New version of General Relativity that unifies Mass and Gravitation in a common 4D Theory

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Eintein's General Relativity does not explain recent enigmas of astrophysics such as dark energy or accelerating universe. A thorough examination of the Einstein Field Equations (EFE) highlights four inconsistencies. Solving these inconsistencies and bringing closer the EFE to the Higgs mechanism fully explains the mass and gravity phenomena. The main interest of this study is to propose a formula of mass in 4D, $m = f_{(x,y,z,t)}$, that solves several basic enigmas fully demonstrated with mathematics but still not explained with logic and good sense such as "How spacetime can be curved by mass?", or 'What is the nature of the gravitational force?", or "What is the mechanism of conversion of mass into energy ($E = mc^2$)?". This paper also solves the main enigmas of astrophysics and leads to interesting explanations of quantum mechanics.

1. Introduction

When Einstein developed General Relativity (GR), knowledge in physics was very poor. In the 1910s, the atom was regarded as a pudding in which raisins represented electrons. The proton and neutron were discovered later, in, respectively, 1919 and 1932. So, GR was grounded on an erroneous view of reality.

In the 1910s, Einstein had an incredible intuition: "Mass curves spacetime". He developed GR upon this assumption and his theory was fully verified by experiment in 1919. However, despite this validation, no one is able to explain how mass curves spacetime. The mathematics of GR are well known, but the mechanism of curvature of spacetime by mass remains obscure.

Since the 1910s, physics has evolved considerably. So, we must reconsider GR in the current context, excluding the "pudding model" that does not exist anymore, but including the Higgs theory. This leads to a new version of GR which fully explains mass and gravity, and several enigmas of physics, such as:

- How can mass curve spacetime?
- What, exactly, is the mechanism of the creation of mass?
- What is the mechanism of the attractive force of gravitation?
- Does the Higgs field permeate all space? (not proven).
- What is the mechanism of the increase of the mass of relativistic particles?
- How mass can be converted into energy in $E = mc^2$ and conversely? ...

2. Relation Between Volume and Mass

One tends to assume that there exists only one kind of volume. This view must be dropped because some volumes curve spacetime, others do not. Since spacetime is at the heart of GR, then to understand the mechanism of curvature of spacetime, the first thing to do is to separate the volumes that curve spacetime from those that do not.

In reality, we have not only one type of volume, as is thought, but five (see Appendix H). The three major classes of volumes are described below.

2.1 Matter Volumes

These volumes are those of nuclear matter and have been well known since the 1930s. Experiments have shown that the radius of a nucleus is $R = A^{1/3}r_0$, where r_0 is about 1.25 fm. Taking the mass number A = 1 and the nucleon mass as 1.67×10^{27} kg, we get

$$M = kV = f_{(x,y,z,t)} \qquad (k = 2.3 \cdot 10^{17} \text{kg/m}^3) \qquad (1)$$

A more accurate formula which gives better results is proposed in Appendix C:

$$M = \frac{c^2}{G_0} \epsilon_v \frac{V}{S} \,\delta = f_{(x,y,z,t)} \qquad (2)$$

where

$$\begin{split} M &= \text{mass (kg)}, \\ V &= \text{volume (m^3)}, \\ S &= \text{surface (m^2)}, \\ c &= \text{speed of light (m/s)}, \\ \epsilon_v &= \text{elasticity of spacetime,} \\ G_0 &= \text{universal constant of gravitation,} \\ \delta &= \text{density of spacetime, relative to a flat spacetime.} \end{split}$$

It is important to see that whatever the formula (1) or (2), we have

$$M = f_{(x,y,z,t)} \tag{3}$$

2.2 Empty Volumes

Empty Volumes are simply vacua, but can also be found in various forms, such as the volumes of the orbitals of atoms. They exist but are transparent regarding spacetime. Empty Volumes have no mass and they do not curve spacetime (remember that GR considers that the curvature of spacetime is produced by mass).

2.3 Apparent volumes

Objects we use daily are mainly made of molecules, which are combinations of "Matter Volumes" (nuclei and electrons) that have mass, and "Empty Volumes" (orbitals, gaps between atoms or molecules...), without mass:

Apparent Volumes =
$$\sum$$
 Matter Volumes + \sum Empty Volumes.

3. What We Call "Volumes"

3.1 Definition

These three classes of volumes have been well known since the 1930s. However, it seems that we have great difficulty separating the volumes that curve spacetime from the volumes that do not. The result is that our current view of "Volumes" is correct in daily life, but wrong in physics, more particularly in astrophysics.

1/ In daily life, we use exclusively "Apparent Volumes". We have the feeling that Mass and Volume are two different quantities because in "Apparent Volumes", the proportion of (*Matter Volumes*)/(*Empty Volumes*) varies from one atom to another, from one molecule to another, from one object to another.

2/ In physics, we must exclude "Empty Volumes" by a thought experiment since this kind of volume is inert regarding spacetime. They are massless. There remains "Matter Volumes", i.e. the only volumes that curve spacetime. Therefore, it appears that Mass and "Matter Volume" are two different views of the same quantity, since Mass can be converted into "Matter Volume" and conversely, Eqs (1–3).

$$M = f_{(MatterVolume)} \tag{4}$$

Since this article is addressed to physicists, we are concerned only with the second definition. In other words, we must disregard massless "Empty Volumes".

3.2 Substituting Mass by (Matter) Volume

Since $M = f_{(MatterVolume)}$, Eqs (1-3), it becomes possible to replace mass by (matter) volume. This means that, instead of saying

"Spacetime is curved by mass",

we should say:

"Spacetime is curved by (Matter) Volume"

3.3 Interpretation

The substitution of M by V is very important because it explains the curvature of spacetime and, consequently, many enigmas of modern physics.

To date, the mechanism of curvature of spacetime by mass is obscure and no one can give a rational explanation of this strange phenomenon. On the contrary, it is easy to understand how spacetime can be curved by a (matter) volume.

For example, if we plunge an object into a glass filled with water, as expected, it is the volume of the object (not its mass) that displaces the water.

This example can be transposed to spacetime as follows.

Consider a particle of volume V crossing a cube of volume 1,000,000 V. Here, the particle is a (matter) volume whereas the cube contains a vacuum, i.e. (empty) volumes. Both kinds of volumes have their own spacetime. The internal spacetime of the particle will "push" the surrounding spacetime to make room. So, as in our preceding example, it is the volume of the particle that curves spacetime (not its mass), but since $Mass = f_{(MatterVolume)}$, Eqs (1–3), the result is identical.

To summarize, since all volumes are not equal regarding spacetime, to understand the curvature of spacetime, we must separate the volumes that curve spacetime from the volumes that do not. In other words, we must think "*Matter Volume*" instead of "*Mass*".

4. Towards a New GR

GR can be studied in different ways. This article is grounded on the 1910s original Einstein–Grossmann version. For background on GR, see Appendix E. This section explains and solves four inconsistencies found in GR in the original Eintein–Grossmann version.

4.1 Inconsistency #1: Lack of Homogeneity

The Einstein Field Equations (EFE), without the cosmological constant Δ , are:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G_0}{c^4}T_{\mu\nu} \qquad (5)$$

The left hand side of the EFE is developed in 4D, $f_{(x,y,z,t)}$, whereas the right hand side contains an additional quantity, m, which is hidden in the tensor $T_{\mu\nu}$. Since the m variable is not defined in 4D such as $m = f_{(x,y,z,t)}$, we can consider that m is a 5th dimension. The empirical Einstein Constant $8\pi G_0/c^4$ homogeizes the right hand side of EFE but does not solves the nature of m: 4D variable or 5th dimension?

Such a formulation 4D = 5D is not mathematically wrong per se if a 4D expression of m exists, such as $m = f_{(x,y,z,t)}$. This is not the case. In the EFE, replacing "Mass" by "Matter Volume" solves this inconsistency. More

In the EFE, replacing "Mass" by "Matter Volume" solves this inconsistency. More precisely, substituing m by its expression $m = f_{(x,y,z,t)}$ Eq. (2) in the right hand member of the EFE leads to a correct 4D = 4D formulation of the EFE (Appendix E):

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi\delta_v}{S}J_{\mu\nu} \qquad (6)$$

4.2 Inconsistency #2: Direction of Curvature

In fluid mechanics, the curvature of the fluid made by an object is *convex* (Fig. 1A). The scientific literature shows that the curvature of spacetime produced by a mass such as a black hole is *concave* (Fig. 1B). Since the stress tensor of fluid mechanics is at the origin of the energy–momentum tensor of EFE, it would seem logical to think that the curvature of spacetime is *convex*, not *concave*.



Figure 1. Since the stress tensor and the energy–momentum tensor are identical, there is no reason to replace the convex curvature of the original stress tensor by a concave curvature in GR.

In 1927, Von Laue also thought that the curvature of spacetime could be convex, but he was not followed by his peers. Why did Einstein–Grossmann invert the direction of the curvature (convex in fluid, concave in spacetime)? Why did the community of physicists reject Von Laue's suggestion? The reasons remain obscure.

4.3 Inconsistency #3: Mass or Volume?

In fluid mechanics, the displacement of the fluid is made by the *volume* of the object, not by its *mass*, as shown in Fig. 1A, previous page. Here too, we must keep the original meaning of the stress tensor in the EFE and replace "*Mass*" by "(*Matter*) Volume".

Note: In the 1910s, the atom was built on the plum pudding model. This probably explains why Einstein and Grossmann made this error.

4.4 Inconsistency #4: Pressure or attractive force?

All the elements of the Cauchy stress tensor have a pressure-like dimension (Fig. 2A). Indeed, the fluid exerts a *pressure* on objects (brown arrows in Fig. 1A). Since Einstein and Grossmann built the energy–momentum tensor from the Cauchy stress tensor, each term must keep its original meaning. This is not the case. Question: *Why has the original pressure force been replaced by an attractive force (gravitation)?*



Figure 2. The origin of the energy–momentum tensor of GR, in the original 1910s fluid mechanics version, is the Cauchy stress tensor. Since the stress tensor is included in the energy–momentum tensor, the two tensors must have the same meaning, at least in the spatial part (in pink in the figure).

5. Recapitulation

There is not only one definition of volume, but three. In physics, we must replace Mass by (Matter) Volume. The results are identical since $M = f_{(MatterVolume)}$ but this substitution explains several enigmas of physics.

On the other hand, the energy–momentum tensor is built from the stress tensor. This means that these two tensors must have the same foundation and principle of construction. A close examination of the EFE highlights the following inconsistencies:

#1: 4D=5D. This error disappears if Mass is replaced by (Matter) Volume $M = f_{(x,y,z,t)}$. #2: In the EFE, the curvature of spacetime must be *convex*, not *concave* (Fig. 1A-1B). #3: Here we show that spacetime is curved by (*matter*) volume, not by mass (Fig. 1A). #4: Spacetime exerts a pressure force on objects, not an attractive force (Section 7).

6. Expression of Mass

The above conclusions show that mass is not a fifth dimension $f_{(m,x,y,z,t)}$ but a 4D virtual quantity $f_{(x,y,z,t)}$ (Eqs 1–3), which is created by any "Matter Volume". The mechanism is summarized as follows (Figs. 3A and 5, next pages):

- A "Matter Volume" (not a mass) curves spacetime;
- This curvature is convex, not concave;
- The result is that spacetime exerts a pressure on the surface of the matter volume;
- This pressure increases the resistance to movement of the matter volume;
- ...which leads to a "mass effect": $M = f_{(x,y,z,t)}$.

This explanation is close to the Higgs mechanism. Section 10 covers this subject.

7. What is Gravitation?

The mechanism for gravitation becomes clear and easy to understand if we think "Matter Volume" instead of "Mass" (Fig. 3B):

Gravitation is not an attractive force between masses, but a pressure force exerted by the curvature of spacetime on "Matter Volumes", which tends to bring them closer to each other.

We have seen that the curvature of spacetime is convex, not concave. This convex curvature of spacetime exerts a pressure force on objects, not an attractive force. This is confirmed by a close examination of the two tensors (Fig. 2). We can combine these two assertions into the following expression:

Concave curvature + Attractive force \equiv Convex curvature + Pressure force.

If we assign by convention a positive sign to a convex curvature and to a pressure force, the above assertion gives (--) = (++). As we see, this does not change the mathematics of GR. The result is identical. For example, a pressure on one side of a sheet of paper (concave curvature) produces the same effects as an attractive force (convex curvature) on the other side. In this example we get: (+-) = (-+).



Figure 3. Mass and gravitation are two similar phenomena. They are the consequence of the elasticity of spacetime (Einstein) which exerts a pressure on one (mass) or more (gravity) "Matter Volumes".

8. Curvature of Spacetime vs. Pressure

Figure 4a shows the curvature of spacetime produced by the "Matter Volumes" of two spheres. On the blue sphere, spacetime has two different magnitudes of curvature:

- Point L (Left side): The curvatures of spacetime of the two spheres are added.

- **Point R** (Right side): The curvature of spacetime of the red sphere is subtracted from that of the blue one because the two curvatures are opposed.



Figure 4. (a) The curvature of spacetime produced by the red sphere interacts with that produced by the blue one. (b) As a consequence, a difference of pressure on the surface of each sphere appears (gravity).

The difference of the curvature of spacetime on each side of the two spheres leads to a difference of pressure (black and orange arrows on Fig. 4b). Fig. 4a and 4b are two different views of the same phenomenon: *curvature of spacetime vs. pressure*.

The same principle also applies if the two spheres are the Earth and the Moon. At the Lagrangian point (green area), the curvature of spacetime produced by the Earth is canceled by that produced by the Moon. The curvature of spacetime is null at this point, and gravity disappears.

9. Conclusions

This article shows that spacetime is not curved by mass as we think, but by a special kind of volume: "Matter Volume". This leads to a relation $M = f_{(x,y,z,t)}$. As a result, we can replace mass by "Matter Volume". This substitution explains why the curvature of spacetime is convex, not concave, and why it makes a pressure on the surface of volumes. The whole mass and gravity mechanism is shown in Fig. 5.



Figure 5. Mass and gravity mechanism. Relation between the curvature of spacetime, "Matter Volumes", "Empty Volumes", "Apparent Volumes" (object), mass, and gravitation

Therefore, the originality of this paper lies in these simple observations:

1/ Mass can be replaced by "Matter Volumes": $M = f_{(V)}$;

2 / The EFE have a lack of homogeneity: 4D = 5D (here solved)

3 / The curvature of spacetime is convex, not concave;

4 / Mass is the result of the pressure of spacetime on "Matter Volumes";

5 / Gravitation is a pressure force made by spacetime, not an attractive force.

10. Higgs Boson

The Higgs boson is a reality but the Higgs theory raises some questions:

1 – The Higgs field does not have an associated potential (Poisson's law);

2 – Not all physicists agree with the Big Bang theory, therefore, saying "*The Higgs field was created during the Big Bang*" is far to be proven;

3 – The Higgs boson has always been detected near particles: Question: does the Higgs field permeate all space (not proven), or is it locally created by the interaction?

4 – The Higgs field does not explain the origin of mass (here, "Matter Volumes");

5 – It does not explain gravity (a pressure force, not an attractive force),

6 – The formulation of the Higgs field (HF) is obscure. Is it in 4D such as $HF = f_{(x,y,z,t)}$?, or is it a fifth dimension (HF, x, y, z, t)?

Here we solve all these inconsistencies of the Higgs theory. This leads to consider that the Higgs field is nothing but a 4D Newtonian field in spacetime created by the "Matter Volume" of the particle expected to get mass (Fig. 3), which creates a potential. This also means that, out of any mass, the Higgs Field does not exist.

11. Appendices

The appendices are located in the following pages. They cover:

A – New version of Newton's Law calculated from Matter Volumes;

B – The Schwarzschild metric calculated from Matter Volumes;

C – Expression of the mass effect in 4D: $m = f_{(x,y,z,t)}$;

D – Explanation of the increase of the mass of relativistic particles;

E – New formulation of the EFE: $R_{\mu\nu} - 1/2g_{\mu\nu}R = (8\pi\delta_v)/S$. $J_{\mu\nu}$;

F – Connexion with the Von Laue geodesics;

G – New version of the Equivalence Principle;

H – Black Holes Simulation;

I - Classes of Volumes;

J – Applications.

12. References

The bibliography is located at the end of the following appendices.

[1] This paper has been registered at INPI under the following references: 238268, 238633, 244221, 248427, 258796, 261255, 268327, 297706, 297751, 297811, 297928, 298079, 298080, 329638, 332647, 335152, 335153, 339797, 12-01112.

Mass and Gravitation

Mathematical Demonstrations

The following appendices cover:

- A New Version of the Newton's Law,
- B New Version of the Schwarzschild Metric,
- C Expression of the Mass Effect in 4D,
- D Explanation of the Increase of the Mass of Relativistic Particles,
- E Partial Rewriting of the Einstein Field Equations (EFE),
- F Von Laue Geodesics,
- G New Version of the Equivalence Principle,
- H Black Holes Simulation
- I Classes of Volumes
- J Applications

Bibliography

New Version of the Newton's Law

A-1 Introduction

The Newton's Law is obtained from the Einstein Field Equations (EFE). It is a particular solution for a spherical static symmetry object using the weak field approximation. Here we show that the Newton's Law can be easily obtained replacing mass by matter volume. This method is much more simpler than calculating the Newton's Law from the EFE.

A-2 Bulk Modulus

The bulk modulus K_B of a substance measures the substance's resistance to uniform compression. It is defined as the pressure increase needed to cause a given relative decrease in volume (fig. A-1).



Figure A-1: Bulk modulus

$$K_B = -V \frac{\Delta P}{\Delta V} \qquad (1)$$

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Starting with the Navier-Stokes Equations of the Fluid Mechanics, Einstein, helped by Grossman, demonstrated in the 1910s that spacetime:

- Can be identified to a newtonian fluid,
- Returns to its rest shape after having applied a stress (properties of elasticity).

Therefore, the Bulk Modulus (Equ. 1), which is a version of the Cauchy Tensor of the Fluid Mechanics, can also be applied to spacetime. It means that the displacement of spacetime made by a matter volume exerts a pressure on the surface of the latter (fig. A-1).

