\partial z} &= \lambda_7 \rho g_x \\
B_x \frac{\partial V_z}{\partial z} &= \lambda_7 \rho g_x \\
B_x \frac{\partial V_z}{\partial z} &= \lambda_7 \rho g_x \\
V_z B_z &= \lambda_7 \rho g_s z + \psi_z(B_z) \\
V_z &= \frac{\lambda_7 \rho g_s z + \psi_z(B_z)}{B_z} \\
B_z &= 0 \\
\end{align*}
\]

#### 26.

\[
\begin{align*}
-\frac{V_x}{\mu} \frac{\partial B_z}{\partial z} &= \lambda_8 \rho g_x \\
V_x \frac{\partial B_z}{\partial z} &= -\lambda_8 \rho g_x \\
V_x \frac{\partial B_z}{\partial z} &= -\lambda_8 \rho g_x \\
V_x B_z &= -\lambda_8 \rho g_s z + \psi_z(V_z) \\
B_z &= \frac{\lambda_8 \rho g_s z + \psi_z(V_z)}{V_z} \\
V_z &= 0 \\
\end{align*}
\]

#### 27.

\[
\begin{align*}
-\frac{B_z}{\mu} \frac{\partial V_x}{\partial z} &= \lambda_9 \rho g_x \\
B_z \frac{\partial V_x}{\partial z} &= -\lambda_9 \rho g_x \\
B_z \frac{\partial V_x}{\partial z} &= -\lambda_9 \rho g_x \\
B_z V_x &= -\lambda_9 \rho g_s z + \psi_z(B_z) \\
V_x &= \frac{-\lambda_9 \rho g_s z + \psi_z(B_z)}{B_z} \\
B_z &= 0 \\
\end{align*}
\]
Step 6: One collects the integrals of the sub-equations, above, for $V_x, V_y, V_z, B_x, B_y, B_z, P(x)$:

$$
V_x(x,y,z,t) = \beta_1 \rho g_{x,x} + a g_{x,t} \pm \sqrt{2 \omega_1 g_{x,x} + \frac{\omega_2 g_{y,y}}{V_y} - \frac{\lambda_3 \rho g_{x,y}}{B_y} + \frac{\omega_3 g_{x,z}}{V_z} - \frac{\lambda_9 \rho g_{x,z}}{B_z} + \psi_x(V_y) + \psi_y(B_x) + \psi_y(V_y) + \psi_z(B_z) + C_1;
$$

$$
P(x) = b \rho g_{x,x} + C_2;
$$

$$
V_y(y) = \beta_2 \rho g_{x,y} + \frac{\lambda_5 \rho g_{x,y}}{B_y} + \frac{\psi_x(B_x)}{B_y} + C_3;
$$

$$
V_z(z) = \beta_3 \rho g_{x,z} + \frac{\lambda_7 \rho g_{x,z}}{B_x} + \frac{\psi_x(B_x)}{B_x} + C_4;
$$

$$
B_x(x,y,z,t) = -\frac{\rho g_{x,x}}{2 \eta} (dx^2 + dy^2 + dz^2) + q \rho g_{x,x} + C_2 x + C_4 y + C_6 z + c g_{x,t} t - \frac{\lambda_4 \rho g_{y,y}}{V_y} + \frac{\psi_y(B_y)}{B_y} + \frac{\psi_y(V_y)}{V_y} + \frac{\psi_z(B_z)}{V_z} + C_7;
$$

$$
B_y(y) = r \rho g_{x,y} \pm \sqrt{2 \omega_6 \mu g_{x,x} - \frac{\lambda_2 \rho g_{y,y}}{V_x} + \frac{\psi_x(V_x)}{V_x} + C_8}
$$

$$
B_z(z) = s \rho g_{x,z} \pm \sqrt{2 \omega_5 \mu g_{x,x} - \frac{\lambda_8 \rho g_{x,z}}{V_x} + \frac{\psi_x(V_x)}{V_x} + C_9}.
$$
Step 7: Find the test derivatives for the linear part

1. \( \frac{\partial V_x}{\partial x} = (\beta_1 \rho g_x) \)
2. \( \frac{\partial V_y}{\partial y} = (\beta_2 \rho g_x) \)
3. \( \frac{\partial V_z}{\partial z} = (\beta_3 \rho g_x) \)
4. \( \frac{\partial \rho}{\partial t} = (a g_x) \)
5. \( \frac{\partial p}{\partial x} = (b \rho g_x) \)
6. \( \frac{dB_z}{dt} = (c g_x) \)

