

# Concept of Gravitational Acceleration Field and its Consequences for the Compact Stellar Objects

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**Abstract:** In previous papers [1,2] relating to the Combined Gravitational Action (CGA), we have exclusively studied the orbital motion without spin. In the present paper we apply the CGA to any self-rotating material body, *i.e.*, axially spinning massive object, which itself may be locally seen as a gravitrotational source because it is capable of generating the gravitrotational acceleration field, which seems unknown in the previously existing gravity theories. The consequences of such an acceleration field are very interesting particularly for the compact stellar objects.

**Keywords:** CGA, gravitrotational acceleration field, gravitrotational energy, neutron stars, pulsars

## 1. Introduction

We have previously [1,2] shown that the Combined Gravitational Action (CGA) as an alternative gravity theory is very capable of predicting and explaining some old and new gravitational phenomena. For example, in [2], we have investigated the CGA-spin-orbit coupling precession and applied CGA to large-scale structures and the problem of galactic rotation curves has been resolved. Also the Modified Newtonian Dynamics (MOND) [3,4,5,6] as an alternative theory to the dark matter (DM) paradigm became by means of CGA [2] an additional support for DM!

Conceptually, the CGA is basically founded on the concept of the combined gravitational potential energy (CGPE) defined by the expression

$$U \equiv U(r, v) = -\frac{k}{r} \left( 1 + \frac{v^2}{w^2} \right), \quad (1)$$

where  $k = GMm$ ;  $G$  being the Newton's gravitational constant;  $M$  and  $m$  are the masses of the gravitational source  $A$  and the moving test-body  $B$ ;  $r = \sqrt{(x - x_0)^2 + (y - y_0)^2 + (z - z_0)^2}$  is the relative distance between  $A$  and  $B$ ;  $v = \sqrt{v_x^2 + v_y^2 + v_z^2}$  is the velocity of the test-body  $B$  relative to the inertial reference frame of source  $A$ ; and  $w$  is a specific kinematical parameter having the physical dimensions of a constant velocity defined by

$$w = \begin{cases} c_0, & \text{if } B \text{ is in relative motion inside the vicinity of } A \\ v_{\text{esc}} = \sqrt{2GM/R}, & \text{if } B \text{ is in relative motion outside the vicinity of } A \end{cases}, \quad (2)$$

where  $c_0$  is the light speed in local vacuum and  $v_{\text{esc}}$  is the escape velocity at the surface of the gravitational source  $A$ .

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Hence, starting from the CGPE and using only the very familiar tools of classical gravitomechanics and Euler-Lagrange equations, we have established the CGA-formalism [1,2]. The main consequence of CGA is the dynamic gravitational field (DGF),  $\Lambda$ , which is phenomenologically an induced field, it is more precisely a sort of gravitational induction due to the relative motion of material body in the vicinity of the gravitational source[1,2]. The magnitude of DGF is of the form

$$\Lambda = \pm \frac{GM}{r^2} \left( \frac{v}{w} \right)^2. \quad (3)$$

Eq.(3) means that DGF may play a double role, that is to say, when perceived/interpreted as an extra-gravitational acceleration,  $\Lambda > 0$ , or an extra-gravitational deceleration,  $\Lambda < 0$ , (see Ref. [1] for a detailed discussion).

In previous papers [1,2], we have focused our interest on the orbital motion and gravitational two-body problem. In the present paper, we shall apply CGA to any self-rotating (spinning) material body, *i.e.*, axially rotating massive object that is itself may be locally seen as a gravitorotational source since it is capable of generating the gravitorotational acceleration field,  $\lambda$ , which seems unknown in the previously existing gravity theories.

## 2. Concept of the gravitorotational acceleration field

Phenomenologically speaking, the concept of the gravitorotational acceleration field vector (GRA),  $\lambda$ , is very similar to DGF, that is if  $\Lambda$  is mainly induced by the relative motion of the massive test body in the vicinity of the principal gravitational source, the GRA is intrinsically generated by any massive body in rotational motion independently of the principal gravitational source, which itself may be characterized by its proper GRA during its axial-rotation, therefore, the gravitorotational acceleration field is, in fact, a combination of *gravity* and *rotation*.