Note: This subject is also covered in Appendix E.

A-3 Elasticity Law

Elasticity phenomena follow the well-known logarithmic law:

$$\epsilon = \ln\left(\frac{R - \Delta R}{R}\right) \qquad (2)$$

with ϵ = coefficient of elasticity.

The Schwarzschild Metric gives an order of magnitude of the curvature of spacetime, which is infinitesimal. For example, the ratio curvature of spacetime/radius, or $\Delta R/R = GM/Rc^2$, is 1.4166E-39 for the proton, with M = 1.672E-27 kg, R = 8.768E-16 m, G = 6.674E-11, c = 8.987E+16. See Appendix B for the meaning of $\Delta R/R$ and of the factor 2 in $2GM/Rc^2$.

Under these conditions, whatever the formula used, logarithmic or not, the curvature of spacetime can be considered as a linear function since we are working on an infinitesimal segment near to the point zero. So, Equ. (2) becomes in first order approximation:

$$\epsilon_R \approx \frac{\Delta R}{R}$$
 (3)

or, with volumes:

$$\epsilon_V \approx \frac{\Delta V}{V}$$
 (4)

For the moment, the linear ϵ_R and volumetric ϵ_v coefficients of elasticity of spacetime are unknown.

A-4 Curvature of Spacetime

A matter volume V inserted into spacetime pushes it to make room (Fig. A-2). So the following volumes are identical:

$$V = V_1 = V_2 \dots = V_n$$
 (5)



Figure A-2: Volumes V, V1, V2, V3... are identical

Since the curvature of spacetime is a linear function, the coefficient of elasticity of spacetime ϵv can be considered constant. So, combining (4) and (5) gives:

$$\Delta V = \Delta V_1 = \Delta V_2 \dots = \Delta V_n \tag{6}$$

However, there should not be any confusion between a simple displacement of spacetime, Vx, produced by the insertion of a matter volume into a flat spacetime, and the curvature $\Delta V_x = \epsilon_v V_x$ due to the elasticity of spacetime (fig. A-3).



Figure A-3: Simple displacement (V_x) vs curvature of spacetime (ΔV_x)

A-5 Solving $\Delta x = f_{(\Delta R)}$

Since the ΔV 's are infinitesimal, the volume ΔVx is simply the product of Δx by the surface Sx (fig. A-4):

$$\Delta V_x = S_x \ \Delta x = 4\pi d^2 \Delta x \tag{7}$$

The volume ΔV_R is also the product of ΔR by the surface S_R :

$$\Delta V_R = S_R \ \Delta R = 4\pi R^2 \Delta R \tag{8}$$



Figure A-4: Displacement and curvature at distances R and d

From (6) we have:

$$\Delta V_R = \Delta V_x \qquad (9)$$

Combining (7), (8) and (9) gives:

$$4\pi R^2 \Delta R = 4\pi d^2 \Delta x \qquad (10)$$

Finally, we get:

$$\Delta x = \frac{R^2}{d^2} \Delta R \qquad (11)$$

Where:

- R is the radius of the matter volume V_R ,
- ΔR is the curvature of spacetime on the surface of the matter volume V_R ,
- d is the distance of the point of measurement,
- Δx is the curvature of spacetime at distance d.

A-6 Curvature Δx vs Mass M

As explained in the main text and in the first part of this appendix, a relation exists between the curvature of spacetime, ΔR (or Δx at a distance "d" from R), and the "mass effect" of the object:

$$\Delta R = f_{(M)} \tag{12}$$

This is the pressure that produces the mass effect. This suggests that the latter is inversely proportional to the surface S, or $[1/L^2]$, as any pressure does. On the other hand, it is obvious that the mass effect is also proportional to the volume, or $[L^3]$. Therefore, the dimensional quantity of the mass effect is $[1/L^2][L^3] = [L]$. In other words, $[M] \equiv [L]$.

At this point, we don't know the relation between ΔR and M but, in referring to Einstein's works, we have good reasons to believe that this relation is a simple linear function like:

$$\Delta R = KM \tag{13}$$

 \ldots where K is an constant having the dimensional quantity of [L/M].

Here we show that [K] = [L/M], but we will see later that $K = G/c^2$. This result is in line with the dimensional quantity of some terms of the Schwarzschild Metric: $2\epsilon = 2\Delta r/r = 2GM/rc^2$ (see Appendix B). The challenge, now, is to calculate K to get the Newton's Law.

A-7 The Newton Law

Porting (13) in (11) gives:

$$\Delta x = \frac{R^2}{d^2} \ KM \qquad (14)$$

or

$$\frac{\Delta x}{R^2} = K \frac{M}{d^2} \qquad (15)$$

Since x = ct, replacing R^2 by c^2t^2 gives:

$$\frac{\Delta x}{c^2 t^2} = K \frac{M}{d^2} \qquad (16)$$

or :

$$\frac{\Delta x}{t^2} = c^2 K \frac{M}{d^2} \qquad (17)$$

The value $\Delta x/t^2$ has the dimensional quantity of an acceleration [L/T²]. So, replacing this fraction by the acceleration symbol "a" (see the note below), we get:

$$a = c^2 K \frac{M}{d^2} \tag{18}$$

Notes: Δx is an infinitesimal quantity, not a differential quantity such as dx. Moreover, we are working in a linear segment of the elasticity of spacetime. In such a situation, $\Delta x/\Delta t \approx x/t$.

On the other hand, the multiplication of a constant c^2 by a second constant K gives another constant. So, we can replace the product c^2K by a new and unknown constant for the moment, G for example:

$$c^2 K = G \tag{19}$$

or (this equation isn't necessary here but will be used for the calculation of the Schwarzschild Metric in Appendix B.)

$$K = \frac{G}{c^2} \qquad (20)$$

Porting (19) in (18) gives:

$$a = G\frac{M}{d^2} \qquad (21)$$

To be consistent, this unknown constant G must have the same dimensional quantity of the product $c^2 K$ (Equ. 19):

- c^2 : Dimensional quantity $\Rightarrow [L^2/T^2]$
- K : Dimensional quantity $\Rightarrow [L/M]$ (see paragraph A-6)

So,

The dimensional quantity of this new constant G is

$$[c^2K] = [L^2/T^2][L/M] = [L^3/MT^2].$$

On the other hand, we know that a force is the product of an acceleration by a mass, here "m". Therefore, Equ. (21) can be written as follows:

$$F = G \ \frac{Mm}{d^2} \tag{22}$$

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For the moment, G is unknown but we must note that:

- G is a constant,
- Its dimensional quantity is $[L^3/MT^2]$.

So, we can identify G to the constant of gravitation issued from experimentation:

$$G = 6,67428.E - 11 \tag{23}$$

In other words,

Equ. (22) can be identified to the Newton Law of Universal Gravitation.

New Version of the Schwarzschild Metric

B-1 Introduction

Here we show that the Schwarzschild Metric can be easily obtained using matter volumes instead of masses. The following demonstration doesn't require tensor knowledge.

B-2 The Minkowski Metric

The expression of the Minkowski Metric, in spherical coordinates, is:

$$ds^{2} = -c^{2}dt^{2} + dr^{2} + r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2}) \qquad (1)$$

The Schwarzschild Metric refers to a static object with a spherical symmetry. It is built from a Minkowski Metric, in spherical coordinates, with two unknown functions: A(r) and B(r):

$$ds^{2} = -B_{(r)}c^{2}dt^{2} + A_{(r)}dr^{2} + r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2}) \qquad (2)$$

The Minkowski Equation must follow the Lorentz Invariance in Special Relativity (SR) or General Relativity (GR). To get this invariance, we must set $A_{(r)} = 1/B_{(r)}$. Details of calculus are described in Appendix E and in books concerning GR. So:

$$B_{(r)}A_{(r)} = 1$$
 (3)

B-3 The Schwarzschild Metric

To calculate the Schwarzschild Metric, we can start with fig. B-1 (next page), which is issued from the theory described in the main article, where:

- d_{rout} is an elementary differential radial variation outside of any mass,
- d_{rin} is an elementary differential radial variation inside a Schwarzschild spacetime,
- r is the point of measurement.



Figure B-1: Spacetime has been reduced to 1D

Supplementary Information A gives:

$$\epsilon = \ln\left(\frac{r - \Delta R}{r}\right) \qquad (4)$$

where:

- ϵ is a coefficient of the increase of spacetime curvature at distance r,
- ΔR is the initial curvature of spacetime produced by the matter volume,

The order of magnitude of ϵ is 10E-39. So, we can use the first order approximation from Supplementary Information A, Equ. (3):

$$\epsilon \approx \frac{\Delta R}{r} \qquad (5)$$

Since ϵ is a simple coefficient, we can calculate the relation between two differential elementary radius dr(out) and dr(in), out and in a gravitational field:

$$dr_{in} = (1+\epsilon) \ dr_{out} \tag{6}$$

Since $\epsilon \ll 1$, Equ. (6) becomes:

$$dr_{in} = \frac{1}{(1-\epsilon)} dr_{out} \qquad (7)$$

or, elevating in square:

$$dr_{in}^2 = \frac{1}{(1-\epsilon)^2} \ dr_{out}^2 \qquad (8)$$

Developing the denominator $(1-\epsilon)^2 = 1-2\epsilon+\epsilon^2$ and ignoring the last term ϵ^2 , we obtain:

$$dr_{in}^2 = \frac{1}{(1-2\epsilon)} dr_{out}^2$$
 (9)

This result is nothing but the radial component of the Schwarzschild Metric, that is to say the function A(r) of dr^2 in (2). Then, the calculation of B(r) is immediate, taking into account that A(r)B(r) = 1 from Equ. (3). So:

$$A_{(r)} = \frac{1}{(1 - 2\epsilon)} \qquad (10)$$
$$B_{(r)} = (1 - 2\epsilon) \qquad (11)$$

Thus, Equ. (2) becomes:

$$ds^{2} = -(1 - 2\epsilon)c^{2}dt^{2} + \frac{1}{(1 - 2\epsilon)}dr^{2} + r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2}) \quad (12)$$

In Appendix A "The Newton Law", we have got $\Delta R = KM$ (Equ. 13). Since $K = G/c^2$ (Appendix A Equ. 20), Equ. (5) can be rewritten as:

$$\epsilon = \frac{\Delta R}{r} = \frac{KM}{r} = \frac{GM}{rc^2} \qquad (13)$$

Finally, porting this expression in Equ. (12) gives:

$$ds^{2} = -\left(1 - \frac{2GM}{rc^{2}}\right)c^{2}dt^{2} + \frac{1}{\left(1 - \frac{2GM}{rc^{2}}\right)}dr^{2} + r^{2}(d\theta^{2} + \sin^{2}\theta d\phi^{2}) \quad (14)$$

B-4 Conclusions

This new calculus of the Schwarzschild Metric, which is exclusively based on matter volumes, gives identical results than developing the EFE in the special case of a static spherical symmetry. This is due to the fact that the origin of EFE is the Fluid Mechanics, which is itself based on volumes, not on masses (see Appendix E).

Expression of the Mass Effect in 4D

C-1 Expression of "m"

In the main paper and appendix A, we have seen that the displacement of spacetime V_R is equal to that of the matter volume, V, which produces this displacement (fig. C-1):



 $V_R = V \qquad (1)$

Figure C-1: The curvature of spacetime

On the other hand, the curvature of spacetime is:

$$\Delta V_R = \epsilon_v \ V_R \qquad (2)$$

Porting (1) in (2) gives:

$$\Delta V_R = \epsilon_v \ V \qquad (3)$$

The radial curvature of spacetime, ΔR , at the surface of M, is calculated dividing the volume by the surface:

$$\Delta R = \frac{\Delta V_R}{S} \qquad (4)$$

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Replacing ΔV_R by its expression (3) gives:

$$\Delta R = \epsilon_v \frac{V}{S} \qquad (5)$$

On the other hand, we have calculated the Newton's Law starting with equation (13) of Appendix A:

$$\Delta R = KM \qquad (6)$$

...where K is an unknown constant having the dimensional quantity of [L/M]. Porting (6) in (5) gives:

or

$$KM = \epsilon_v \frac{V}{S} \qquad (7)$$
$$M = \frac{\epsilon_v V}{KS} \qquad (8)$$

Porting in equation (8) the expression of K given in equation (20) of Appendix A gives the expression of the "mass effect". Here, we have added a new coefficient, ρ (see explanation below):

$$M = \frac{c^2}{G_0} \epsilon_v \frac{V}{S} \rho \tag{9}$$

with:

M = Mass effect (kg)

V = Volume of the matter volume (m³)

S = Surface of the matter volume (m²)

 ϵ_v = Coefficient of volumetric elasticity of spacetime in a flat spacetime. This parameter is unknown but can be calculated from the mass/diameter of spherical particles such as some leptons or "magic" nuclei. See the next sections.

c = Speed of the light (m/s)

 $G_0 =$ Universal constant of gravitation

 ρ = Density of surrounding spacetime relative to a flat spacetime. This parameter is equal to 1 in a Minkowsky Spacetime.

It seems useful to differentiate ϵ_v , the coefficient of elasticity of spacetime in a flat spacetime, and ρ , the density of surrounding spacetime. We could merge these two parameters in one common parameter since we are faced with two dimensionless coefficients. In both cases, result is identical. However, we must note that the "proper coefficient of elasticity of spacetime", as proper length in special relativity, must be measured in a flat spacetime. This is why the two parameters have been separated. For example, the particle may be located in a riemanian spacetime, i.e. in a field produced by another particle. In this case, since the mass effect is a function of the curvature of spacetime, we need to know the latter before any calculation.

C-2 Case of a sphere

In the particular case of a sphere, we have $V = 4/3\pi R^3$ and $S = 4\pi R^2$. Thus, equation (9) becomes:

$$M = \frac{\epsilon_v R c^2}{3 G_0} \rho \tag{10}$$

C-3 Nuclei

Nuclei aren't spherical generally. Since we don't know exactly their shape, it is not possible to apply equation (10) to calculate their mass effect.

The semi-empirical mass formula (SEMF), sometimes called the Bethe-Weizsacker mass formula, is used to approximate various properties of an atomic nucleus. It is based partly on the liquid drop model proposed by George Gamow, and partly on empirical measurements. From the SEMF formula, the radius R is defined as $R = 1.2A^{1/3}$, A being the mass number. In reality, the right member does not mean that the mass vs. the radius follows a $A^{1/3}$ law. It is the arrangement of nucleons inside the nucleus that follows this rule. The result is equivalent but the significance is different.

We must also note that a surface component also exists in the Bethe-Weizsacker expression. It means that, early in 1937, Bethe and Weizsacker predicted the equation (9) here demonstrated.

We must keep in mind that nuclei are made of empty and matter volumes. The space between nucleons may vary from one nucleus to another. Thus, it is necessary to know the arrangement of nucleons with accuracy before any calculations. A particular case are magic nuclei (nuclei having a null quadripolar moment) because they are supposed to be spherical. It will be interesting to make accurate experiments on the relation volume/surface/mass effect of magic nuclei.

On the other hand, some nuclei have a halo made of empty volumes that are not relevant in mass calculation. This is the case for example of the ¹¹Li (3p8n), which has empty volumes between the ⁹Li and the 2n orbitals. These exceptions highlight the difficulty to make accurate calculations of the mass effect. In all cases, before any calculation, we must know exactly the geometry of matter and empty volumes inside the nucleus. It means that the calculus of the mass from the geometry of the nucleus (equation 9) is not as simple as it sounds.

As a direct consequence of the proposed theory, it could be possible that a relationship exists between the sphericity of particles or nuclei and the accuracy of measurements. This deduction suggests that leptons could be spherical since their mass effect is known with an excellent accuracy. This is also the case of some particles such as the proton, neutron, or π meson. Inversely, this is not the case of quarks. This means that quarks could have a non-spherical or complex shape.

Explanation of the Increase of the Mass of Relativistic Particles

D-1 Introduction

The increase of the mass of relativistic particles is covered by special relativity. However, this phenomenon remains particularly obscure and, to date, we are still unable to explain with simple words, i.e. without using mathematics, why does the mass increase with velocity. The proposed theory gives a simple and rational explanation of this strange phenomenon.

D-2 Length contraction

Special relativity states that, at relativistic speed, times *expand*, lengths *contract* and angles are *modified*. A simple demonstration is given in 1923 by Einstein himself in his book "The Theory of Special and General Relativities". The length contraction is defined by the formula

$$l_m = l_0 \sqrt{1 - \frac{v^2}{c^2}} \qquad (1)$$

with

- l_m = Measured length
- $l_0 =$ Proper length
- v = Speed of object
- c = Speed of light

D-3 Mass increase

Lets consider a particle at rest (fig. D-1a next page). Its matter volume produces a curvature of the spacetime. Geodesics of spacetime are spaced of l_0 .

If this particle moves at a relativistic speed "v" (fig. D-1b), spacetime geodesics seems to shrink. This is the well-known phenomenon of length contraction. The closer the geodesics are to each other, the more important spacetime density is, according to equation (9) of Appendix C. In other words, the curvature of spacetime is inversely proportional to the space between two geodesics (see note 1). So, relation (1) becomes:

$$\Delta R_m = \frac{\Delta R_0}{\sqrt{1 - v^2/c^2}} \qquad (2)$$

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with

- $\Delta R_m =$ Infinitesimal element of the curvature of spacetime measured (external observer)
- $\Delta R_0 =$ Infinitesimal element of the proper curvature of spacetime
- v = Speed of the particle
- c = Speed of light



Figure D-1: Since the spacetime is more dense in (b), the mass effect increases.

Since the curvature is a function of mass $\Delta R = KM$ (see Appendix A, equations 13 and 20), we can replace the curvature of spacetime ΔR_m of equation (2) by the mass effect m, and the proper curvature of spacetime ΔR_0 by the proper mass effect m_0 (see notes 2 and 3).

$$m = \frac{m_0}{\sqrt{1 - v^2/c^2}}$$
 (3)

Thus, the theory presented in this paper based on matter volumes instead of masses, and figure D-1, give a very simple and rational explanation of the mass increase of relativistic particles.