7. \( \frac{\partial^2 B_x}{\partial x^2} = -\frac{dp g_x}{\eta} \)
8. \( \frac{\partial^2 B_x}{\partial y^2} = -\frac{fp g_x}{\eta} \)
9. \( \frac{\partial^2 B_x}{\partial z^2} = -\frac{m g_x}{\eta} \)
10. \( \frac{\partial B_x}{\partial y} = q g_x \)
11. \( \frac{\partial B_x}{\partial z} = r g_x \)
12. \( \frac{\partial B_z}{\partial z} = s g_x \)

Test derivatives for the nonlinear part

13. \( \frac{\partial V_x}{\partial x} = \frac{\omega g_x}{V_x} \)
14. \( \frac{\partial V_y}{\partial y} = \frac{\omega g_x}{V_y} \)
15. \( \frac{\partial V_z}{\partial z} = \frac{\omega g_x}{V_z} \)
16. \( \frac{\partial B_x}{\partial y} = -\frac{\lambda_1 \mu g_x}{B_y} \)
17. \( \frac{\partial B_x}{\partial z} = -\frac{\lambda_2 \mu g_x}{B_z} \)
18. \( \frac{\partial B_y}{\partial z} = \frac{\omega_6 \mu g_x}{B_y} \)
19. \( \frac{\partial B_z}{\partial y} = -\frac{\lambda_1 \mu g_x}{B_z} \)
20. \( \frac{\partial B_z}{\partial z} = -\frac{\lambda_2 g_x}{B_z} \)
21. \( \frac{\partial V_x}{\partial y} = -\frac{\lambda_3 \rho g_x}{V_y} \)
22. \( \frac{\partial V_y}{\partial z} = -\frac{\lambda_4 \rho g_x}{V_y} \)
23. \( \frac{\partial V_z}{\partial x} = \frac{\lambda_5 \rho g_x}{B_x} \)
24. \( \frac{\partial B_y}{\partial x} = \frac{\lambda_6 \rho g_x}{B_x} \)
25. \( \frac{\lambda_7 \rho g_x}{B_x} \)
26. \( \frac{\partial B_z}{\partial x} = \frac{\lambda_8 \rho g_x}{B_z} \)
27. \( \frac{\partial V_y}{\partial x} = \frac{\lambda_9 \rho g_x}{B_x} \)

Step 8: Substitute the above test derivatives respectively in the following 28-term equation

\[
\begin{align*}
\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z} + \rho \frac{\partial V_x}{\partial t} + \frac{\partial \rho}{\partial x} &+ \eta \frac{\partial^2 B_x}{\partial x^2} - \eta \frac{\partial^2 B_x}{\partial y^2} - \eta \frac{\partial^2 B_x}{\partial z^2} + \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} \\
+ \rho V_x \frac{\partial V_x}{\partial x} + \rho V_y \frac{\partial V_y}{\partial y} + \rho V_z \frac{\partial V_z}{\partial z} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial z} + \frac{1}{\mu} B_z \frac{\partial B_y}{\partial z} + \frac{1}{\mu} B_y \frac{\partial B_x}{\partial y} - \frac{1}{\mu} B_z \frac{\partial B_x}{\partial y} - V_x \frac{\partial B_x}{\partial y} - B_y \frac{\partial V_x}{\partial y} \\
+ V_y \frac{\partial B_x}{\partial y} + B_x \frac{\partial V_y}{\partial y} + V_z \frac{\partial B_x}{\partial z} - V_x \frac{\partial B_x}{\partial z} - B_z \frac{\partial V_z}{\partial z} = \rho g_x
\end{align*}
\]

(Three lines per equation)

\[
\begin{align*}
(\beta_1 \rho g_x) + (\beta_2 \rho g_x) + (\beta_3 \rho g_x) + (\rho a g_x) + (b \rho g_x) + (\rho c g_x) - \eta (-\frac{dp g_x}{\eta}) - \eta (-\frac{fp g_x}{\eta}) - \eta (-\frac{mg g_x}{\eta}) \\
(q \rho g_x) + (r \rho g_x) + (s \rho g_x) + \rho V_x \frac{\omega g_x}{V_x} + \rho V_y \frac{\omega g_x}{V_y} + \rho V_z \frac{\omega g_x}{V_z} - \frac{1}{\mu} B_z \frac{\omega g_x}{B_z} \\
\frac{1}{\mu} B_z \frac{\omega g_x}{B_z} + \frac{1}{\mu} B_y \frac{\omega g_x}{B_y} \frac{1}{\mu} B_y \frac{\lambda_1 \mu g_x}{B_y} - V_x \frac{\lambda_2 \rho g_x}{V_x} - B_y \frac{\lambda_3 \rho g_x}{B_y} + V_y \frac{\lambda_4 \rho g_x}{V_y} \\
B_x \frac{\lambda_5 \rho g_x}{B_x} + V_z \frac{\lambda_6 \rho g_x}{V_z} + B_x \frac{\lambda_7 \rho g_x}{B_x} - V_x \frac{\lambda_8 \rho g_x}{V_x} - B_z \frac{\lambda_9 \rho g_x}{B_z} = \rho g_x
\end{align*}
\]