## 3. Expression of GRA

In order to derive an explicit expression for GRA, let us first rewrite Eq.(3) for the case when  $\Lambda > 0$ , that is

$$\Lambda = \frac{GM}{r^2} \left( \frac{v}{w} \right)^2, \quad (4)$$

and considering a massive body of mass  $M$  and radius  $R$ , which is intrinsically in axial-rotation in its proper reference frame at rotational velocity of magnitude  $v_{\text{rot}} = \Omega R$  independently of the presence of any other gravitational source. Therefore, according to the concept of GRA, in which a case, the rotating massive body should be locally seen as a gravitorotational source when  $\|\Lambda\| \rightarrow \|\lambda\| \equiv \lambda$  as  $r \rightarrow R$ ,  $v \rightarrow v_{\text{rot}}$  and  $w \rightarrow c_0$ , thus (4) becomes after substitution

$$\lambda = \frac{GM}{R^2} \left( \frac{\Omega R}{c_0} \right)^2. \quad (5)$$

Since  $\Omega = 2\pi P^{-1}$ , where  $P$  is the rotational period, hence we get after substitution in (5) the expected expression of GRA

$$\lambda = GM \left( \frac{2\pi}{c_0 P} \right)^2. \quad (6)$$

It is clear from Eq.(6), GRA  $\lambda$  depends exclusively on the mass and rotational period, therefore, mathematically may be treated as a function of the form

$$\lambda \equiv \lambda(M, P). \quad (7)$$

The structure of Eq.(6) allows us to affirm that for any astrophysical massive object, the magnitude of  $\lambda$  should be infinitesimally small for slowly rotating massive stellar object and enormous for rapidly rotating ones. Further, in order to confirm numerically our affirmation, we have selected seven well-known (binary) pulsars and calculated their GRAs, and compared them with the Sun's GRA. The values are listed in Table 1.

OBJECT	$P$	$M$	$\lambda$	REF.
Sun + PRS	(s)	( $M_{\odot}$ )	( $\text{m s}^{-2}$ )	
Sun	$2.358720 \times 10^6$	1	$1.047211 \times 10^{-8}$	
B 1913+16	$5.903000 \times 10^{-2}$	1.4410	$2.409380 \times 10^7$	a
B 1534+12	$3.790000 \times 10^{-2}$	1.3400	$5.435171 \times 10^7$	b,c
B 2127+11C	$3.053000 \times 10^{-2}$	1.3600	$8.501044 \times 10^7$	d
B 1257+12	$6.200000 \times 10^{-3}$	1.4000	$2.121932 \times 10^9$	e
J 0737-3039	$2.280000 \times 10^{-2}$	1.3381	$1.500000 \times 10^8$	f
B 1937+21	$1.557800 \times 10^{-3}$	1.4000	$3.364000 \times 10^{10}$	g
J 1748-2446ad	$1.395000 \times 10^{-3}$	1.4000	$4.194982 \times 10^{10}$	h

**Table 1:** The values of GRA for seven well-known (binary) pulsars compared with the Sun's GRA value.

**Ref.:** a) Taylor and Weisberg [7]; b) Arzoumanian [8]; c) Wolszcan [9]; d) Deich and Kulkarni [10]; e) Konacki and Wolszcan [11]; f) Kramer and Wex [12]; g) Takahashi *et al.* [13]; h) Hessels *et al.* [14].

**Note:** To calculate these values, we have used  $G = 6.67384 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$ ;  $c_0 = 299792458 \text{ ms}^{-1}$ ;

$M_{\odot} = 1.9891 \times 10^{30} \text{ kg}$  and  $P_{\odot} = 27.30 \text{ d}$ .

The analysis of Table 1 gives us the following results: 1) The magnitude of the Sun's GRA,  $\lambda_{\odot} = 1.047211 \times 10^{-8} \text{ ms}^{-2}$ , is extremely weak that's why its effect on the solar system is unobservable, but perhaps only the Sun's immediate vicinity that should be concerned by it. Since GRA is explicitly independent of the radius of rotating massive object, thus the extreme weakness of the Sun's GRA is mainly due to the huge value of the rotational period,  $P_{\odot} = 2.358720 \times 10^6 \text{ s}$ , compared with those of the pulsars. 2) In spite of the fact that the PSR's masses are nearly equal, the PSR's rotational periods show a neat inequality between them. Also, the different values of GRA for each celestial object show us how the GRA is so sensitive to the variation in rotational period.

#### 4. Mutual dependence between the mass and the rotational period

Since GRA may be treated as a function of the form (7), thus we can show more clearly the existence of the mutual dependence between the mass and the rotational period of the same rotating body *via* GRA. For this purpose, we deduce from Eq.(6) the following expression

$$\frac{M}{P^2} = \left( \frac{c_0^2}{4\pi^2 G} \right) \lambda. \quad (8)$$

Obviously, Eq.(8) shows us the expected mutual dependence between the mass and rotational period *via* GRA. Furthermore, because the rotational period is an intrinsic physical quantity, here, according to Eq.(8) the spin of any massive celestial body should vary with mass independently of cosmic time.