Note 1: The spacetime curvature is the difference of displacement ΔR of a geodesic vs. to the same geodesic in a Minkowski space. As shown in this article, the Schwarzschild metric gives an order of magnitude of this spacetime curvature: 1.4166 E-39 meters for the proton on its surface. This value is much smaller at distance r. Thus, regardless of the function used, the portion of the curve on which we work is linear. Taking this linearity into account, there is no objection to consider that the curvature is inversely proportional to the space between two geodesics.

Note 2: The nature of expression $\Delta R = KM$ is not relevant because this section covers exclusively the calculation of the coefficient to be applied to a proper value to get the measured value. This coefficient is noted γ (or $1/\gamma$) in scientific literature. This means that the relationship between the spacetime curvature and the mass effect is not affected by this study. For example, if we make 5 measurements of the curvature of spacetime at different speeds, we will not have 5 different relationships between ΔR and M, but only one applicable in all cases ...but we will have 5 different coefficients γ .

Note 3: The principles of special relativity state that the measurement of the mass of a relativistic particle increases. However, the converse is also true if we swap the reference systems. If we could pick up a measuring device on a particle in movement, this device would indicate that our spacetime, that in which we live, is much more dense as we see it. Thus, a section of the LHC for example, with a mass of 3 tons, measured from a device located on the particle in motion, would have a mass of 3000 tons if $\gamma = 1000$. From our point of view, the mass of a relativistic particle increases, but from the particle's point of view, it is our world that increases. In all cases, the proper mass of the particle or that of our world remains unchanged. This "relative view" is often misunderstood.

Note 4: Many physicists think that the mass, so the "matter volume", of relativistic particles really increases. In reality, it is the mass effect due to the apparent compression of spacetime that increases. The volume remains unchanged. On Earth, we consider that "mass" is an intrinsic value of a particle, such as the volume. It is not true. Since the mass effect comes from the pressure of spacetime on the particle, it is a virtual effect, such as pressure, speed, force, energy... On the other hand, the mass effect depends of the surrounding density of spacetime. Thus, for example, if the spacetime density was two times higher, the mass effect would be twice as important as well, but the intrinsic characteristics of the particle would remain unchanged. This explanation is shown in the graphic of fig. D-1.

Partial Rewriting of the Einstein Field Equations (EFE)

E-1 General Relativity Origins

In the 1910s, Einstein studied gravity. Following the reasoning of Faraday and Maxwell, he thought that if two objects are attracted to each other, there would be some medium. The only medium he knew in 1910 was spacetime. He then deduced that the gravitational force is an indirect effect carried by spacetime. He concluded that any mass perturbs spacetime, and that the spacetime, in turn, has an effect on mass, which is gravitation. So, when an object enters in the volume of the curvature of spacetime made by a mass, i.e. the volume of a gravitational field, it is subject to an attracting force. In other words, Einstein assumed that the carrier of gravitation is the *curvature of spacetime*. Thus, he tried to find an equation connecting:

1. The curvature of spacetime. This mathematical object, called the "Einstein tensor", is the left hand side of the EFE (Eq. 1).

2. The properties of the object that curves spacetime. This quantity, called the "Energy–Momentum tensor", is $T_{\mu\nu}$ in the right hand side of the EFE.

Curvature of spacetime \equiv Object producing this curvature

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu} \qquad (1)$$

Einstein Tensor (curvature of spacetime)

Thus, Einstein early understood that gravity is a consequence of the curvature of spacetime. Without knowing the mechanism of this curvature, he posed the question of the special relativity in a curved space. He left aside the flat Minkowski space to move to a Gaussian curved space. The latter leads to a more general concept, the "Riemannian space". On the other hand, he identified the gravitational acceleration to the inertial acceleration (see the Appendix G "New Version of the Equivalent Principle").

The curvature of a space is not a single number, though. It is described by "tensors", which are a kind of matrices. For a 4D space, the curvature is given by the Riemann-Christoffel tensor which becomes the Ricci Tensor after reductions. From here, Einstein created another tensor called "Einstein Tensor" (left hand side of equation 1) which combines the Ricci Curvature Tensor $R_{\mu\nu}$, the metric tensor $g_{\mu\nu}$ and the scalar curvature R (see the explanations below).

Fluid Mechanics

In fluid mechanics, the medium has effects on objects. For example, the air (the medium) makes pressures on airplanes (objects), and also produces perturbations around them. So, Einstein thought that the fluid mechanics could be adapted to gravity. He found that the Cauchy-Stress Tensor was close to what he was looking for. Thus, he identified 1/ the "volume" in fluid mechanics to "mass", and 2/ the "fluid" to "spacetime".

Energy-Momentum Tensor

The last thing to do is to include the characteristics of the object that curves spacetime in the global formulation. To find the physical equation, Einstein started with the elementary volume dx.dy.dz in fluid mechanics. The tensor that describes the forces on the surface of this elementary volume is the Cauchy Tensor, often called Stress Tensor. However, this tensor is in 3D. To convert it to 4D, Einstein used the "Four-Vectors" in Special Relativity. More precisely, he used the "Four-Momentum" vector Px, Py, Pz and Pt. The relativistic "Four-Vectors" in 4D (x, y, z and t) are an extension of the well-known non-relativistic 3D (x, y and z) spatial vectors. Thus, the original 3D Stress Tensor of the fluid mechanics became the 4D Energy-Momentum Tensor of EFE.

Einstein Field Equations (EFE)

Finally, Einstein identified its tensor that describes the curvature of spacetime to the Energy-Momentum tensor that describes the characteristics of the object which curves spacetime. He added an empirical coefficient, $8\pi G/c^4$, to homogeize the right hand side of the EFE. Today, this coefficient is calculated to get back the Newton's Law from EFE in the case of a static sphere in a weak field. If no matter is present, the energy-momentum tensor vanishes, and we come back to a flat spacetime without gravitational field.

The Proposed Theory

However, some unsolved questions exist in the EFE, despite the fact that they work perfectly. For example, Einstein built the EFE without knowing 1/ what is mass, 2/ the mechanism of gravity, 3/ the mechanism by which spacetime is curved by mass ... To date, these enigmas remain. Considering that "mass curves spacetime" does not explain anything. No one knows by which strange phenomenon a mass can curve spacetime. It seems obvious that *if a process makes a deformation of spacetime, it may reasonably be expected to provide information about the nature of this phenomenon*. Therefore, the main purpose of the present paper is to try to solve these enigmas, i.e. to give a rational explanation of mass, gravity and spacetime curvature. The different steps to achieve this goal are:

- Special Relativity (SR). This section gives an overview of SR.
- *Einstein Tensor.* Explains the construction of the Einstein Tensor.
- Energy-Momentum Tensor. Covers the calculus of this tensor.
- *Einstein Constant*. Explains the construction of the Einstein Constant.
- EFE. This section assembles the three precedent parts to build the EFE.
- **EFE Inconsistencies.** Shows and solves four inconsistencies of EFE.
- Metrics. Explains how to build special metrics using the new theory.

E-2 Special Relativity (Background)

Lorentz Factor

$$\beta = \frac{v}{c} \qquad (2)$$
$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}} = \frac{1}{\sqrt{1 - \beta^2}} \qquad (3)$$

Minkowski Metric with signature (-, +, +, +):

$$\eta_{\mu\nu} \equiv \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$
(4)
$$ds^{2} = -c^{2}dt^{2} + dx^{2} + dy^{2} + dz^{2} = \eta_{\mu\nu}dx^{\mu}dx^{\nu}$$
(5)

Minkowski Metric with signature (+, -, -, -):

$$\eta_{\mu\nu} \equiv \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$
(6)
$$ds^{2} = c^{2}dt^{2} - dx^{2} - dy^{2} - dz^{2} = \eta_{\mu\nu}dx^{\mu}dx^{\nu}$$
(7)

Time Dilatation

" τ " is the proper time.

$$ds^{2} = c^{2}dt^{2} - dx^{2} - dy^{2} - dz^{2} = c^{2}d\tau^{2}$$
 (8)

$$d\tau^{2} = \frac{ds^{2}}{c^{2}} \quad \Rightarrow \quad d\tau^{2} = dt^{2} \left(1 - \frac{dx^{2} + dy^{2} + dz^{2}}{c^{2}dt^{2}} \right)$$
(9)
$$v^{2} = \frac{dx^{2} + dy^{2} + dz^{2}}{dt^{2}}$$
(10)
$$d\tau = dt \sqrt{1 - \beta^{2}}$$
(11)

Lenght Contractions

"dx'" is the proper lenght.

$$dx' = \frac{dx}{\sqrt{1 - \beta^2}} \qquad (12)$$

Lorentz Transformation

$$ct' = \gamma(ct - \beta x)$$
(13)

$$x' = \gamma(x - \beta ct)$$

$$y' = y$$

$$z' = z$$

$$\begin{bmatrix} ct' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} ct \\ x \\ y \\ z \end{bmatrix}$$
(14)

Inverse transformation on the x-direction:

$$\begin{aligned} ct &= \gamma(ct' + \beta x') & (15) \\ x &= \gamma(x' + \beta ct') \\ y &= y' \\ z &= z' \end{aligned}$$

$$\begin{aligned} \begin{bmatrix} c t \\ x \\ y \\ z \end{bmatrix} &= \begin{bmatrix} \gamma & +\beta\gamma & 0 & 0 \\ +\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} c t' \\ x' \\ y' \\ z' \end{bmatrix}$$

$$(16)$$

Four-Position

Event in a Minkowski space:

$$\mathbf{X} = x_{\mu} = (x_0, x_1, x_2, x_3) = (ct, x, y, z)$$
(17)

Displacement:

$$\Delta X_{\mu} = (\Delta x_0, \Delta x_1, \Delta x_2, \Delta x_3) = (c\Delta t, \Delta x, \Delta y, \Delta z)$$
(18)

$$dx_{\mu} = (dx_0, dx_1, dx_2, dx_3) = (cdt, dx, dy, dz)$$
(19)

Four-Velocity

 v_x, v_y, v_z = Traditional speed in 3D.

$$\mathbf{U} = u_{\mu} = (u_0, u_1, u_2, u_3) = \left(c\frac{dt}{d\tau}, \frac{dx}{d\tau}, \frac{dy}{d\tau}, \frac{dz}{d\tau}\right)$$
(20)

from (7):

$$ds^{2} = c^{2}dt^{2} \left(1 - \frac{dx^{2} + dy^{2} + dz^{2}}{c^{2}dt^{2}}\right)$$
(21)
$$v^{2} = \frac{dx^{2} + dy^{2} + dz^{2}}{dt^{2}}$$
(22)

$$ds^{2} = c^{2} dt^{2} \left(1 - \frac{v^{2}}{c^{2}} \right) \quad \Rightarrow \quad ds = c dt \sqrt{1 - \beta^{2}}$$
(23)

Condensed form:

$$u_{\mu} = \frac{dx_{\mu}}{d\tau} = \frac{dx_{\mu}}{dt}\frac{dt}{d\tau} \qquad (24)$$

Thus:

$$u_{0} = \frac{dx_{0}}{d\tau} = \frac{cdt}{dt \sqrt{1 - \beta^{2}}} = \frac{c}{\sqrt{1 - \beta^{2}}} = \gamma c \qquad (25)$$

$$u_{1} = \frac{dx_{1}}{d\tau} = \frac{dx}{dt \sqrt{1 - \beta^{2}}} = \frac{v_{x}}{\sqrt{1 - \beta^{2}}} = \gamma v_{x} \qquad (26)$$

$$u_{2} = \frac{dx_{2}}{d\tau} = \frac{dy}{dt \sqrt{1-\beta^{2}}} = \frac{v_{y}}{\sqrt{1-\beta^{2}}} = \gamma v_{y} \qquad (27)$$

$$u_{3} = \frac{dx_{3}}{d\tau} = \frac{dz}{dt \sqrt{1 - \beta^{2}}} = \frac{v_{z}}{\sqrt{1 - \beta^{2}}} = \gamma v_{z} \qquad (28)$$

Four-Acceleration

$$a_i = \frac{du_i}{d\tau} = \frac{d^2x_i}{d\tau^2} \qquad (29)$$

Four-Momentum

 p_x, p_y, p_z = Traditional momentum in 3D U = Four-velocity

$$\mathbf{P} = m\mathbf{U} = m(u_0, u_1, u_2, u_3) = \gamma(mc, p_x, p_y, p_z)$$
(30)

(p = mc = E/c)

$$\|\mathbf{P}\|^{2} = \frac{E^{2}}{c^{2}} - |\vec{p}|^{2} = m^{2}c^{2} \qquad (31)$$
$$p_{\mu} = mu_{\mu} = m\frac{dx_{\mu}}{d\tau} \qquad (32)$$

hence

$$p_{0} = mc \frac{dt}{d\tau} = mu_{0} = \gamma mc = \gamma E/c \quad (33)$$

$$p_{1} = m \frac{dx}{d\tau} = mu_{1} = \gamma mv_{x} \quad (34)$$

$$p_{2} = m \frac{dy}{d\tau} = mu_{2} = \gamma mv_{y} \quad (35)$$

$$p_{3} = m \frac{dz}{d\tau} = mu_{3} = \gamma mv_{z} \quad (36)$$

Force

$$\mathbf{F} = F_{\mu} = \frac{dp_{\mu}}{d\tau} = \left(\frac{d(mu_0)}{d\tau}, \frac{d(mu_1)}{d\tau}, \frac{d(mu_2)}{d\tau}, \frac{d(mu_3)}{d\tau}\right)$$
(37)

E-3 Einstein Tensor

Since the Einstein Tensor is not affected by the presented theory, one could think that it is not useful to study it in the framework of this document. However, the knowledge of the construction of the Einstein Tensor is necessary to fully understand the four inconsistencies highlighted and solved in this document. Therefore, this section is only a summary of the Einstein Tensor. A more accurate development of EFE can be obtained on books or on the Internet.

The Gauss Coordinates

Consider a curvilinear surface with coordinates u and v (Fig. E-1A). The distance between two points, M(u, v) and M'(u + du, v + dv), has been calculated by Gauss. Using the g_{ij} coefficients, this distance is:

$$ds^{2} = g_{11}du^{2} + g_{12}dudv + g_{21}dvdu + g_{22}dv^{2}$$
(38)

The Euclidean space is a particular case of the Gauss Coordinates that reproduces the Pythagorean theorem (Fig. E-1B). In this case, the Gauss coefficients are $g_{11} = 1$, $g_{12} = g_{21} = 0$, and $g_{22} = 1$.

$$ds^2 = du^2 + dv^2 \qquad (39)$$

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Figure E-1: Gauss coordinates in a curvilinear space (A) and in an Euclidean space (B).

Equation (39) may be condensed using the Kronecker Symbol δ , which is 0 for $i \neq j$ and 1 for i = j, and replacing du and dv by du_1 and du_2 . For indexes i, j = 0 and 1, we have:

$$ds^2 = \delta_{ij} du_i du_j \qquad (40)$$

The Metric Tensor

Generalizing the Gaussian Coordinates to "n" dimensions, equation (38) can be rewritten as:

$$ds^2 = g_{\mu\nu} du_\mu du_\nu \qquad (41)$$

or, with indexes μ and ν that run from 1 to 3 (example of x, y and z coordinates):

$$ds^{2} = g_{11}du_{1}^{2} + g_{12}du_{1}du_{2} + g_{21}du_{2}du_{1}\dots + g_{32}du_{3}du_{2} + g_{33}du_{3}^{2}$$
(42)

This expression is often called the "Metric" and the associated tensor, $g_{\mu\nu}$, the "Metric Tensor". In the spacetime manifold of RG, μ and ν are indexes which run from 0 to 3 (t, x, y and z). Each component can be viewed as a multiplication factor which must be placed in front of the differential displacements. Therefore, the matrix of coefficients $g_{\mu\nu}$ are a tensor 4×4 , i.e. a set of 16 real-valued functions defined at all points of the spacetime manifold.

$$g_{\mu\nu} = \begin{bmatrix} g_{00} & g_{01} & g_{02} & g_{03} \\ g_{10} & g_{11} & g_{12} & g_{13} \\ g_{20} & g_{21} & g_{22} & g_{23} \\ g_{30} & g_{31} & g_{32} & g_{33} \end{bmatrix}$$
(43)

However, in order for the metric to be symmetric, we must have:

$$g_{\mu\nu} = g_{\nu\mu} \qquad (44)$$

...which reduces to 10 independent coefficients, 4 for the diagonal in bold face in equation (45), g_{00} , g_{11} , g_{22} , g_{33} , and 6 for the half part above - or under - the diagonal, i.e. $g_{01} = g_{10}$, $g_{02} = g_{20}$, $g_{03} = g_{30}$, $g_{12} = g_{21}$, $g_{13} = g_{31}$, $g_{23} = g_{32}$. This gives:

$$g_{\mu\nu} = \begin{bmatrix} \mathbf{g_{00}} & g_{01} & g_{02} & g_{03} \\ (g_{10} = g_{01}) & \mathbf{g_{11}} & g_{12} & g_{13} \\ (g_{20} = g_{02}) & (g_{21} = g_{12}) & \mathbf{g_{22}} & g_{23} \\ (g_{30} = g_{03}) & (g_{31} = g_{13}) & (g_{32} = g_{23}) & \mathbf{g_{33}} \end{bmatrix}$$
(45)

To summarize, the metric tensor $g_{\mu\nu}$ in equations (43) and (45) is a matrix of functions which tells how to compute the distance between any two points in a given space. The metric components obviously depend on the chosen local coordinate system.

The Riemann Curvature Tensor

The Riemann curvature tensor $R^{\alpha}_{\beta\gamma\delta}$ is a four-index tensor. It is the most standard way to express curvature of Riemann manifolds. In spacetime, a 2-index tensor is associated to each point of a 2-index Riemannian manifold. For example, the Riemann curvature tensor represents the force experienced by a rigid body moving along a geodesic.