(Four lines per equation)

\[
\begin{align*}
\beta_1 \rho g_x + \beta_2 \rho g_x + \beta_3 \rho g_x + a \rho g_x + b \rho g_x + c \rho g_x + d \rho g_x + f \rho g_x + m \rho g_x + q \rho g_x + r \rho g_x + s \rho g_x + \omega_1 \rho g_x \\
+ \omega_3 \rho g_x + \omega_5 \rho g_x + \omega_6 \rho g_x + \lambda_1 \mu g_x + \lambda_2 \rho g_x + \lambda_3 \rho g_x + \lambda_4 \rho g_x + \lambda_5 \rho g_x + \omega_2 \rho g_x + \omega_3 \rho g_x \\
+ \lambda_6 \rho g_x + \lambda_7 \rho g_x + \lambda_8 \rho g_x + \lambda_9 \rho g_x = \rho g_x
\end{align*}
\]

(Three lines per equation)
Magnetohydrodynamic Equations

\[ \begin{align*}
\beta_1 g_x + \beta_2 g_x + \beta_3 g_x + a g_x + b g_x + c g_x + d g_x + f g_x + m g_x q g_x + r g_x + s g_x + \omega_1 g_x + \omega_3 g_x + \omega_5 g_x \\
+ \omega_6 g_x + \lambda_1 g_x + \lambda_2 g_x + \lambda_3 g_x + \lambda_4 g_x + \lambda_5 g_x + \omega_2 g_x + \omega_3 g_x + \lambda_6 g_x + \lambda_7 g_x + \lambda_8 g_x + \lambda_9 g_x = g_x
\end{align*} \]  
(2 lines)

\[ \begin{align*}
g_x (1) = g_x \\
(\text{Sum of the ratio terms} = 1)
\end{align*} \]

Since an identity is obtained, the solutions to the 28-term equation are as follows

<table>
<thead>
<tr>
<th>Equation</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>( V_x(x, y, z, t) = )</td>
<td>(sum of integrals from sub-equations #1, #4, #13, #14, #15, #21, #27)</td>
</tr>
<tr>
<td>( P(x) = b \rho g_x x + C_2 )</td>
<td>(integral from sub-equation #5)</td>
</tr>
<tr>
<td>( V_y = \beta_2 g_y + \frac{\lambda_3 \rho g_y y}{B_x} + \frac{\psi_y (B_x)}{B_x} + C_3 )</td>
<td>(sum of integrals from sub-equations #2, #23)</td>
</tr>
<tr>
<td>( V_z = \beta_3 g_z + \frac{\lambda_7 \rho g_z z}{B_x} + \frac{\psi_z (B_x)}{B_x} + C_4 )</td>
<td>(sum of integrals from sub-equations #3, #25)</td>
</tr>
<tr>
<td>( B_x (x, y, z, t) = )</td>
<td>(sum of integrals from sub-equations #6, #7, #8, #9, #10, #16, #19, #22, #24)</td>
</tr>
<tr>
<td>( B_y = r \rho g_y y + \frac{\lambda_4 \rho g_y y}{V_x} + \frac{\psi_y (V_x)}{V_x} + C_8 )</td>
<td>(sum of integrals from sub-equations #11, #18, #20)</td>
</tr>
<tr>
<td>( B_z = s \rho g_z z + \frac{\lambda_8 \rho g_z z}{V_x} + \frac{\psi_z (V_z)}{V_z} + C_21 )</td>
<td>(sum of integrals from sub-equations #12, #17, #26)</td>
</tr>
</tbody>
</table>
Step 9: The linear part of the relation satisfies the linear part of the equation (in Step 8; and the non-linear part of the relation satisfies the non-linear part of the equati. The solutions are above.

Analogy for the Identity Checking Method: If one goes shopping with American dollars and Japanese yens (without any currency conversion) and after shopping, if one wants to check the cost of the items purchased, one would check the cost of the items purchased with dollars against the receipts for the dollars; and one would also check the cost of the items purchased with yens against the receipts for the yens purchase. However, if one converts one currency to the other, one would only have to check the receipts for only a single currency, dollars or yens. This conversion case is similar to the linearized N-S equations, where there was no partitioning in identity checking.

Important insight
One observes above that the most important insight of the above solutions is the indispensability of the gravity term in MHD. Observe that if gravity, \( g_x \), were zero, all the non-constant terms in each solution would be zero. These results can be stated emphatically that without gravity forces on earth, there would be no magnetohydrodynamics on earth as is known. It would not therefore be meaningful to write a system of MHD equations without the gravity term, since there would be no magnetohydrodynamics.