#### 5. Link between GRA and rotational acceleration

Now, returning to Eq.(6) and showing that GRA and the rotational acceleration

$$a_{\text{rot}} = \Omega^2 R, \quad (9)$$

are in fact proportional,  $\lambda \propto a_{\text{rot}}$ , and the constant of proportionality is precisely the compactness factor  $\varepsilon = GM/c_0^2 R$  that characterizes any massive celestial body. To this end, it suffices to multiply and divide by the radius,  $R$ , the right hand side of Eq.(6) to get the expected expression

$$\lambda = \varepsilon a_{\text{rot}}. \quad (10)$$

According to the expression (10), GRA is at the same time an *old* and a *new* natural physical quantity that should play a crucial role, specially, for the compact stellar objects like, *e.g.*, the rotating neutron stars and pulsars to which the compactness,  $\varepsilon$ , has a large value compared to that of normal stellar objects. As illustration, the Sun's compactness has the value  $\varepsilon_{\odot} = 4.926858 \times 10^{-6}$ .

## 6. Consequences of GRA

In what follows, we will show that, in the context of CGA, the transitional state, stability and instability of the uniformly rotating neutron star (NS) depending on the ‘*antagonism*’ between centrifugal force and gravitational force or in energetic terms between rotational kinetic energy (RKE) and gravitational binding energy (GBE).

Usually, the physics of NS considered the source of the emitted energy is essentially the RKE, however, such a consideration should immediately imply that at least, in the medium term, the gravitational binding energy should absolutely dominate RKE and as a result the NS should be prematurely in a state of gravitational collapse. Hence, as we will see, the main source of the emitted energy is not the RKE but the gravitorotational energy (GRE), a sort of new physical quantity which is a direct consequence of GRA.

Now, let us determine the conditions of transitional state, dynamical stability and dynamical instability that may be characterized any NS at least in the medium term. To this end, we assume a uniformly rotating NS as a homogeneous rigid spherical body of mass  $M$ , radius  $R$  and rotational velocity  $\Omega = 2\pi/P$ , where  $P$  is the rotational period. Its RKE and GBE are, respectively, defined by the well-known formulae:

$$E_{\text{rot}} = I\Omega^2/2, \quad (11)$$

and

$$E_{\text{G}} = -\frac{3}{5}\frac{GM^2}{R}, \quad (12)$$

where  $I = 2MR^2/5$  is the moment of inertia of NS under consideration.

The total energy is

$$W = E_{\text{rot}} + E_{\text{G}}, \quad (13)$$

this puts forward the following conditions:

- a)  $W < 0$ , NS is in a state of dynamical stability,
- b)  $W = 0$ , NS is in a state of transition,
- c)  $W > 0$ , NS is in a state of dynamical instability.

It is worth noting that the three suggested conditions a, b and c are taken in the medium term because NS may be suddenly in a state of dynamical perturbation or in a state of transition from stability to instability and vice versa.

## 7. Critical rotational period

Knowing the critical rotational period (CRP) of NS is highly important thing because CRP should be treated as a parameter of reference on which the temporal evolution of NS depending. Furthermore, since the change from stability to instability and vice versa should obligatory pass *via* the transitional state, therefore, from the latter we deduce an expression for theoretical period, and we find, after performing a simple algebraic calculation

$$P_c = 2\pi R \sqrt{\frac{R}{3GM}} . \quad (14)$$

We can numerically evaluate the CRP by taking, through this paper, the standard NS mass and radius are, respectively:  $M = 1.4M_\odot$  and  $R = 10\text{km}$ , and we get

$$P_c = 2.660963 \times 10^{-4} \cong 0.2661 \text{ ms} , \quad (15)$$

which is a tiny fraction of the smallest yet observed rotational period,  $P = 1.3950 \text{ ms}$ , of PRS J1748-2446ab [14]

## 8. Gravitational energy

Now, we are arriving at the most important consequence of GRA, namely, the gravitrotational energy (GRE), which should qualitatively and quantitatively characterize any massive rotating body. As we will see, GRE is quantitatively comparable to the amount of RKE, particularly, for NS and pulsars.