The Riemann tensor is the only tensor that can be constructed from the metric tensor and its first and second derivatives. These derivatives must exist if we are in a Riemann manifold. They are also necessary to keep homogeneity with the right hand side of EFE which can have first derivative such as the velocity dx/dt, or second derivative such as an acceleration d^2x/dt^2 .

Christoffel Symbols

The Christoffel symbols are tensor-like objects derived from a Riemannian metric $g_{\mu\nu}$. They are used to study the geometry of the metric. There are two closely related kinds of Christoffel symbols, the first kind Γ_{ijk} , and the second kind Γ_{ij}^k , also known as "affine connections" or "connection coefficients".

At each point of the underlying n-dimensional manifold, the Christoffel symbols are numerical arrays of real numbers that describe, in coordinates, the effects of parallel transport in curved surfaces and, more generally, manifolds. The Christoffel symbols may be used for performing practical calculations in differential geometry. In particular, the Christoffel symbols are used in the construction of the Riemann Curvature Tensor.

In many practical problems, most components of the Christoffel symbols are equal to zero, provided the coordinate system and the metric tensor possesses some common symmetries.

Comma Derivative

The following convention is often used in the writing of Christoffel Symbols. The components of the gradient dA are denoted $A_{,k}$ (a comma is placed before the index) and are given by:

$$A_{,k} = \frac{\partial A}{\partial x^k} \qquad (46)$$

Christoffel Symbols in spherical coordinates

The best way to understand the Christoffel symbols is to start with an example. Let's consider vectorial space \mathbb{E}_3 associated to a punctual space in spherical coordinates \mathcal{E}_3 . A vector **OM** in a fixed Cartesian coordinate system $(0, e_i^0)$ is defined as:

$$\mathbf{OM} = x^i e_i^0 \qquad (47)$$

or

$$\mathbf{OM} = r \, \sin\theta \, \cos\varphi \, e_1^0 + r \, \sin\theta \, \sin\varphi \, e_2^0 + r \, \cos\theta \, e_3^0 \qquad (48)$$

Calling e_k the evolution of **OM**, we can write:

$$e_k = \partial_k (x^i e_i^0) \qquad (49)$$

We can calculate the evolution of each vector e_k . For example, the vector e_1 (equation 50) is simply the partial derivative regarding r of equation (48). It means that the vector e_1 will be supported by a line OM oriented from zero to infinity. We can calculate the partial derivatives for θ and φ by the same manner. This gives for the three vectors e_1 , e_2 and e_3 :

$$e_{1} = \partial_{1}\mathbf{M} = \sin\theta \,\cos\varphi \,e_{1}^{0} + \sin\theta \,\sin\varphi \,e_{2}^{0} + \cos\theta \,e_{3}^{0}$$
(50)

$$e_{2} = \partial_{2}\mathbf{M} = r \,\cos\theta \,\cos\varphi \,e_{1}^{0} + r \,\cos\theta \,\sin\varphi \,e_{2}^{0} - r \,\sin\theta \,e_{3}^{0}$$
(51)

$$e_{3} = \partial_{3}\mathbf{M} = -r \,\sin\theta \,\sin\varphi \,e_{1}^{0} + r \,\sin\theta \,\cos\varphi \,e_{2}^{0}$$
(52)

The vectors e_1^0 , e_2^0 and e_3^0 are constant in module and direction. Therefore the differential of vectors e_1 , e_2 and e_3 are:

$$de_1 = (\cos\theta \, \cos\varphi \, e_1^0 + \cos\theta \, \sin\varphi \, e_2^0 - \sin\theta \, e_3^0) d\theta \dots \\ \dots + (-\sin\theta \, \sin\varphi \, e_1^0 + \sin\theta \, \cos\varphi \, e_2^0) d\varphi$$
(53)

$$de_{2} = (-r \sin\theta \cos\varphi \ e_{1}^{0} - r \sin\theta \sin\varphi \ e_{2}^{0} - r \cos\theta \ e_{3}^{0})d\theta \dots$$

$$\cdots + (-r \cos\theta \sin\varphi \ e_{1}^{0} + r \cos\theta \cos\varphi \ e_{2}^{0})d\varphi \dots$$

$$\cdots + (\cos\theta \cos\varphi \ e_{1}^{0} + \cos\theta \sin\varphi \ e_{2}^{0} - \sin\theta \ e_{3}^{0})dr \qquad (54)$$

$$de_{3} = (-r \cos\theta \sin\varphi \ e_{1}^{0} + r \cos\theta \cos\varphi \ e_{2}^{0})d\theta \dots$$

$$\cdots + (-r \sin\theta \cos\varphi \ e_{1}^{0} - r \sin\theta \sin\varphi \ e_{2}^{0})d\varphi \dots$$

$$\cdots + (-sin\theta \sin\varphi \ e_{1}^{0} + sin\theta \cos\varphi \ e_{2}^{0})dr \qquad (55)$$

We can remark that the terms in parenthesis are nothing but vectors e_1/r , e_2/r and e_3/r . This gives, after simplifications:

$$de_1 = (d\theta/r)e_2 + (d\varphi/r)e_3 \tag{56}$$

$$de_2 = (-r \ d\theta)e_1 + (dr/r)e_2 + (cotang\theta \ d\varphi)e_3$$
(57)

$$de_3 = (-r \sin^2\theta \ d\varphi)e_1 + (-\sin\theta \ \cos\theta \ d\varphi)e_2 + ((dr/r) + \cot ang\theta \ d\theta)e_3 \quad (58)$$

In a general manner, we can simplify the writing of this set of equation writing ω_i^j the contravariant components vectors de_i . The development of each term is given in the next section. The general expression, in 3D or more, is:

$$de_i = \omega_i^j e_j \tag{59}$$

Christoffel Symbols of the second kind

If we replace the variables r, θ and φ by u^1 , u^2 , and u^3 as follows:

$$u^1 = r; \quad u^2 = \theta; \quad u^3 = \varphi \tag{60}$$

... the differentials of the coordinates are:

$$du^1 = dr; \quad du^2 = d\theta; \quad du^3 = d\varphi$$
 (61)

... and the ω_i^j components become, using the Christoffel symbol Γ_{ki}^j :

$$\omega_i^j = \Gamma_{ki}^j \, du^k \tag{62}$$
In the case of our example, quantities Γ_{ki}^{j} are functions of r, θ and φ . These functions can be explicitly obtained by an identification of each component of ω_{i}^{j} with Γ_{ki}^{j} . The full development of the precedent expressions of our example is detailed as follows:

$$\begin{cases} \omega_1^1 = 0 \\ \omega_1^2 = 1/r \ d\theta \\ \omega_1^3 = 1/r \ d\varphi \\ \omega_2^1 = -r \ d\theta \\ \omega_2^2 = 1/r \ dr \\ \omega_2^3 = \cot ang \ \theta \ d\varphi \\ \omega_3^3 = -r \ sin^2\theta \ d\theta \\ \omega_3^3 = 1/r \ dr + \cot ang\theta \ d\theta \end{cases}$$
(63)

Replacing dr by du^1 , $d\theta$ by du^2 , and $d\varphi$ by du^3 as indicated in equation (61), gives:

$$\begin{cases} \omega_{1}^{1} = 0 \\ \omega_{1}^{2} = 1/r \ du^{2} \\ \omega_{1}^{3} = 1/r \ du^{3} \\ \omega_{2}^{1} = -r \ du^{2} \\ \omega_{2}^{2} = 1/r \ du^{1} \\ \omega_{2}^{3} = \cot ang \ \theta \ du^{3} \\ \omega_{3}^{1} = -r \ sin^{2}\theta \ du^{2} \\ \omega_{3}^{2} = -sin\theta \ cos\theta \ du^{3} \\ \omega_{3}^{3} = 1/r \ du^{1} + \cot ang \ \theta \ du^{2} \end{cases}$$
(64)

On the other hand, the development of Christoffel symbols are:

$$\begin{aligned}
\omega_{1}^{1} &= \Gamma_{11}^{1} \, du^{1} + \Gamma_{21}^{1} \, du^{2} + \Gamma_{31}^{1} \, du^{3} \\
\omega_{1}^{2} &= \Gamma_{21}^{2} \, du^{1} + \Gamma_{21}^{2} \, du^{2} + \Gamma_{31}^{2} \, du^{3} \\
\omega_{1}^{3} &= \Gamma_{11}^{3} \, du^{1} + \Gamma_{21}^{3} \, du^{2} + \Gamma_{31}^{3} \, du^{3} \\
\omega_{2}^{1} &= \Gamma_{12}^{1} \, du^{1} + \Gamma_{22}^{1} \, du^{2} + \Gamma_{32}^{1} \, du^{3} \\
\omega_{2}^{2} &= \Gamma_{22}^{2} \, du^{1} + \Gamma_{22}^{2} \, du^{2} + \Gamma_{32}^{2} \, du^{3} \\
\omega_{3}^{3} &= \Gamma_{13}^{3} \, du^{1} + \Gamma_{23}^{3} \, du^{2} + \Gamma_{33}^{3} \, du^{3} \\
\omega_{3}^{3} &= \Gamma_{13}^{2} \, du^{1} + \Gamma_{23}^{2} \, du^{2} + \Gamma_{33}^{3} \, du^{3} \\
\omega_{3}^{3} &= \Gamma_{13}^{3} \, du^{1} + \Gamma_{23}^{2} \, du^{2} + \Gamma_{33}^{3} \, du^{3} \\
\omega_{3}^{3} &= \Gamma_{13}^{3} \, du^{1} + \Gamma_{23}^{2} \, du^{2} + \Gamma_{33}^{3} \, du^{3}
\end{aligned}$$

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Finally, identifying the two equations array (64) and (65) gives the 27 Christoffel Symbols.

$\Gamma^{1}_{11} = 0$	$\Gamma^{1}_{21} = 0$	$\Gamma^{1}_{31} = 0$	
$\Gamma_{11}^2 = 0$	$\Gamma_{21}^2 = 1/r$	$\Gamma_{31}^2 = 0$	
$\Gamma^{3}_{11} = 0$	$\Gamma^{3}_{21} = 0$	$\Gamma^{3}_{31} = 1/r$	
$\Gamma^1_{12} = 0$	$\Gamma^1_{22} = -r$	$\Gamma^1_{32} = 0$	
$\Gamma_{12}^2 = 0$	$\Gamma_{22}^2 = 0$	$\Gamma_{32}^2 = 0$	(66)
$\Gamma_{12}^3 = 0$	$\Gamma_{22}^3 = 0$	$\Gamma_{32}^3 = cotang \ \theta$	
$\Gamma^1_{13} = 0$	$\Gamma^1_{23} = -r \ sin^2 \theta$	$\Gamma^{1}_{33} = 0$	
$\Gamma_{13}^2 = 0$	$\Gamma_{23}^2 = 0$	$\Gamma_{33}^2 = -\sin\theta \cos\theta$	
$\Gamma_{13}^3 = 1/r$	$\Gamma^3_{23} = cotang \ \theta$	$\Gamma^3_{33} = 0$	

These quantities Γ_{ki}^{j} are the Christoffel Symbols of the second kind. Identifying equations (59) and (62) gives the general expression of the Christoffel Symbols of the second kind:

$$de_i = \omega_i^j e_j = \Gamma_{ki}^j \ du^k \ e_j \tag{67}$$

Christoffel Symbols of the first kind

We have seen in the precedent example that we can directly get the quantities Γ_{ki}^{j} by identification. These quantities can also be obtained from the components g_{ij} of the metric tensor. This calculus leads to another kind of Christoffel Symbols.

Lets write the covariant components, noted ω_{ii} , of the differentials de_i :

$$\omega_{ji} = e_j \ de_i \qquad (68)$$

The covariant components ω_{ji} are also linear combinations of differentials du^i that can be written as follows, using the Christoffel Symbol of the first kind Γ_{kji} :

$$\omega_{ji} = \Gamma_{kji} \ du^k \qquad (69)$$

On the other hand, we know the basic relation:

$$\omega_{ji} = g_{jl} \; \omega_i^l \tag{70}$$

Porting equation (69) in equation (70) gives:

$$\Gamma_{kji} \ du^k = g_{jl} \ \omega_i^l \tag{71}$$

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Let's change the name of index j to l of equation (62):

$$\omega_i^l = \Gamma_{ki}^l \ du^k \tag{72}$$

Porting equation (117) in equation (116) gives the calculus of the Christoffel Symbols of the first kind from the Christoffel Symbols of the second kind:

$$\Gamma_{kji} \ du^k = g_{jl} \ \Gamma^j_{ki} \ du^k \tag{73}$$

Geodesic Equations

Let's take a curve M_0 -C- M_1 . If the parametric equations of the curvilinear abscissa are $u^i(s)$, the length of the curve will be:

$$l = \int_{M_0}^{M_1} \left(g_{ij} \frac{du^i}{ds} \frac{du^j}{ds} \right)^{1/2} ds \qquad (74)$$

If we pose $u'^i = du^i/ds$ and $u'^j = du^j/ds$ we get:

$$l = \int_{M_0}^{M_1} \left(g_{ij} u'^i u'^j \right)^{1/2} ds \qquad (75)$$

Here, the u'^{j} are the direction cosines of the unit vector supported by the tangent to the curve. Thus, we can pose:

$$f(u^k, u'^j) = g_{ij}u'^i u'^j = 1$$
(76)

The length l of the curve defined by equation (74) has a minimum and a maximum that can be calculated by the Euler-Lagrange Equation which is:

$$\frac{\partial \mathcal{L}}{\partial f_i} - \frac{d}{dx} \left(\frac{\partial \mathcal{L}}{\partial f'_i} \right) = 0 \qquad (77)$$

In the case of equation (74), the Euler-Lagrange equation gives:

$$\frac{d}{ds}(g_{ij}u'^{j}) - \frac{1}{2}\partial_i g_{jk} \ u'^{j} \ u'^{k} = 0 \qquad (78)$$

or

$$g_{ij}u'^{j} + (\partial_k g_{ij} - \frac{1}{2}\partial_i g_{jk}) \ u'^{j} \ u'^{k} = 0$$
 (79)

After developing derivative and using the Christoffel Symbol of the first kind, we get:

$$g_{ij}\frac{du'^{j}}{ds} + \Gamma_{jik} \ u'^{j} \ u'^{k} = 0 \qquad (80)$$

The contracted multiplication of equation (79) by g^{il} gives, with $g_{ij}g^{il}$ and $g^{il}\Gamma_{ijk} = \Gamma^l_{jk}$:

$$\frac{d^2u^l}{ds^2} + \Gamma^l_{jk}\frac{du^j}{ds}\frac{du^k}{ds} = 0 \qquad (81)$$

Parallel Transport

Figure E-2 shows two points M and M' infinitely close to each other in polar coordinates.



Figure E-2: Parallel transport of a vector.

In polar coordinates, the vector V_1 will become V_2 . To calculate the difference between two vectors V_1 and V_2 , we must before make a parallel transport of the vector V_2 from point M' to point M. This gives the vector V_3 . The absolute differential is defined by:

$$dV = V_3 - V_1$$
 (82)

Variations along a Geodesic

The variation along a 4D geodesic follows the same principle. For any curvilinear coordinates system y^i , we have (from equation 81):

$$\frac{d^2y^i}{ds^2} + \Gamma^i_{kj}\frac{dy^j}{ds}\frac{dy^k}{ds} = 0 \qquad (83)$$

where s is the abscissa of any point of the straight line from an origin such as M in figure E-2.

Let's consider now a vector \vec{v} having covariant components v_i . We can calculate the scalar product of \vec{v} and $\vec{n} = dy^k/ds$ as follows:

$$\vec{v} \cdot \vec{n} = v_i \frac{dy^i}{ds} \qquad (84)$$

During a displacement from M to M' (figure E-2), the scalar is subjected to a variation of:

$$d\left(v_i\frac{dy^i}{ds}\right) = dv_k\frac{dy^k}{ds} + v_i \ d\left(\frac{d^2y^i}{ds^2}\right) \tag{85}$$

or:

$$d\left(v_i\frac{dy^i}{ds}\right) = dv_k\frac{dy^k}{ds} + v_i\frac{d^2y^i}{ds^2}ds \qquad (86)$$

On one hand, the differential dv_k can be written as:

$$dv_k = \partial_j v_k \frac{dy^j}{ds} ds \qquad (87)$$

On the other hand, the second derivative can be extracted from equation (83) as follows:

$$\frac{d^2y^i}{ds^2} = -\Gamma^i_{kj}\frac{dy^j}{ds}\frac{dy^k}{ds} \qquad (88)$$

Porting equations (87) and (88) in (86) gives:

$$d\left(v_i\frac{dy^i}{ds}\right) = \partial_j v_k \frac{dy^j}{ds}\frac{dy^k}{ds}ds - v_i\Gamma^i_{kj}\frac{dy^j}{ds}\frac{dy^k}{ds}ds \qquad (89)$$

or:

$$d\left(v_i\frac{dy^i}{ds}\right) = \left(\partial_j v_k - v_i\Gamma^i_{kj}\right)\frac{dy^j}{ds}\frac{dy^k}{ds}ds \qquad (90)$$

Since $(dy^j/ds)ds = dy^j$, this expression can also be written as follows:

$$d(\vec{v}\cdot\vec{n}) = (\partial_j v_k - v_i \Gamma^i_{kj}) \ dy^j \frac{dy^k}{ds} \qquad (91)$$

The absolute differentials of the covariant components of vector \vec{v} are defined as:

$$Dv_k \frac{dy^k}{ds} = (\partial_j v_k - v_i \Gamma^i_{kj}) \ dy^j \qquad (92)$$

Finally, the quantity in parenthesis is called "affine connection" and is defined as follows:

$$\nabla_j v_k = \partial_j v_k - v_i \Gamma^i_{kj} \tag{93}$$

Some countries in the world use ";" for the covariant derivative and "," for the partial derivative. Using this convention, equation (93) can be written as:

$$v_{k;j} = v_{k,j} - v_i \Gamma_{kj}^i \qquad (94)$$

To summarize, given a function f, the covariant derivative $\nabla_v f$ coincides with the normal differentiation of a real function in the direction of the vector \vec{v} , usually denoted by $\vec{v}f$ and $df(\vec{v})$.