Supporter Equation Contributions (see also viXra:1405.0251)
Note above that there are 28 terms in the driver equation, and 27 supporter equations. Each supporter equation provides useful information about the driver equation. The more of these supporter equations that are integrated, the more the information one obtains about the driver equation. However, without solving a supporter equation, one can sometimes write down some characteristics of the integration relation of the supporter equation by referring to the subjects of the supporter equations of the Navier-Stokes equations. For example, if one uses \( \eta \partial^2 B_x / \partial x^2 \) as the subject of a supporter equation here, the curve for the integration relation obtained would be parabolic, periodic, and decreasingly exponential. Using \( \rho(\partial V / \partial t) \) as the subject of the supporter equation, the curve would be periodic and decreasingly exponential. Using \( (\partial p / \partial x) \), the curve would be parabolic.

Comparison of Solutions of Navier-Stokes Equations and Solutions of Magnetohydrodynamic Equations

Navier-Stokes \( x \)-direction solution

\[
V_x(x,y,z,t) = -\frac{\rho g_x}{2\mu}(ax^2 + by^2 + cz^2) + C_1x + C_3y + C_5z + fg\pm \sqrt{hgx + \frac{ngy}{V_y} \pm \frac{ngz}{V_z}} \psi_y(V_y) \psi_z(V_z) \\

P(x) = d\rho g_x x \\
\text{arbitrary functions}
\]

For magnetohydrodynamic solutions, see previous page

1. \( V_x \) for MHD system resembles the \( V_x \) for the Euler solution part of N-S solution.
2. \( P(x) \) for N-S and MHD equations are the same.
3. \( V_y \) and \( V_z \) for MHD are different from those of N-S solution.
4. \( B_x \) is parabolic and resembles \( V_x \) for N-S, except for the absence of the square root function.
5. \( B_y \) and \( B_z \) resemble the Euler solution part of the N-S solution.
Conclusion

The author proposed and applied a new law to solve the system of magnetohydrodynamic equations. This law states that in magnetohydrodynamics, all the other terms in the system of equations divide the gravity term in a definite ratio, and each term utilizes gravity to function. The experience gained in solving the Navier-Stokes equations guided the author to solve the MHD equations.

It was shown that without gravity forces on earth, there would be no magnetohydrodynamics on earth as is known. The equations in the system of equations were added to produce a single equation which was then integrated. Ratios were used to split-up the single equation, and the resulting sub-equations were readily integrated; and even, the nonlinear sub-equations were readily integrated. Twenty-seven sub-equations were integrated. The linear part of the relation obtained from the integration of the linear part of the equation satisfied the linear part of the equation; and the relation from the integration of the non-linear part satisfied the non-linear part of the equation. Comparison of the solutions of MHD equations with the solutions of the N-S equations revealed the following:

(a) $V_x$ for MHD system resembles the $V_x$ for the non-linear part of the N-S solution; (b) $P(x)$ for N-S and MHD equations are the same; (c) $V_y$ and $V_z$ for MHD are different from those of N-S solutions; (d) $B_x$ is parabolic and resembles $V_x$ for N-S solution, except for the absence of the square root function; and (e) $B_y$ and $B_z$ resemble the non-linear part of N-S solution.

By solving algebraically and simultaneously for $V_x$, $V_y$, $V_z$, $B_x$, $B_y$, $B_z$, the solutions could be expressed in term of $x, y, z$ and $t$.

In applications, the ratio terms may perhaps be determined using information such as initial and boundary conditions or may have to be determined experimentally. Finally, for any magnetohydrodynamic design, one should always maximize the role of gravity for cost-effectiveness, durability, and dependability. Perhaps, a law for magnetohydrodynamics should read "Sum of everything else equals $\rho g$"; and this would imply the proposed new law that the other terms in the system of equations divide the gravity term in a definite ratio, and each term utilizes gravity to function.

Note: The liquid pressure, $P$ at the bottom of a liquid of depth $h$ units is given by $P = \rho gh$.

From the MHD solutions in this paper, $P(x) = b \rho g x$ from integrating $\frac{dp}{dx} = b \rho g$ where $b$ is ratio term. Each of the other terms in the MHD equation must also be set equal to the product of a ratio term and $\rho g$. This result implies that the approach used in solving the MHD equations is sound.

P.S.
The author spent more time on "vector juggling" than on the integration of the equations, since no complete system without vector notation was available either in textbooks or on-line. The integration took less time because of the experience with the N-S equations. Any error in the vector juggling part, if any, can be integrated within minutes.

References:
For paper edition of the above paper, see Appendix 8 of the book entitled "Power of Ratios" by A. A. Frempong, published by Yellowtextbooks.com. Since ratios were the key to splitting the combined 28-term MHD equation into sub-equations and solving them, the solutions have also been published in the "Power of Ratios" book which covers definition of ratio and applications of ratio in mathematics, science, engineering, economics and business fields.