Since GRE is a direct consequence of GRA, thus GRE should be proportional to GRA, *i.e.*,  $\mathcal{E} \propto \lambda$  or equivalently

$$\mathcal{E} = \kappa \lambda . \quad (16)$$

Let us determine the expression of the proportionality constant,  $\kappa$ , by using the dimensional analysis as follows.

$$[\kappa] = \frac{[\mathcal{E}]}{[\lambda]} = \frac{\text{ML}^2\text{T}^{-2}}{\text{LT}^{-2}} = \frac{\text{ML}^2}{\text{L}} .$$

Remark, the dimensional quantity  $\text{ML}^2$  has the physical dimensions of the moment of inertia, therefore,  $\kappa$  should take the form  $\kappa = I/R$  and consequently, we find the expected expression for GRE:

$$\mathcal{E} = \frac{\lambda I}{R} . \quad (17)$$

In order to show that the amount of GRE  $\mathcal{E}$  is quantitatively comparable to that of RKE, we use the same sample of seven (binary) pulsars plus the Sun listed in Table 1. The numerical values of  $\mathcal{E}$  are listed in Table 2.

OBJECT	$E_{\text{rot}}$	$\mathcal{E}$
Sun + PRS	(J)	(J)
Sun	$1.365519 \times 10^{36}$	$5.794993 \times 10^{30}$
B 1913+16	$6.494783 \times 10^{41}$	$2.764710 \times 10^{41}$
B 1534+12	$1.465117 \times 10^{42}$	$5.799600 \times 10^{41}$
B 2127+11C	$2.291561 \times 10^{42}$	$9.206431 \times 10^{41}$
B 1257+12	$5.719933 \times 10^{43}$	$2.365592 \times 10^{43}$
J 0737-3039	$4.042638 \times 10^{42}$	$1.597991 \times 10^{42}$
B 1937+21	$9.060475 \times 10^{44}$	$3.747140 \times 10^{44}$
J 1748-2446ad	$1.129863 \times 10^{45}$	$4.672774 \times 10^{44}$

**Table 2:** comparison between the numerical values of  $E_{\text{rot}}$  and  $\mathcal{E}$  for the Sun and seven well known (binary) pulsars.

-Analysis of Table 2: The numerical values listed in Table 2 show us, excepting the Sun' values, all the values of  $E_{\text{rot}}$  and  $\mathcal{E}$  are comparable for the seven (binary) pulsars. This fact is mainly due, at the same time, to the rotational period and the compactness. To illustrate this fact, the expression (17) may be written as follows:

$$\mathcal{E} = \left( \frac{GM}{c_0^2 R} \right) I \Omega^2 = 2 \left( \frac{GM}{c_0^2 R} \right) E_{\text{rot}} . \quad (18)$$

From all that we arrive at the following result: In the context of CGA, the RKE cannot be considered as the main source of the emitted energy for rotating NS and pulsars because its own role is to balance, approximately, the GBE, at least in the medium term. Therefore, the veritable principal source of the emitted energy should be undoubtedly GRE. This affirmation is supposed by GRE's numerical values listed in Table 2, which are quantitatively comparable to those of RKE for pulsars.

## 9. Rotating magnetars

Rotating magnetized NS (magnetars) are also important compact stellar objects, that's why it is possible to exploit GRE as an energetic reservoir for rotating magnetars by assuming that there is a certain physical mechanism that can convert all GRE into maximum magnetic energy:

$$E_{\text{max}} = B_{\text{max}}^2 R^3 = \mathcal{E} , \quad (19)$$

this could, of course, produce a maximum magnetic field strength

$$B_{\max} = \sqrt{\mathcal{E} R^{-3}}, \quad (20)$$

where  $B_{\max}$ ,  $\mathcal{E}$  and  $R$  should be expressed in gauss, erg and cm, respectively. And as illustration, let us evaluate the maximum magnetic field strength of radio pulsar B 1931+24. We have according to Ref. [15] the following parameters:  $P = 0.813\text{s}$  and  $B_0 \cong 3 \times 10^{12} \text{G}$ . By taking, as usual,  $M = 1.4M_{\odot}$  and  $R = 10\text{km}$ , we find for GRE  $\mathcal{E} = 1.375758 \times 10^{46} \text{erg}$ , and after substitution in (20), we get

$$B_{\max} = 1.173 \times 10^{14} \text{G}. \quad (21)$$

## 10. Conclusion

Basing on our gravity model, Combined Gravitational Action, we have derived an explicit expression for the concept of gravitorotational acceleration field (GRA), which is unknown in the previously established gravity theories. The most significant result of GRA is the gravitorotational energy (GRE), which should qualitatively and quantitatively characterize any massive rotating body. Furthermore, GRE should be exploited as an energetic reservoir, particularly, for NS and pulsars.

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