Second Covariant Derivatives of a Vector

Remembering that the derivative of the product of two functions is the sum of partial derivatives, we have:

$$\nabla_a(t_b r_c) = r_c \cdot \nabla_a t_b + t_b \cdot \nabla_a r_c \qquad (95)$$

Porting equation (93) in equation (95) gives:

$$\nabla_a(t_b r_c) = r_c \left(\partial_a t_b - t_l \Gamma^l_{ab} \right) + t_b \left(\partial_a r_c - r_l \Gamma^l_{ac} \right) \tag{96}$$

or:

$$\nabla_a(t_b r_c) = r_c \partial_a t_b - r_c t_l \Gamma^l_{ab} + t_b \partial_a r_c - t_b r_l \Gamma^l_{ac} \qquad (97)$$

Hence:

$$\nabla_a(t_b r_c) = r_c \partial_a t_b + t_b \partial_a r_c - r_c t_l \Gamma^l_{ab} - t_b r_l \Gamma^l_{ac} \qquad (98)$$

Finally:

$$\nabla_a(t_b r_c) = \partial_a(t_b r_c) - r_c t_l \Gamma^l_{ab} - t_b r_l \Gamma^l_{ac} \qquad (99)$$

Posing $t_b r_c = \nabla_j v_i$ gives:

$$\nabla_k (\nabla_j v_i) = \partial_k (\nabla_j v_i) - (\nabla_j v_r) \Gamma_{ik}^r - (\nabla_r v_i) \Gamma_{jk}^r$$
(100)

Porting equation (93) in equation (100) gives:

$$\nabla_k(\nabla_j v_i) = \partial_k(\partial_j v_i - v_l \Gamma_{ji}^l) - (\partial_j v_r - v_l \Gamma_{jr}^l) \Gamma_{ik}^r - (\partial_r v_i - v_l \Gamma_{ri}^l) \Gamma_{jk}^r$$
(101)

Hence:

$$\nabla_k (\nabla_j v_i) = \partial_{kj} v_i - (\partial_k \Gamma^l_{ji}) v_l - \Gamma^l_{ji} \partial_k v_l - \Gamma^r_{ik} \partial_j v_r + \Gamma^r_{ik} \Gamma^l_{jr} v_l - \Gamma^r_{jk} \partial_r v_i + \Gamma^r_{jk} \Gamma^l_{ri} v_l$$
(102)

The Riemann-Christoffel Tensor

In expression (102), if we make a swapping between the indexes j and k in order to get a differential on another way (i.e. a parallel transport) we get:

$$\nabla_j (\nabla_k v_i) = \partial_{jk} v_i - (\partial_j \Gamma^l_{ki}) v_l - \Gamma^l_{ik} \partial_j v_l - \Gamma^r_{ij} \partial_k v_r + \Gamma^r_{ij} \Gamma^l_{kr} v_l - \Gamma^r_{kj} \partial_r v_i + \Gamma^r_{kj} \Gamma^l_{ri} v_l$$
(103)

A subtraction between expressions (102) and (103) gives, after a rearrangement of some terms:

$$\nabla_k (\nabla_j v_i) - \nabla_j (\nabla_k v_i) = (\partial_{kj} - \partial_{jk}) v_i + (\partial_j \Gamma_{ki}^l - \partial_k \Gamma_{ji}^l) v_l + (\Gamma_{ik}^l \partial_j - \Gamma_{ji}^l \partial_k) v_l \dots$$

$$\dots + (\Gamma_{ij}^r \partial_k - \Gamma_{ik}^r \partial_j) v_r + (\Gamma_{ik}^r \Gamma_{jr}^l - \Gamma_{ij}^r \Gamma_{kr}^l) v_l + (\Gamma_{kj}^r - \Gamma_{jk}^r) \partial_r v_i + (\Gamma_{jk}^r \Gamma_{ri}^l - \Gamma_{kj}^r \Gamma_{ri}^l) v_l \quad (104)$$

On the other hand, since we have:

$$\Gamma_{jk}^r = \Gamma_{kj}^r \qquad (105)$$

Some terms of equation (104) are canceled:

$$\partial_{kj} - \partial_{jk} = 0 \qquad (106)$$
$$\Gamma^{r}_{kj} - \Gamma^{r}_{jk} = 0 \qquad (107)$$
$$\Gamma^{r}_{jk}\Gamma^{l}_{ri} - \Gamma^{r}_{kj}\Gamma^{l}_{ri} = 0 \qquad (108)$$

And consequently:

$$\nabla_k (\nabla_j v_i) - \nabla_j (\nabla_k v_i) = (\partial_j \Gamma_{ki}^l - \partial_k \Gamma_{ji}^l) v_l + (\Gamma_{ik}^l \partial_j - \Gamma_{ji}^l \partial_k) v_l \dots$$
$$\dots + (\Gamma_{ij}^r \partial_k - \Gamma_{ik}^r \partial_j) v_r + (\Gamma_{ik}^r \Gamma_{jr}^l - \Gamma_{ij}^r \Gamma_{kr}^l) v_l \quad (109)$$

Since the parallel transport is done on small portions of geodesics infinitely close to each other, we can take the limit:

$$\partial_j v_l, \quad \partial_k v_l, \quad \partial_j v_r, \quad \partial_k v_r \to 0$$
 (110)

This means that the velocity field is considered equal in two points of two geodesics infinitely close to each other. Then we can write:

$$\nabla_k (\nabla_j v_i) - \nabla_j (\nabla_k v_i) \cong (\partial_j \Gamma_{ki}^l - \partial_k \Gamma_{ji}^l + \Gamma_{ik}^r \Gamma_{jr}^l - \Gamma_{ij}^r \Gamma_{kr}^l) v_l \quad (111)$$

As a result of the tensorial properties of covariant derivatives and of the components v_l , the quantity in parenthesis is a four-order tensor defined as:

$$R_{i,jk}^{l} = \partial_{j}\Gamma_{ki}^{l} - \partial_{k}\Gamma_{ji}^{l} + \Gamma_{ik}^{r}\Gamma_{jr}^{l} - \Gamma_{ij}^{r}\Gamma_{kr}^{l}$$
(112)

In this expression, the comma in Christoffel Symbols means a partial derivative. The tensor $R_{i,jk}^{l}$ is called *Riemann-Christoffel Tensor* or *Curvature Tensor* which characterizes the curvature of a Riemann Space.

The Ricci Tensor

The contraction of the Riemann-Christoffel Tensor $R_{i,jk}^l$ defined by equation (112) relative to indexes l and j leads to a new tensor:

$$R_{ik} = R_{i,lk}^l = \partial_l \Gamma_{ki}^l - \partial_k \Gamma_{li}^l + \Gamma_{ik}^r \Gamma_{lr}^l - \Gamma_{il}^r \Gamma_{kr}^l$$
(113)

This tensor R_{ik} is called the "Ricci Tensor". Its mixed components are given by:

$$R^{jk} = g^{ji}R_{ik} \quad (114)$$

The Scalar Curvature

The Scalar Curvature, also called the "Curvature Scalar" or "Ricci Scalar", is given by:

$$R = R_i^i = g^{ij} R_{ij} \tag{115}$$

The Bianchi Second Identities

The Riemann-Christoffel tensor verifies a particular differential identity called the "Bianchi Identity". This identity involves that the Einstein Tensor has a null divergence, which leads to a constraint. The goal is to reduce the degrees of freedom of the Einstein Equations. To calculate the second Bianchi Identities, we must derivate the Riemann-Christoffel Tensor defined in equation (112):

$$\nabla_t R_{i,rs}^l = \partial_{rt} \Gamma_{si}^l - \partial_{st} \Gamma_{ri}^l \qquad (116)$$

A circular permutation of indexes r, s and t gives:

$$\nabla_r R_{i,st}^l = \partial_{sr} \Gamma_{ti}^l - \partial_{tr} \Gamma_{si}^l \qquad (117)$$
$$\nabla_s R_{i,tr}^l = \partial_{ts} \Gamma_{ri}^l - \partial_{rs} \Gamma_{ti}^l \qquad (118)$$

Since the derivation order is interchangeable, adding equations (116), (117) and (118) gives:

$$\nabla_t R_{i,rs}^l + \nabla_r R_{i,st}^l + \nabla_s R_{i,tr}^l = 0 \qquad (119)$$

The Einstein Tensor

If we make a contraction of the second Bianchi Identities (equation 119) for t = l, we get:

$$\nabla_l R_{i,rs}^l + \nabla_r R_{i,sl}^l + \nabla_s R_{i,lr}^l = 0 \qquad (120)$$

Hence, taking into account and the definition of the Ricci Tensor of equation (113) and that $R_{i,sl}^l = -R_{i,ls}^l$, we get:

$$\nabla_l R_{i,rs}^l + \nabla_s R_{ir} - \nabla_r R_{is} = 0 \qquad (121)$$

The variance change with g_{ij} gives:

$$\nabla_s R_{ir} = \nabla_s (g_{ik} R_r^k) \qquad (122)$$

or

$$\nabla_s R_{ir} = g_{ik} \nabla_s R_r^k \qquad (123)$$

Multiplying equation (121) by g^{ik} gives:

$$g^{ik}\nabla_l R^l_{i,rs} + g^{ik}\nabla_s R_{ir} - g^{ik}\nabla_r R_{is} = 0 \qquad (124)$$

Using the property of equation (123), we finally get:

$$\nabla_l R^{kl}_{,rs} + \nabla_s R^k_r - \nabla_r R^k_s = 0 \qquad (125)$$

Let's make a contraction on indexes k and s:

$$\nabla_k R^{kl}_{,rk} + \nabla_k R^k_r - \nabla_r R^k_k = 0 \qquad (126)$$

The first term becomes:

$$\nabla_k R_r^k + \nabla_k R_r^k - \nabla_r R_k^k = 0 \qquad (127)$$

After a contraction of the third term we get:

$$2\nabla_k R_r^k - \nabla_r R = 0 \qquad (128)$$

Dividing this expression by two gives:

$$\nabla_k R_r^k - \frac{1}{2} \nabla_r R = 0 \qquad (129)$$

or:

$$\nabla_k \left(R_r^k - \frac{1}{2} \delta_r^k R \right) = 0 \qquad (130)$$

A new tensor may be written as follows:

$$G_r^k = R_r^k - \frac{1}{2}\delta_r^k R \qquad (131)$$

The covariant components of this tensor are:

$$G_{ij} = g_{ik} \ G_j^k \qquad (132)$$

or

$$G_{ij} = g_{ik} \left(R_j^k - \frac{1}{2} \delta_j^k R \right)$$
(133)

Finally get the Einstein Tensor which is defined by:

$$G_{ij} = R_{ij} - \frac{1}{2}g_{ij}R \tag{134}$$

The Einstein Constant

The Einstein Tensor G_{ij} of equation (134) must match the Energy-Momentum Tensor T_{uv} defined later. This can be done with a constant κ so that:

$$G_{\mu\nu} = \kappa T_{\mu\nu} \qquad (135)$$

This constant κ is called "Einstein Constant" or "Constant of Proportionality". To calculate it, the Einstein Equation (134) must be identified to the Poisson's classical field equation, which is the mathematical form of the Newton Law. So, the weak field approximation is used to calculate the Einstein Constant. Three criteria are used to get this "Newtonian Limit":

1 - The speed is low regarding that of the light c.

2 - The gravitational field is static.

3 - The gravitational field is weak and can be seen as a weak perturbation $h_{\mu\nu}$ added to a flat spacetime $\eta_{\mu\nu}$ as follows:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} \qquad (136)$$

We start with the equation of geodesics (83):

$$\frac{d^2x^{\mu}}{ds^2} + \Gamma^{\mu}_{\alpha\beta}\frac{dx^{\alpha}}{ds}\frac{dx^{\beta}}{ds} = 0 \qquad (137)$$

This equation can be simplified in accordance with the first condition:

$$\frac{d^2 x^{\mu}}{ds^2} + \Gamma^{\mu}_{00} \left(\frac{dx^0}{ds}\right)^2 = 0 \qquad (138)$$

The two other conditions lead to a simplification of Christoffel Symbols of the second kind as follows:

$$\Gamma^{\mu}_{00} = \frac{1}{2}g^{\mu\lambda}(\partial_0 g_{\lambda 0} + \partial_0 g_{0\lambda} - \partial_\lambda g_{00}) \qquad (139)$$

Or, considering the second condition:

$$\Gamma^{\mu}_{00} = -\frac{1}{2}g^{\mu\lambda}\partial_{\lambda} g_{00} \qquad (140)$$

And also considering the third condition:

$$\Gamma^{\mu}_{00} \approx -\frac{1}{2} (\eta^{\mu\lambda} + h^{\mu\lambda}) (\partial_{\lambda}\eta_{00} + \partial_{\lambda}h_{00}) \qquad (141)$$

In accordance with the third condition, the term $\partial_{\lambda}\eta_{00}$ is canceled since it is a flat space:

$$\Gamma^{\mu}_{00} \approx -\frac{1}{2} (\eta^{\mu\lambda} + h^{\mu\lambda}) \partial_{\lambda} h_{00} \qquad (142)$$

Another simplification due to the approximation gives:

$$\Gamma^{\mu}_{00} \approx -\frac{1}{2} \eta^{\mu\lambda} \partial_{\lambda} h_{00} \qquad (143)$$

The equation of geodesics then becomes:

$$\frac{d^2x^{\mu}}{dt^2} - \frac{1}{2}\eta^{\mu\lambda}\partial_{\lambda}h_{00}\left(\frac{dx^0}{dt}\right)^2 = 0 \qquad (144)$$

Reduced to the time component ($\mu = 0$), equation (144) becomes:

$$\frac{d^2x^{\mu}}{dt^2} - \frac{1}{2}\eta^{0\lambda}\partial_{\lambda}h_{00}\left(\frac{dx^0}{dt}\right)^2 = 0 \qquad (145)$$

The Minkowski Metric shows that $\eta_{0\lambda} = 0$ for $\lambda > 0$. On the other hand, a static metric (third condition) gives $\partial_0 h_{00} = 0$ for $\lambda = 0$. So, the 3x3 matrix leads to:

$$\frac{d^2x^i}{dt^2} - \frac{1}{2}\partial_i h_{00} \left(\frac{dx^0}{dt}\right)^2 = 0 \qquad (146)$$

Replacing dt by the proper time $d\tau$ gives:

$$\frac{d^2x^i}{d\tau^2} - \frac{1}{2}\partial_i h_{00} \left(\frac{dx^0}{d\tau}\right)^2 = 0 \qquad (147)$$

Dividing by $(dx_0/d\tau)^2$ leads to:

$$\frac{d^2x^i}{d\tau^2} \left(\frac{d\tau}{dx^0}\right)^2 = \frac{1}{2}\partial_i h_{00} \qquad (148)$$

$$\frac{d^2x^i}{(dx^0)^2} = \frac{1}{2}\partial_i h_{00} \qquad (149)$$

Replacing x^0 by ct gives:

$$\frac{d^2x^i}{d(ct)^2} = \frac{1}{2}\partial_i h_{00} \qquad (150)$$

or:

$$\frac{d^2x^i}{dt^2} = \frac{c^2}{2}\partial_i h_{00}$$
 (151)

Let us pose:

$$h_{00} = -\frac{2}{c^2}\Phi \qquad (152)$$

or

$$\Delta h_{00} = -\frac{2}{c^2} \Delta \Phi \qquad (153)$$

Since the approximation is in an euclidean space, the Laplace operator can be written as:

$$\Delta h_{00} = -\frac{2}{c^2} \nabla^2 \Phi \qquad (154)$$
$$\Delta h_{00} = -\frac{2}{c^2} (4\pi G_0 \rho) \qquad (155)$$
$$\Delta h_{00} = -\frac{8\pi G_0}{c^2} \rho \qquad (156)$$

On the other hand, the element T_{00} defined later is:

$$T_{00} = \rho c^2$$
 (157)

or

$$\rho = \frac{T_{00}}{c^2} \qquad (158)$$

Porting equation (158) in equation (156) gives:

$$\Delta h_{00} = -\frac{8\pi G_0}{c^2} \frac{T_{00}}{c^2} \qquad (159)$$

or:

$$\Delta h_{00} = -\frac{8\pi G_0}{c^4} T_{00} \qquad (160)$$

The left hand side of equation (160) is the perturbation part of the Einstein Tensor in the case of a static and weak field approximation. It directly gives a constant of proportionality which also verifies the homogeneity of the EFE (equation 1). This equation will be fully explained later in this document. Thus:

$$Einstein \ Constant = \frac{8\pi G_0}{c^4} \tag{161}$$

E-4 Movement Equations in a Newtonian Fluid

The figure E-3 (next page) shows an elementary parallelepiped of dimensions dx, dy, dz which is a part of a fluid in static equilibrium. This cube is generally subject to volume forces in all directions, as the Pascal Theorem states. The components of these forces are oriented in the three orthogonal axis. The six sides of the cube are: A-A', B-B' and C-C'.



Figure E-3: Elementary parallelepiped dx, dy, dz

E-5 Normal Constraints

On figure E-4 (next page), the normal constraints to each surface are noted " σ ". The tangential constraints to each surface are noted " τ ". Since we have six sides, we have six sets of equations. In the following equations, ' σ " and " τ ' are constraints (a constraint is a pressure), dF is an elementary force, and dS is an elementary surface:

$$\frac{d\vec{F}_{+x}}{dS} = \vec{\sigma}_x + \vec{\tau}'_{yx} + \vec{\tau}_{xz} \qquad (162)$$

$$\frac{dF_{-x}}{dS} = \vec{\sigma}'_x + \vec{\tau}'_{xy} + \vec{\tau}_{zx} \qquad (163)$$

$$\frac{dF_{+y}}{dS} = \vec{\sigma}_y + \vec{\tau}_{yz} + \vec{\tau}_{yx}^{\prime\prime} \qquad (164)$$

$$\frac{d\vec{F}_{-y}}{dS} = \vec{\sigma}'_y + \vec{\tau}_{zy} + \vec{\tau}''_{xy} \qquad (165)$$

$$\frac{dF_{+z}}{dS} = \vec{\sigma}_z + \vec{\tau}''_{xz} + \vec{\tau}'_{yz} \qquad (166)$$

$$\frac{d\vec{F}_{-z}}{dS} = \vec{\sigma}'_{z} + \vec{\tau}''_{zx} + \vec{\tau}'_{zy} \qquad (167)$$



Figure E-4: Forces on the elementary parallelepiped sides.

We can simplify these equations as follows:

$$\vec{\sigma}_x + \vec{\sigma}'_x = \vec{\sigma}_X \qquad (168)$$
$$\vec{\sigma}_y + \vec{\sigma}'_y = \vec{\sigma}_Y \qquad (169)$$
$$\vec{\sigma}_z + \vec{\sigma}'_z = \vec{\sigma}_Z \qquad (170)$$

So, only three components are used to define the normal constraint forces, i.e. one per axis.

E-6 Tangential Constraints

If we calculate the force's momentum regarding the gravity center of the parallelepiped, we have 12 tangential components (two per side). Since some forces are in opposition to each other, only 6 are sufficient to describe the system. Here, we calculate the three momenta for each plan, XOY, XOZ and YOZ, passing through the gravity center of the elementary parallelepiped:

For the XOY plan:

$$\mathcal{M}_{XOY} = (\vec{\tau}_{zy} dz dx) \frac{dy}{2} + (\vec{\tau}_{zx} dy dz) \frac{dx}{2} + (\vec{\tau}_{yz} dx dz) \frac{dy}{2} + (\vec{\tau}_{xz} dz dy) \frac{dx}{2}$$
(171)
$$\mathcal{M}_{XOY} = \frac{1}{2} \left[(\vec{\tau}_{zy} + \vec{\tau}_{yz}) + (\vec{\tau}_{zx} + \vec{\tau}_{xz}) \right]$$
(172)
$$\mathcal{M}_{XOY} = \frac{1}{2} dV \left[(\vec{\tau}_{ZY} + \vec{\tau}_{ZX}) \right]$$
(173)
$$\mathcal{M}_{XOY} = \frac{1}{2} dV \vec{\tau}_{XOY}$$
(174)

For the XOZ plan:

$$\mathcal{M}_{XOZ} = (\vec{\tau}'_{yx}dydz)\frac{dx}{2} + (\vec{\tau}'_{yz}dxdy)\frac{dz}{2} + (\vec{\tau}'_{xy}dzdy)\frac{dx}{2} + (\vec{\tau}'_{zy}dydx)\frac{dz}{2}$$
(175)

$$\mathcal{M}_{XOZ} = \frac{1}{2}dV \left[(\vec{\tau}'_{yx} + \vec{\tau}'_{xy}) + (\vec{\tau}'_{yz} + \vec{\tau}'_{zy}) \right]$$
(176)

$$\mathcal{M}_{XOZ} = \frac{1}{2}dV \left[(\vec{\tau}_{YX} + \vec{\tau}_{YZ}) \right]$$
(177)

$$\mathcal{M}_{XOZ} = \frac{1}{2}dV\vec{\tau}_{XOZ}$$
(178)

For the ZOY plan:

$$\mathcal{M}_{ZOY} = (\vec{\tau}_{xz}'' dx dy) \frac{dz}{2} + (\vec{\tau}_{yx}'' dx dz) \frac{dy}{2} + (\vec{\tau}_{zx}'' dy dx) \frac{dz}{2} + (\vec{\tau}_{xy}'' dz dx) \frac{dy}{2}$$
(179)
$$\mathcal{M}_{ZOY} = \frac{1}{2} dV \left[(\vec{\tau}_{xz}'' + \vec{\tau}_{zx}'') + (\vec{\tau}_{yx}'' + \vec{\tau}_{xy}'') \right]$$
(180)
$$\mathcal{M}_{ZOY} = \frac{1}{2} dV \left[(\vec{\tau}_{ZX} + \vec{\tau}_{XY}) \right]$$
(181)
$$\mathcal{M}_{ZOY} = \frac{1}{2} dV \vec{\tau}_{ZOY}$$
(182)

So, for each plan, only one component is necessary to define the set of momentum forces. Since the elementary volume dV is equal to a^*a^*a (a = dx, dy or dz), we can come back to the constraint equations dividing each result (174), (178) and (182) by dV/2.

E-7 Constraint Tensor

Finally, the normal and tangential constraints can be reduced to only 6 terms with $\vec{\tau}_{xy} = \vec{\tau}_{yx}$, $\vec{\tau}_{xz} = \vec{\tau}_{zx}$, and $\vec{\tau}_{zy} = \vec{\tau}_{yz}$:

$$\vec{\sigma}_X = \vec{\sigma}_{xx} \qquad (183)$$
$$\vec{\sigma}_Y = \vec{\sigma}_{yy} \qquad (184)$$
$$\vec{\sigma}_Z = \vec{\sigma}_{zz} \qquad (185)$$
$$\vec{\tau}_{XOY} = \vec{\tau}_{xy} \qquad (186)$$
$$\vec{\tau}_{XOZ} = \vec{\tau}_{xz} \qquad (187)$$
$$\vec{\tau}_{ZOY} = \vec{\tau}_{zy} \qquad (188)$$

Using a matrix representation, we get:

$$\begin{bmatrix} F_x \\ F_y \\ F_z \end{bmatrix} = \begin{bmatrix} \sigma_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \sigma_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \sigma_{zz} \end{bmatrix} \begin{bmatrix} S_x \\ S_y \\ S_z \end{bmatrix}$$
(189)

The constraint tensor at point M becomes:

$$T_{(M)} = \begin{bmatrix} \sigma_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \sigma_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & \tau_{12} & \tau_{13} \\ \tau_{21} & \sigma_{22} & \tau_{23} \\ \tau_{31} & \tau_{32} & \sigma_{33} \end{bmatrix}$$
(190)

This tensor is symmetric and its meaning is shown in Figure E-5.



Figure E-5: Meaning of the Constraint Tensor.

Since all the components of the tensor are pressures (more exactly constraints), we can represent it by the following equation where F_i are forces and s_i are surfaces:

$$F_i = \sigma_{ij} s_j = \sum_j \sigma_{ij} s_j \tag{191}$$

E-8 Energy-Momentum Tensor

We can write equation (191) as:

$$\sigma_{ij} = \frac{F_i}{s_j} = \frac{\Delta F_i}{\Delta s_j} = \frac{\Delta(ma_i)}{\Delta s_j} \qquad (192)$$

Here we suppose that only volumes (x, y and z) and time (t) make the force vary. Therefore, the mass (m) can be replaced by volume (V) using a constant density (ρ):

$$m = \rho V \qquad (193)$$

Porting (193) in (192) gives:

$$\sigma_{ij} = \frac{\Delta(m.a_i)}{\Delta s_j} = \frac{\Delta(\rho V.a_i)}{\Delta s_j} = \frac{\rho V}{\Delta s_j} \Delta a_i \qquad (194)$$

As shown in figure E-3, V and S concern an elementary parallelepiped.

$$V = (\Delta X_j)^3$$
 and $S_j = (\Delta X_j)^2$ (195)

Thus:

$$\frac{V}{S_j} = \frac{(\Delta X_j)^3}{(\Delta X_j)^2} = \Delta X_j \qquad (196)$$

Hence

$$\sigma_{ij} = \rho \ \Delta X_j \Delta a_i \qquad (197)$$

$$\sigma_{ij} = \rho \Delta X_j \frac{v_i}{\Delta t} = \rho \frac{\Delta X_j}{\Delta t} v_i \qquad (198)$$

Finally

$$\sigma_{ij} = \rho v_j v_i \qquad (199)$$

This tensor comes from the Fluid Mechanics and uses traditional variables v_x, v_y and v_z . We can extend these 3D variables to 4D in accordance with Special Relativity (see above). The new 4D tensor created, called the "Energy-Momentum Tensor', has the same properties as the old one, in particular symmetry.

To avoid confusion, let's replace σ_{ij} by $T_{\mu\nu}$ as follows:

$$T_{\mu\nu} = \rho u_{\mu} u_{\nu} \tag{200}$$

or

$$T_{\mu\nu} = \begin{bmatrix} T_{00} & T_{01} & T_{02} & T_{03} \\ T_{10} & T_{11} & T_{12} & T_{13} \\ T_{20} & T_{21} & T_{22} & T_{23} \\ T_{30} & T_{31} & T_{32} & T_{33} \end{bmatrix} = \begin{bmatrix} \rho u_0 u_0 & \rho u_0 u_1 & \rho u_0 u_2 & \rho u_0 u_3 \\ \rho u_1 u_0 & \rho u_1 u_1 & \rho u_1 u_2 & \rho u_1 u_3 \\ \rho u_2 u_0 & \rho u_2 u_1 & \rho u_2 u_2 & \rho u_2 u_3 \\ \rho u_3 u_0 & \rho u_3 u_1 & \rho u_3 u_2 & \rho u_3 u_3 \end{bmatrix}$$
(201)

...with u_{μ} and u_{ν} as defined in equations (25) to (28).

This tensor may be written in a more explicit form, using the "traditional" velocity v_x, v_y and v_z instead of the relativistic velocities, as shown in equations (25) to (28):

$$T_{\mu\nu} = \begin{bmatrix} \rho\gamma^2 c^2 & \rho\gamma^2 cv_x & \rho\gamma^2 cv_y & \rho\gamma^2 cv_z \\ \rho\gamma^2 cv_x & \rho\gamma^2 v_x v_x & \rho\gamma^2 v_x v_y & \rho\gamma^2 v_x v_z \\ \rho\gamma^2 cv_y & \rho\gamma^2 v_y v_x & \rho\gamma^2 v_y v_y & \rho\gamma^2 v_y v_z \\ \rho\gamma^2 cv_z & \rho\gamma^2 v_z v_x & \rho\gamma^2 v_z v_y & \rho\gamma^2 v_z v_z \end{bmatrix}$$
(202)

For low velocities, we have $\gamma = 1$:

$$T_{\mu\nu} = \begin{bmatrix} \rho c^2 & \rho c v_x & \rho c v_y & \rho c v_z \\ \rho c v_x & \rho v_x v_x & \rho v_x v_y & \rho v_x v_z \\ \rho c v_y & \rho v_y v_x & \rho v_y v_y & \rho v_y v_z \\ \rho c v_z & \rho v_z v_x & \rho v_z v_y & \rho v_z v_z \end{bmatrix}$$
(203)

Replacing the spatial part of this tensor by the old definitions (equation 189) gives:

$$T_{\mu\nu} = \begin{bmatrix} \rho c^2 & \rho c v_x & \rho c v_y & \rho c v_z \\ \rho c v_x & \sigma_x & \tau_{xy} & \tau_{xz} \\ \rho c v_y & \tau_{yx} & \sigma_y & \tau_{yz} \\ \rho c v_z & \tau_{zx} & \tau_{zy} & \sigma_z \end{bmatrix}$$
(204)

Replacing ρ by mc^2 (equation 193) and mc^2 by E leads to one of the most commons form of the Energy-Momentum Tensor, V being the volume:

$$T_{\mu\nu} = \begin{bmatrix} T_{00} & T_{01} & T_{02} & T_{03} \\ T_{10} & T_{11} & T_{12} & T_{13} \\ T_{20} & T_{21} & T_{22} & T_{23} \\ T_{30} & T_{31} & T_{32} & T_{33} \end{bmatrix} = \begin{bmatrix} E/V & \rho cv_x & \rho cv_y & \rho cv_z \\ \rho cv_x & \sigma_x & \tau_{xy} & \tau_{xz} \\ \rho cv_y & \tau_{yx} & \sigma_y & \tau_{yz} \\ \rho cv_z & \tau_{zx} & \tau_{zy} & \sigma_z \end{bmatrix}$$
(205)



Figure E-6: Meaning of the Energy-Momentum Tensor.

E-9 Einstein Field Equations

Finally, the association of the Einstein Tensor $G_{\mu\nu} = R_{\mu\nu} - 1/2g_{\mu\nu}R$ (equation 134), the Einstein Constant $8\pi G_0/c^4$ (equation 161), and the Energy-Momentum Tensor $T_{\mu\nu}$ (equations 204/205), gives the full Einstein Field Equations, excluding the cosmological constant Λ which is not proven:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G_0}{c^4}T_{\mu\nu}$$
 (206)

To summarize, the Einstein field equations are 16 nonlinear partial differential equations that describe the curvature of spacetime, i.e. the gravitational field, produced by a given mass. As a result of the symmetry of $G_{\mu\nu}$ and $T_{\mu\nu}$, the actual number of equations are reduced to 10, although there are an additional four differential identities (the Bianchi identities) satisfied by $G_{\mu\nu}$, one for each coordinate.

E-10 Dimensional Analysis

The dimensional analysis verifies the homogeneity of equations. The most common dimensional quantities used in this chapter are:

 $\begin{array}{l} \textbf{Speed V} \Rightarrow [L/T] \\ \textbf{Energy E} \Rightarrow [ML^2/T^2] \\ \textbf{Force F} \Rightarrow [ML/T^2] \\ \textbf{Pressure P} \Rightarrow [M/LT^2] \\ \textbf{Momentum M} \Rightarrow [ML/T] \\ \textbf{Gravitation constant G} \Rightarrow [L^3/MT^2] \end{array}$

Here are the dimensional analysis of the Energy-Momentum Tensor:

 \mathbf{T}_{00} is the energy density, i.e. the amount of energy stored in a given region of space per unit volume. The dimensional quantity of E is $[ML^2/T^2]$ and V is $[L^3]$. So, the dimensional quantity of E/V is $[M/LT^2]$. Energy density has the same physical units as pressure which is $[M/LT^2]$.

 $\mathbf{T_{01}}, \mathbf{T_{02}}, \mathbf{T_{03}}$ are the energy flux, i.e. the rate of transfer of energy through a surface. The quantity is defined in different ways, depending on the context. Here, ρ is the density [M/V], and c and v_i (i = 1 to 3) are velocities [L/T]. So, the dimensional quantity of T_{0i} is $[M/L^3][L/T][L/T]$. It is that of a pressure $[M/LT^2]$.

 $\mathbf{T_{10}}, \mathbf{T_{20}}, \mathbf{T_{30}}$ are the momentum density, which is the momentum per unit volume. The dimensional quantity of T_{i0} (i = 1 to 3) is identical to T_{0i} , i.e. that of a pressure $[M/LT^2]$.

 $T_{12}, T_{13}, T_{23}, T_{21}, T_{31}, T_{32}$ are the shear stress, or a pressure $[M/LT^2]$.

 T_{11}, T_{22}, T_{33} are the normal stress or isostatic pressure $[M/LT^2]$.

Note: The Momentum flux is the sum of the shear stresses and the normal stresses.

As we see, all the components of the Energy-Momentum Tensor have a pressure-like dimensional quantity:

$$T_{\mu\nu} \Rightarrow \text{Pressure } [M/LT^2]$$
 (207)

E-11 Inconsistencies of EFE

Here we have demonstrated that the Energy-Momentum Tensor is nothing but the extension in spacetime of the Stress Tensor of the Fluid Mechanics. We also have demonstrated that all elements of the Energy-Momentum Tensor have a pressure-like dimension. This is not abnormal since the Energy-Momentum Tensor is built on the Stress Tensor, which describes pressures. This is why the Energy-Momentum Tensor is often called "Stress-Energy-Momentum Tensor" and sometimes is also used in the Fluid Mechanics. The generalization of the Stress Tensor to the Energy-Momentum Tensor is shown in figure E-7 (next page). The related mathematics have been described in the precedent sections.

The Stress Tensor \mathbb{P} in Fluid Mechanics is a part of the NavierStokes equations that describe the motion of fluid substances such as liquids and gases: $\rho Dv/Dt = \nabla \cdot \mathbb{P} + \rho \mathbf{f}$. This tensor represents the external pressures, more exactly normal and tangential constraints, acting on the fluid, as shown in figure E-8 (next page). For pedagogical purposes, only the upper normal pressure is shown on this figure. Other pressures, as defined in figure E-4, are not represented.



Figure E-7: Stress Tensor vs. Energy-Momentum Tensor.



Figure E-8: External pressures on a volume.

QUESTION #1

Since the Energy-Momentum Tensor represents a set of pressures made by spacetime on objects, and since the normal component (i.e. gravity) is an isostatic pressure, why is gravitation considered an ATTRACTIVE force instead of a PRESSURE force?

QUESTION #2

Since the Fluid Mechanics exerts a pressure on the VOLUME of the object, why has EFE replaced the volume by MASS?

QUESTION #3

In GR, the curvature of spacetime is supposed to be concave, as shown by Figure E-9. This figure is taken from one of the thousands of graphics representing the curvature of spacetime. On the contrary, as shown on Figure E-8, the curvature of the fluid in the Fluid Mechanics is convex. So, why has GR replaced the CONVEX curvature by a CONCAVE curvature?



Figure E-9: Curvature of spacetime by a black hole.

QUESTION #4

What is mass? To date, no one knows. So, we are faced with another inconsistency due to an unknown variable, "m", present in the right hand side of EFE (equation 193) but not in the left hand side. This leads to the following alternative:

1/ The mass variable "m" is not defined in 4D. In this case, we need an additional dimension to define it. In other words, the 4D Energy-Momentum tensor must be converted in a 5D tensor including a new dimension, "m". The problem is that the expression of the left hand side of EFE is in 4D:

$$f_{(x,y,z,t)} = g_{(x,y,z,t,m)}$$
(208)

As we see, the homogeneity of the EFE is not respected since the number of dimensions of the left hand side is different to that of the right hand side.

2/ The mass variable "m" is fully defined in 4D. In this case, the EFE should explicitly include the nature of mass and its expression: $m = f_{(x,y,z,t)}$. This information is missing inside the EFE. In other words, the EFE uses a variable "m" without knowing its meaning nor its expression.

Summary

Since the Energy-Momentum Tensor is built from the Stress Tensor, it is obvious that these two tensors must share the same reasoning and principle of construction. As we see, this is not the case. Einstein and Grossman brought some modifications to the Stress Tensor on several points. This is why the EFE have these inconsistencies. To summarize, the questions are:

- 1/ Why has the Attractive force replaced the original Pressure force?
- 2/ Why has mass replaced volume?
- 3/ Why a concave curvature replaced the convex curvature of the fluid mechanics?
- 4/ Why was the nature of mass ignored in the EFE?

E-12 Explanations to these Enigmas

The theory here presented gives the solution to these four enigmas.

ENIGMA #1

In the 1910s and even today, physicists thought that gravitation was an attractive force. However, in this paper we have shown that an alternative exists. *"Why would Gravity not be a pressure force ?"*. Indeed, a pressure on one side of a sheet of paper produces the same effects as an attractive force on the other side. Here we have demonstrated that the mechanism of pressure of spacetime on matter volumes is much more credible than an attractive force between masses that no one can explain. Moreover, this pressure has a scientific origin: the Stress Tensor.

ENIGMA #2

In the 1910s, the constitution of atoms was unknown. The proton was discovered in 1918 by Rutherford and the neutron in 1932 by Chadwick. In the 1910s, physicists thought that the atom was made only of electrons, and that these electrons were distributed like raisins in a pudding. Therefore, is was impossible for Einstein to separate matter to empty volumes, i.e. massive volumes (protons, neutrons, electrons) to massless volumes (orbitals). This is probably why Einstein took mass instead of matter volumes in the construction of its EFE.

Here we have demonstrated that spacetime is not curved by masses but by matter volumes. Please note that "matter volumes", "Empty volumes", "Standard apparent volumes", "Hermetic apparent volumes", and "Special apparent volumes" are five different definitions of "Volumes" (see the main article and appendix H). Replacing Mass $m = f_{(???)}$ by the Mass Effect $m = f_{(x,y,z,t)}$ defined in Appendix C gives identical results in EFE while solving this second enigma.

ENIGMA #3

Figure E-8 shows that the curvature of the medium in Fluid Mechanics is CONVEX, not CONCAVE. Therefore, the figure E-9 is wrong, despite its popularity. Since the Energy-Momentum Tensor is built from the Stress Tensor, it is obvious that we must keep the basic principles of construction, as shown in figure E-7. To date, no one can explain why has Einstein and Grossman replaced the original CONVEX curvature by a CONCAVE curvature. The most likely explanation is given on the next page: "Consistency of ENIGMAS #1 and #3".

ENIGMA #4

Each side of EFE is built on a 4D space. Since EFE works perfectly, it means that only the second possibility of the precedent "QUESTION #4" can be accepted. The first possibility, i.e. the Energy-Momentum Tensor must be converted to a 5D space, must be rejected. As we know, in any expression, it is not acceptable to have a side in 4D and the other in 5D. Replacing the unknown variable of the mass m = ??? by the Mass Effect m = f(x, y, z, t) gives identical results and solves this fourth enigma. The expression is given in Appendix C:

$$M = \frac{c^2}{G_0} \frac{V}{S} \delta_v = f_{(x,y,z,t)}$$
(209)

$$\begin{split} \mathbf{M} &= \text{Mass effect (kg)} \\ \mathbf{V} &= \text{Volume of the matter volume (m^3)} \\ \mathbf{S} &= \text{Surface of the matter volume (m^2)} \\ \delta_v &= \text{Coefficient of volumetric elasticity of spacetime } \\ \mathbf{c} &= \text{Speed of the light (m/s)} \\ \mathbf{G}_0 &= \text{Universal constant of gravitation} \end{split}$$

Note: To calculate the mass effect as shown in Appendix C, it is simpler to have two constants, 1/ the coefficient of volumetric elasticity of spacetime ϵ_v , and 2/ the density of the surrounding spacetime ρ . On the contrary, in EFE, it is simpler to have only one constant defined as: $\delta_v = \epsilon_v + \rho$. In both cases, the result is identical.

Consistency of ENIGMAS #1 and #3

In physics, a normal attracting constraint is positive by convention, and a pressure constraint is negative. We can also consider that a concave curvature has the "-" sign, and a convex curvature has the "+" sign. Since the signs are given "by convention", this leads to four combinations:

- 1 Gravitation is an attractive force in a concave curvature of spacetime: (+ -)
- 2 Gravitation is an attractive force in a convex curvature of spacetime: (+ +)
- 3 Gravitation is a pressure force in a concave curvature of spacetime: (- -)
- 4 Gravitation is a pressure force in a convex curvature of spacetime: (-+)

These four combinations can be interpreted as:

- 1 (+-) This combination is that of Newton-Einstein. It works perfectly and doesn't need validation. However, despite its popularity, this combination does not explain the origin of the mass, gravitation, and curvature of spacetime,
- 2 (++) This combination does not explain anything and must be rejected,
- 3 (- -) This combination does not explain anything and must be rejected,
- 4 (-+) This combination is that of the proposed theory. It conducts to an identical result as the first combination because (+ -) = (- +). However, this combination is much more credible than the first one because, for identical results, it gives a rational explanation of the preceding enigmas #1 and #3.

E-13 Objections

Simplicity or Complexity?

We could think that the basic laws of physics are extremely complex since the mathematics of physics are. Such is not the case.

Let us consider, for example, a drum. A 5 years old child intuitively knows the principle, namely that by striking it he makes noise. On the other hand, the mathematical description of the surface waves requires the knowledge of Bessel Functions. It is thus advisable to distinguish the principle, **always very simple**, from the laws governing it, which may be extremely complex.

Modern physics shares the same principle. The basic laws of the universe are not included in increasingly complex theories but, on the contrary, in simplicity. Most scientists throughout the world agree with this point of view. For example, to detect any trace of life on Mars, the biologists will not seek complex living organisms but elementary molecules like H_2O .

This view must also be applied to mass, gravitation, and the curvature of spacetime. Here we show that the basic principles of these phenomena are very simple but their formulation complex.

Association of basic ideas

A new theory may be issued from a genius idea, such as GR, or simply the association of some basic ideas that everybody knows. This is the case of the theory here described, which is the association of three basic and well-known ideas:

1/ Differentiating matter volumes (with mass) to empty volumes (without mass),

2/ Considering that the the origin of the curvature of spacetime is not mass but a matter volume. Consequently, the concave curvature becomes a convex curvature (stress tensor),

3/ Replacing the attractive force of gravitation by a pressure force exerted by the curvature of spacetime on matter volumes.

About Einstein

One might wonder why Einstein did not think to the four enigmas discussed above. The main reason is that in the 1910s physicists thought that the atom was built on the "pudding model". Therefore, the volume definition was wrong. However, other reasons can also be retained.

Science is not the matter of only one man such as Einstein, but by thousand scientists. GR was also indirectly built by Gauss, Faraday, Maxwell, Riemann, Christoffel, Ricci, Hooke, Cauchy, Navier, Stokes, Minkowski, Lorentz, Poincarre, Michelson, Morlay, Grosmann, Schwarzschild...

For example: Why in 1915 Einstein did not think to solve the EFE in the case of a simple and basic static sphere? No one knows... This very simple solution was devised by Schwarzschild in 1916, one year after the publishing of the EFE by Einstein.

This example, like many others, shows that a genius like Einstein can build great theories but, for different reasons, can also miss very simple ideas such as the Schwarzschild solution. The theory described here follows the same principle. It is built on basic ideas that many physicists know but do not apply, such as separating matter volumes to empty volumes. Why since the 1910s no one thought to these simple ideas? ...No one knows.

Multiple Solutions

We can also think that, since GR works perfectly, any other theories must be rejected. Imagine, for example, a theory with an equation such as $y = x^2$. This theory is verified if the result is 4. It means that x = 2 is the solution. This result may be widely accepted because it is the most obvious solution, but another solution that could be a better solution also exists: x = -2.

This example shows that we must not systematically reject new ideas, even if they are contrary to theories already established. This is the case of gravitation which may be a pressure force instead of an attractive force.

E-14 New version of the EFE

In equation (208), the time component t is present in c^2 , in $[L^2/T^2]$, and G_0 , in $[L^3/MT^2]$. However, it disappears in the ratio c^2/G_0 , or $[L^2/T^2][L^3/MT^2]$, which gives [L/M] after simplification. In this expression, the quantity M is not a part of the basic four dimensions but is an expression of the kind $m = f_{(x,y,z,t)}$. It means that in the energy-momentum tensor, in the case of non-relativistic objects, the density of matter ρ should be replaced by:

$$\rho = \frac{m}{V} = \frac{c^2}{G_0} \frac{V}{S} \ \delta_v \cdot \frac{1}{V} \qquad (210)$$

After simplification by V, we get:

$$\rho = \frac{c^2}{G_0 S} \delta_v \qquad (211)$$

From these three equations (209), (210) and (211), we can calculate the elements of the Energy-Momentum Tensor. For teaching purpose, the symmetric elements have been duplicated here.

Element T_{00} .

$$T_{00} = \frac{E}{V} = \frac{mc^2}{V} = \frac{c^2}{G_0} \frac{V}{S} \ \delta_v \cdot \frac{c^2}{V}$$
(212)

After simplification by V:

$$T_{00} = \frac{c^4}{G_0 S} \,\delta_v \qquad (213)$$

Elements $T_{01}, T_{02}, T_{03}, T_{10}, T_{20}, T_{30}$. For $\mu\nu = 01, 02, 03, 10, 20$ and 30:

$$T_{\mu\nu} = \rho \ c \ v_{\mu\nu} = \frac{c^2}{G_0 \ S} \delta_v \cdot c v_{\mu\nu} \qquad (214)$$

or:

$$T_{\mu\nu} = \frac{c^4}{G_0 S} \,\delta_v \cdot \frac{v_{\mu\nu}}{c} \qquad (215)$$

Elements T_{11} , T_{22} , T_{33} . σ is a pressure. For $\mu\nu = 11$, 22 and 33 we have:

$$\sigma_{\mu\nu} = \frac{F_{\mu\nu}}{S} = \frac{ma_{\mu\nu}}{S} = \frac{c^2}{G_0} \frac{V}{S} \ \delta_v \cdot \frac{a_{\mu\nu}}{S}$$
(216)

or:

$$\sigma_{\mu\nu} = \frac{c^4}{G_0 S} \,\delta_v \cdot \frac{V}{S c^2} a_{\mu\nu} \qquad (217)$$

Elements $T_{12}, T_{13}, T_{23}, T_{21}, T_{31}, T_{32}$. τ is a pressure. For $\mu\nu = 12, 13, 23, 21, 31$ and 32 we have:

$$\tau_{\mu\nu} = \frac{F_{\mu\nu}}{S} = \frac{ma_{\mu\nu}}{S} = \frac{c^2}{G_0} \frac{V}{S} \ \delta_v \cdot \frac{a_{\mu\nu}}{S}$$
(218)

or:

$$\tau_{\mu\nu} = \frac{c^4}{G_0 S} \ \delta_v \cdot \frac{V}{S \ c^2} a_{\mu\nu} \qquad (219)$$

New Energy-Momentum Tensor

A new Energy-Momentum tensor can be built with the elements of equations (213), (215), (217), and (219). However, we can extract the quantity from each equation:

$$\frac{c^4}{G_0 S} \delta_v \qquad (220)$$

To avoid confusion with the traditional Energy-Momentum tensor $T_{\mu\nu}$, this new tensor is called $J_{\mu\nu}$ (this symbol is different to the Bessel Symbol used in some countries). Including the quantity (220) in the Einstein Constant gives:

$$J_{\mu\nu} = \begin{bmatrix} 1 & v_{01}/c & v_{02}/c & v_{03}/c \\ v_{10}/c & (V/Sc^2)a_{11} & (V/Sc^2)a_{12} & (V/Sc^2)a_{13} \\ v_{20}/c & (V/Sc^2)a_{21} & (V/Sc^2)a_{22} & (V/Sc^2)a_{23} \\ v_{30}/c & (V/Sc^2)a_{31} & (V/Sc^2)a_{32} & (V/Sc^2)a_{33} \end{bmatrix}$$
(221)

Note: Accelerations a_{11} , a_{22} , and a_{33} , are normal accelerations regarding the surface. The other accelerations are tangential.

The coefficient (220) may be merged with the Einstein Constant (161) as follows:

$$\frac{8\pi G_0}{c^4} \cdot \frac{c^4}{G_0 S} \ \delta_v = \frac{8\pi \delta_v}{S} \quad (222)$$

Finally, this new version of EFE can be written as:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi\delta_v}{S}J_{\mu\nu}$$
 (223)

Important note:

As shown in this equation, the two sides of this "new EFE" are in 4D. The variable "m" disappears.

Case of a static sphere

This particular case can be identified to the criteria of construction of the Schwarzschild Metric. We have:

$$V = \frac{4}{3}\pi r^3 \qquad (224)$$

and

$$S = 4\pi r^2 \qquad (225)$$

hence

$$\frac{V}{S} = \frac{4\pi r^3}{3} \frac{1}{4\pi r^2} = \frac{r}{3} \qquad (226)$$

Porting this expression in the tensor $J_{\mu\nu}$ of equation(221) gives, after simplifications:

$$J'_{\mu\nu} = \begin{bmatrix} 1 & v_{01}/c & v_{02}/c & v_{03}/c \\ v_{10}/c & (r/3c^2)a_{11} & (r/3c^2)a_{12} & (r/3c^2)a_{13} \\ v_{20}/c & (r/3c^2)a_{21} & (r/3c^2)a_{22} & (r/3c^2)a_{23} \\ v_{30}/c & (r/3c^2)a_{31} & (r/3c^2)a_{32} & (r/3c^2)a_{33} \end{bmatrix}$$
(227)

In this particular case, the new version of EFE can be written as:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{2\delta_v}{r^2}J'_{\mu\nu}$$
 (228)

E-15 Conclusions

From a mathematical point of view, the principles of inheritance can be applied: "If the Energy-Momentum Tensor is built from the Stress Tensor, automatically it inherits all its properties". It means that if the Stress Tensor applies to volumes, pressures and convex curvature, the Energy-Momentum Tensor must do likewise. The only difference is the number of dimensions: 3D for the Stress Tensor, 4D for the Energy-Momentum Tensor.

On the other hand, adding an unknown variable, the mass "m", only in the right hand side of the EFE creates a lack of homogeneity between the two sides since 'm" is undefined. This is why the traditional EFE has been reconsidered and rewritten.

To use the new version of EFE, we must know the volume and surface of the matter volumes which produce the curvature of spacetime, and the coefficient of curvature of spacetime δ_{v} . It is more convenient to continue using the intermediate variable m in the Energy-Momentum Tensor as usual. In other words, this new version of EFE is only interesting to provide a correct formulation of the mass and gravitation enigmas and to explain the mechanism of physics phenomena such as dark energy. Its mathematical advantages regarding the actual EFE are:

- Curvature of Spacetime ... Produced by Matter Volumes, •
- Nature of the Curvature ... Convex, not concave,
- **Gravitation** ... Pressure force, not an attractive force,
- Mass Enigma ... Comes from the pressure of spacetime on matter volumes,
- **Mass Variable** ... Fully defined $m = f_{(x,y,z,t)}$ (equation 209), **Homogeneity** ... $4D \equiv 4D$ instead of $4D \equiv 5D$ as in EFE. The unknown variable mhas been converted in a 4D variable $m = f_{(x,y,z,t)}$ (equation 209).

E-16 Special Metrics

We could think that the spacetime curvature "C" depends on mass "m" since the expression of the energy-momentum tensor is C = f(m):

$Mass \rightarrow Spacetime curvature$

Since the mass is unknown, this expression is incomplete and should be written as

$??? \rightarrow Mass \rightarrow Spacetime curvature$

The present theory explains the mass origin and proposes to replace "???" by the following group:

 $\begin{bmatrix} Matter \ Volumes \Rightarrow \\ ...curve \ Spacetime \Rightarrow \\ ...that \ exerts \ a \ Pressure \ on \ the \ Body \Rightarrow \\ ...which \ leads \ to \ a \ "Mass \ Effect" \Rightarrow \end{bmatrix}$

\Rightarrow Mass \Rightarrow Spacetime curvature

We see that the spacetime curvature is redundant. So, it is not necessary to specify the last line since the spacetime curvature has already been calculated (second line in italics).

It is important to note that this scheme is static. Its purpose is nothing but to calculate the mass effect in a flat spacetime, as shown in Appendix C.

If we need to know the dynamic spacetime curvature in a particular situation, the first thing to do is to calculate the static mass effect in a flat spacetime. Then, we must use the following scheme:

Step 1

 $\begin{bmatrix} Matter \ Volumes \Rightarrow \\ ...curve \ Spacetime \Rightarrow \\ ...that \ exerts \ a \ Pressure \ on \ the \ Body \Rightarrow \\ ...which \ leads \ to \ a \ "Mass \ Effect" \Rightarrow \end{bmatrix}$

Step 2

\Rightarrow Dynamic spacetime curvature

Step 1: Calculation of the mass effect of the body from the static curvature of spacetime, as shown in Appendix C.

Step 2: The knowledge of the static mass effect "m" allows the calculation of the spacetime curvature in a particular dynamic context. To do that, we must use EFE with correct parameters or one of its known solutions.

For example, for a rotating sphere, the process is to:

- Step 1: Calculate the static mass effect "m" from matter volumes (equation 9, Appendix C),
- Step 2: Calculate the dynamic spacetime curvature as usual, applying the Kerr Metric.

However, it is simpler to continue to calculate special metrics using the traditional EFE as done since the 1920's, even if the mass variable m is undefined.

E-17 The Schwarzschild Metric

To calculate the spacetime curvature with a static body having a spherical symmetry, we must use the Schwarzschild Metric. In that particular case, step 2 isn't necessary since the spacetime curvature has already been calculated from matter volumes in step 1 (see Appendix B).

Using this principle of calculation in two steps, here we show that the theory presented in this paper is in perfect accordance with EFE but, as stated above, this method is not the simplest. For any calculation, we can continue to use the unknown variable "m", even if this calculus is not "academic".

Von Laue Geodesics

F-1 Introduction

A set of concentric circles is drawn in (fig. F-1a). These lines represent the geodesics of spacetime far from any mass, in a Minkowski spacetime. If a static spherical symmetry matter volume is inserted in the centre (fig. F-1b), spacetime will be curved, as explained in the main text. This figure F-1b has been duplicated in fig. F-2.



Figure F-1: Curvature of spacetime



Figure F-2: Von Laue Geodesics

The Von Laue Geodesics has been drawn over the concentric circles of fig. F-1b.

We see that the Von Laue Geodesics match EXACTLY the concentric circles.

In other words, it seems that Von Laue, early as 1927's, predicted the theory presented here. As we see on this figure, **the Von Laue Diagram shows volumes, not masses,** even if the Von Laue Formulas are related to the mass of the body.

New Version of the Equivalence Principle

G-1 Demonstration

Lets consider a static object on Earth (fig. G-1). Matter volumes of this object cause a curvature of spacetime that, in turn, exerts a gravitational force on this object. $g = 9.81m.s^{-2}$ on the surface of Earth.



Figure G-1: Gravitationnal acceleration

Lets now consider the same object accelerated out of any gravitational field. We can represent this object in two different views (fig. G-2a and G-2b).



Figure G-2: Inertial acceleration

In both cases, acceleration "a" is supposed to be identical to g:

$$a = q = 9.81 m.s^{-2}$$

Without any reference, a local observer cannot know if the curvature of spacetime is due to a pressure on the object (fig. G-2a) or to its acceleration (fig. G-2b). In fact, these two figures are identical and depend on where the observer stands, as described in Special Relativity.

Since:

• By definition, $g = 9.81 \text{ m.s}^{-2}$ (fig. G-1) is identical to $a = 9.81 \text{ m.s}^{-2}$ (fig. G-2a and G-2b).

• These examples use the same object. Therefore, the curvature of spacetime produced by the matter volume of this object is identical.

• The two precedent points show that the "mass effect" produced by these curvatures will be necessarily identical in both figures.

We deduce that the "gravitational mass effect" (fig. G-1) is identical to the "inertial mass effect" (fig. G-2):

Gravitational mass effect = Inertial mass effect = Effect from spacetime curvature
Black Holes Simulation

H-1 Introduction

The purpose of this simulation is to show that replacing mass by (matter) volume leads to a consistent and logical explanation of the Schwarzschild Metric and black holes.

Important: This experiment should not be considered as a validation of the theory. However, it is not less credible as those concerning the Big Bang or other debatable simulations conducted in large laboratories.

H-2 Simulation

Figure B-1 of Appendix B has been duplicated on the left of fig. H-1 below. The right part of fig. H-1 shows a simulation of the left part made with an EPP foam. A set of lines, spaced 5 mm apart, has been drawn on an EPP foam which simulates spacetime (See GR for elasticity properties of spacetime). A half-cylinder with a radius of 22 mm. simulates a matter volume. This volume is inserted into the foam.



Figure H-1: Simulation of spacetime in 1D

The gaps between two adjacent lines are computed from the Schwarzschild Metric (see appendix B), which has been reduced to the radial component:

$$y_{r \ exp} = \frac{1}{\left(1 - \frac{2.2}{r}\right)} \qquad (1)$$

The constant 2.2 is calculated from:

- R = 22 mm. Radius of the matter volume of fig. H-1,
- $\epsilon = 0.1$. Arbitrary EPP foam coefficient,

• $\epsilon R = 0.1 \ge 22 \text{ mm} = 2.2 \text{ mm}.$ ($\epsilon R = \Delta R$, see the Schwarzschild Metric)

Col. 1	Col. 2	Col. 3	Col. 4	Col. 5
n	r ₀ = 5n+22	∆r = 0.1 r ₀	r = r₀ - ∆r	У (г)
1	27	2,70	24,30	1,10
2	32	3,20	28,80	1,08
3	37	3,70	33,30	1,07
4	42	4,20	37,80	1,06
5	47	4,70	42,30	1,05
6	52	5,20	46,80	1,05
7	57	5,70	51,30	1,04
8	62	6,20	55,80	1,04
9	67	6,70	60,30	1,04
10	72	7,20	64,80	1,04
11	77	7,70	69,30	1,03
12	82	8,20	73,80	1,03
13	87	8,70	78,30	1,03
14	92	9,20	82,80	1,03
15	97	9,70	87,30	1,03
16	102	10,20	91,80	1,02
17	107	10,70	96,30	1,02
18	112	11,20	100,80	1,02
19	117	11,70	105,30	1,02
20	122	12,20	109,80	1,02
21	127	12,70	114,30	1,02
22	132	13,20	118,80	1,02
23	137	13,70	123,30	1,02
24	142	14,20	127,80	1,02
25	147	14,70	132,30	1,02
26	152	15,20	136,80	1,02

Figure H-2: Theoretical data from formula (1)

Fig. H-2 (previous page) shows data using an incremental index "n" and the above formula (1) where:

• Col. 1: "n" is the rank of each line,

 \bullet Col. 2: Lines " r_0 " are spaced 5 mm. apart, out of gravitation, with an offset of 22, as that

of the radius of the matter volume of fig. H-1,

• Col. 3: Δr is calculated using the elasticity of spacetime, ϵ , supposed to be 0.1,

• Col. 4: Finally, the distance between the base and the point of measurement of the spacetime curvature, r, is computed from ΔR and r_0 ,

• Col. 5: Result.



Figure H-3: Graphic from table H-2

H-3 Black holes

If the radius of the matter volume 22 mm. is increased to 40 mm., a singularity appears under 40 mm. (fig. H-4 and H-5, next page). Above 40 mm., we get a curve as that of fig. H-3.

The table and the curves (fig. H-4 and H-5) show an asymptote when r = 40 mm.. They match exactly that of the behaviour of the Schwarzschild Metric around Rs (Schwarzschild Radius).

We can also remark that the signature is changed from + to -, as inside a black hole. However, the negative part shown by the curve is a "fictive" part that doesn't exist. It's only a mathematical object since below the radius, 40 mm, nothing happens.

To date, no one knows exactly what's inside a black hole. This simulation tends to prove that a black hole is a large matter volume. Moreover, the Schwarzschild Radius is the limit of the matter volume. Inside a matter volume, as inside a black hole, nothing happens.

The light doesn't exist and, therefore, can't escape from the center of a black hole. However, this explanation, as any explanation of black holes, will never be validated by experimentation. It is obvious that no one will never answer the question: "What's inside a black hole?".

Rank	r	y = f (r)		Rank	r	y = f (r)
	22			15	74,4	2,16
1	24,4	-1,56		16	79,3	2,02
2	26,8	-2,03		17	84,4	1,90
3	29,4	-2,77		18	89,5	1,81
4	32,2	-4,13		19	94,5	1,73
5	35,2	-7,33		20	99,5	1,67
6	38,3	-22,53		21	104,2	1,62
8	41,8	23,22		22	108,8	1,58
9	45,9	7,78		23	113,6	1,54
10	50,4	4,85		24	118,4	1,51
11	55,1	3,65		25	123,2	1,48
12	59,9	3,01		26	128	1,45
13	64,6	2,63		27	132,9	1,43
14	69,3	2,37	'			

Figure H-4: Data, increasing the radius to 40mm



Figure H-5: Graphic related to data of fig. H-4

H-4 Conclusions

These simple simulations and explanations show the black hole behaviour. Results are **EXACTLY identical** as in conventional physics using EFE solutions.

This section shows, one more time, that spacetime is curved by (matter) volume, not by mass.

Effectively, during all these simulations and explanations, only lengths have been considered. Mass has been totally ignored. It is the radius, not the mass, which has been increased to 40 mm. in order to calculate a black hole behaviour.

Classes of Volumes

I-1 New Definitions of Volume

One tends to assume that "a volume is a volume". This assumption must be dropped, because some volumes curve spacetime, but others do not. Since GR is grounded on the curvature of spacetime, the first thing to do is to separate the volumes that curve spacetime from the volumes that do not.

Contrary to what we could think, this study about volumes is very important because it leads to the expression of mass: $M = f_{(volume)}$. In reality, spacetime is not curved by mass (or vice-versa), but by a special class of volume. Replacing mass by "matter" volume (see below) changes all the meaning of GR and opens new horizons in astrophysics.

Therefore, we do not have only one type of volume, but five, as described below.

a/ "Matter Volumes" : Volumes that curve spacetime

Example: leptons. Their internal spacetime "pushes" the surrounding spacetime to make room. Thus, matter volumes produce a convex curvature of spacetime (Fig. 3A). Since the latter has the properties of elasticity (Einstein), it exerts a pressure on the surface of these volumes that increases their resistance to movement. As a result, a *mass effect* appears, i.e., an effect having all the characteristics of mass.

For example, consider a particle of volume V crossing a cube of volume 1,000,000 V. The cube will contain 1,000,001 elements of volumes but its overall volume will remain 1,000,000 V because spacetime has the properties of elasticity. Due to this elasticity, the particle will be subject to a pressure from the 1,000,000 other elements of volume.

This pressure, determined by the constraints of the stress tensor (Fig. 1A), leads to a mass effect. See appendix C to get the mass expression $m = f_{(x,y,z,t)}$.

b/ "Empty Volumes" : Volumes that do not curve spacetime.

Empty volumes are simply vacuum, but are often found in various forms, such as the volumes of orbitals. These volumes exist but are transparent regarding spacetime. Therefore, open volumes are massless since no curvature \equiv no pressure of spacetime, and no pressure of spacetime means no mass effect.

c/ "Apparent Volumes" : Combination of matter and empty volumes

Objects we use daily are defined as:

Apparent Volumes = \sum Matter Volumes + \sum Empty Volumes

These volumes, mainly atoms and molecules, are combinations of volumes with mass (matter volumes) and massless volumes (empty volumes). In atoms, for example, the nucleus and electrons curve spacetime and have mass, and are "matter volumes", whereas the orbitals do not and are "empty volumes".

We have the feeling that mass and volume are two different quantities. In reality, it is the proportion of matter to empty volumes in "apparent volumes" that varies from one atom to another, from one molecule to another, and from one object to another, which gives us this feeling. It means that mass and (matter) volume are two different views of the same quantity, such as particles and waves in the wave-particle duality.

The main problem is not to understand this very simple explanation, but to drop our preconceived ideas saying that mass and volume are two different concepts.

d/ "Hermetic Volumes"

These volumes are also combinations of matter and empty volumes but their global volume is hermetic regarding spacetime. For example, a nucleus is made of nucleons (matter volumes), separated by empty space (empty volume). Whatever the content of this empty space (gluons...), this volume exists. The behaviour of this global volume is that of a matter volume regarding spacetime. Consequently, the whole volume of the nucleus deforms spacetime and, therefore, gets mass. This explains why the volume (\equiv mass) of the nucleus is greater than those of its nucleons (mass excess). This also explains why the volume of protons, 938 MeV/c², is much greater than the volumes of the three quarks, 2.3 + 2.3 + 4.8 = 9.4 MeV/c², taking into account the volumes of gluons and the empty space (empty volumes) between the components.

e/ "Special Volumes"

These kinds of volumes are called "special" because we do not know their behaviour regarding spacetime. This is the case of ⁶He, ⁸He, ¹⁴Be... For example, ¹¹Li has a core with 3p6n and a halo of 2n. Since we do not know the penetration of spacetime inside these nuclei, these volumes can not be classified in any of the preceding categories.

I-2 Can we replace mass by volume?

We must bear in mind there it does not exist only one kind of volume, but five. This leads to the following questions:

1/ Can we replace mass by "volume"?

No, because the word "Volume" is undefined.

2/ Can we replace mass by "apparent volumes"?

No, because apparent volumes are a combination of matter volumes, with mass, and massless empty volumes. Only matter volumes (with mass) are significant.

3/ Can we replace mass by "matter volumes" or "hermetic volumes"?

Yes, because there is a relation between matter volumes or hermetic volumes and mass (see the main article and appendix C). It is obvious that if we have $V = f_{(M)}$ and reciprocally $M = f_{(V)}$, we can use indifferently V or M in our GR calculus.

Applications

The following applications are described in the PDF document [1] or on the official website [2]. Please note that these documents has been written for non-physicist people and therefore contains very simplified explanations and cartoon-like figures. However, these explanations of phenomena of physics are correct.

J-1 Applications

- Explanation of $E=mc^2$,
- Explanation of the variation of mass of relativistic particle,
- Explanation of the deviation of light,
- Explanation of the mass excess,
- Explanation of the mass defect,
- Explanation of the twin paradox,
- Black Holes,
- Galactic Clusters,
- Galaxy Filaments,
- Increase of expansion of the universe,
- Dark Matter.

[1] http://www.spacetime-model.com/files/mass.pdf[2] http://www.mass-gravity.com

References

Articles

[3] Einstein, Albert (1907), "Uber das Relativitatsprinzip", Jahrbuch der Radioaktivitat 4: 411 [4] Einstein, Albert (1915), "Die Feldgleichungen der Gravitation", SPA der Wissenschaften: 844-847 [5] Friedrich, Helmut (2005), "Is general relativity understood?", Annalen Phys. 15 (1-2): 84-108 [6] Brans, C. H. (1961), "Relativistic Theory of Gravitation", Physical Review 124 (3): 925-935 [7] Will, Clifford M. (2006), "General Relativity and Experiment", Living Rev. Relativity 9 [8] Einstein, Albert (1916), "Der allgemeinen Relativitatstheorie", Annalen der Physik 49 [9] Einstein, Albert (1917), "Kosmologische ... Relativitatstheorie", SPA der Wissenschaften 142 [10] Janssen, Michel (2005), "Einstein's bumpy road to GR", Ann. Phys. 14 (S1): 58-85 [11] Ashtekar, Abhay (1987), "New Hamiltonian formulation of GR", Phys. Rev. D 36 (6): 1587-1602 [12] Havas, P. (1964), "Newtonian mechanics their relation to GR", Rev. Mod. Phys. 36 (4): 938-965 [13] Komar, Arthur (1959), "Covariant Conservation Laws in GR", Phys. Rev. 113 (3): 934-936 [14] Lovelock, David (1972), "The 4D Space and the Einstein Tensor", J. Math. Phys. 13 (6): 874-876 [15] Font, Jose A. (2003), "Numerical Hydrodynamics in GR", Living Rev. Relativity 6, doi:10.12942 [16] Ashtekar, Abhay (1986), "New variables for ... gravity", Phys. Rev. Lett. 57 (18): 2244-2247 [17] Schwarzschild, Karl (1916a), "Gravitationsfeld ... Einsteinschen Theorie", SPA Wiss.: 189-196 [18] Schwarzschild, Karl (1916b), "Gravitationsfeld ... Einsteinschen Theorie", SPA Wiss.: 424-434 [19] Penrose, Roger (1969), "Gravitational collapse: the role of GR", Riv. Nuovo Cim. 1: 252-276 [20] Shapiro, Irwin (1964), "Fourth test of general relativity", Phys. Rev. Lett. 13 (26): 789-791 [21] Shapiro, Irwin (1968), "Fourth test of general relativity", Phys. Rev. Lett. 20 (22): 1265-1269 [22] Caldwell, Robert R. (2004), "Dark Energy", Physics World 17 (5): 37-42 Books This paper is also grounded on the following books and documents. This list is not exhaustive. [23] Arnowitt, Richard (1962), "The dynamics of general relativity", Wiley [24] Carroll, Sean, "Lecture Notes on GR", http://preposterousuniverse.com/grnotes/ [25] Damour, Thibault (2005), "Einstein aujourd'hui", EDP Sciences [26] Carroll, Sean (2004), "An Introduction to GR", Addison-Wesley, ISBN 0-8053-8732-3 [27] Ludvigsen, Malcolm (2003), RG, "Une approche geometrique", Dunod, ISBN 2-1004-9688-3 [28] Hladik, Jean (2006), "Introduction a la relativite generale niveau M1", Ellipses [29] Belabar, Francoise (1999), "Einstein, RG, cosmologie et theories unitaires", Seuil [30] Hladik, Jean (1999), "Le calcul tensoriel en physique", Ellipses [31] Brown, Kevin (2005), "Reflections on relativity", Cambridge University Press [32] Ehlers, Jurgen (1973), "Survey of general relativity theory", Reidel [33] Moor, Rafi (2006), "Understanding GR", http://www.rafimoor.com/english/GRE.htm [34] Gron, Oyvind, Hervik (2007), "Einstein's GR", Springer, ISBN 978-0-387-69199-2 [35] B. F. Schutz (2001), "A First Course in GR", ISBN 9780-521-88705-2 [36] Misner, Wheeler (1973), "Gravitation", Freeman, ISBN 0-7167-0344-0 [37] Hartle, James B. (1998), "An Introduction to Einstein's GR", ISBN 0-8053-8662-9 [38] Barreau, Aurelien (2011) "Grain, Julien, Relativite Generale", Dunod [39] Einstein, Albert (1916), "La theorie de la Relativite", Dunod [40] Efstathiou, George (2009), "Relativite Generale", De Boeck [41] Einstein, Albert (1917), "La Relativite", Payot [42] Emile, Jean (1987), "La relativite complexe", Albin Michel [43] Elbaz, Edgard (1998), "Relativite Generale et Gravitation", Ellipses [44] Feynman, Richard (1956), "Cours de Physique", Dunod [45] Eisenstaedt, Jean (2002), "Einstein et la relativite generale", CNRS [46] Hughston, (2003), "Introduction to GR", Cambridge Univ. Press, ISBN 0-1985-2957-0 [47] Nordstrom, Gunnar (1918), "On the Energy of the Gravitational Field", Wetenschap., 1918 [48] Brillouin, Leon (1938), "Les tenseurs", Valentin, 1938 [49] Denis-Papin, Kauffman (1955), "Cours de calcul tensoriel", C4pf, 1953

- [50] Jerome Jacky, (2007), http://www.spacetime-model.com/higgsboson/higgs_boson.htm
- [51] Jerome, Jacky (2007), "The Spacetime model", AC42 Publications, ISBN 978 2953 1234-0-1