The cubic equation's relation to the fine structure constant, the mixing angles, and Weinberg angle

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A special case of the cubic equation is shown to possess three unusually economical solutions. A minimal case associated with these solutions is then shown to yield a congruous set of numbers that fit the fine structure constant, the sines squared of the quark and lepton mixing angles, as well as the Weinberg angle. Had Renaissance mathematicians probed the cubic equation's solutions more deeply these numbers might have formed a well-known part of algebra from the 16th century.

I. INTRODUCTION

In [1] the author showed how a slightly asymmetric equation (that is, an equation whose left- and right-hand sides are very similar) produced the experimental value of the fine structure constant (approximately 1/137.036) [2, 3], which elsewhere was tied to the sines squared of the quark and lepton mixing angles [4, 5]. Here, a special case of the cubic equation is shown to possess three unusually economical solutions, where a minimal case associated with the above solutions produces a set of numbers that fit the fine structure constant, the sines squared of the quark and lepton mixing angles, and the Weinberg angle. This article extends the results of [6].

II. SPECIAL CASE OF THE CUBIC EQUATION

Let

$$Z = \left(\frac{m+x}{n}\right)^3 + (m+x)^2 \quad , \tag{2.1}$$

where x is a variable, Z a positive constant, and m and n are positive integer constants such that

$$m = \frac{n^3}{3} \quad . \tag{2.2}$$

III. FIRST SOLUTION

Then, by defining

$$W = \left(\frac{m}{n}\right)^3 + (m)^2 \tag{3.1}$$

$$u = 2\frac{Z}{W} - 1 \tag{3.2}$$

$$v = \sqrt[3]{u \pm \sqrt{u^2 - 1}} \tag{3.3}$$

$$w = v + \frac{1}{v} - 1 \tag{3.4}$$

while choosing Z and n so that $Z \geq W$ (and therefore $u^2 - 1 \geq 0$), we have

$$x = m\left(w - 1\right) \quad , \tag{3.5}$$

which can be shown to be the first of the three solutions to be given for Eq. (2.1).

IV. MINIMAL CASE

The smallest positive integers fitting Eq. (2.2)

$$m=9$$
 and $n=3$

are notable, simply because they are minimal.

V. FINE STRUCTURE CONSTANT

For the above minimal case the solution x = 1 to Eq. (2.1) requires that

$$Z = \left(\frac{9+1}{3}\right)^3 + (9+1)^2$$
$$= 137.\overline{037} .$$

Here, the constant Z is close enough—within one thousandth of one per cent—to the reciprocal of the fine structure constant 137.036 [2, 3] to suggest that looking for a connection might turn up interesting mathematics.

In fact it does [7]. Now, if we let

$$m = 9$$
 $n = 3$ $Z = 137.036$,

then

$$W = (m/n)^{3} + (m)^{2}$$
$$= (9/3)^{3} + 9^{2}$$
$$= 3^{3} + 3^{4}$$
$$= 108$$

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$$u = 2 \times Z/W - 1$$

$$= 2 \times 137.036/108 - 1$$

$$\approx 2 \times 1.268851851 - 1$$

$$\approx 1.537703703$$

$$v = \sqrt[3]{u \pm \sqrt{u^2 - 1}}$$

$$\approx 1.393479170916^{\pm 1}$$

$$w = v + 1/v - 1$$

$$\approx 1.111107407399$$

Therefore,

$$x = m (w - 1)$$

$$\approx 9 \times (1.111 \, 107 \, 407 \, 399 - 1)$$

$$\approx 0.999 \, 966 \, 666 \, 591$$

$$\approx 1 - \frac{1}{29 \, 999.932 \, 142 \, 743 \, 338}$$

$$\approx 1 - \frac{1}{3 \times 10^4} \quad . \tag{5.1}$$

Substituting the values for m, etc. into Eq. (2.1) gives

$$137.036 = \left[\left(\frac{9+1}{3} \right) - \frac{1}{3 \times 29999.932...} \right]^{3} + \left[(9+1) - \frac{1}{29999.932...} \right]^{2} . \quad (5.2)$$

VI. AN IMPORTANT MINIMUM

Now suppose that Z takes the form

$$Z = \frac{M^3 - M^{-3}}{n^3} + M^2 - M^{-3} \quad , \tag{6.1}$$

where

$$M = m + 1 \quad .$$

Then, as proven in Theorem 2 in [7], a surprisingly precise, but simple, approximate solution to Eq. (2.1) follows

$$x \approx 1 - \frac{1}{3 \times M^4}$$
 (6.2)

For M = 10 this equation recovers Eq. (5.1), so it should come as no surprise that, for the minimal case assumed in the previous section, Eq. (6.1) gives

$$Z = \frac{10^3 - 10^{-3}}{3^3} + 10^2 - 10^{-3}$$
$$= 137\,036$$

In the previous section for the minimal case we had to assign 137.036 to Z "by hand"; now in this section we learn that the value 137.036 would have been assigned automatically had we combined the minimal case with Eq. (6.1). In this way we see that 137.036 is itself an important minimum, and hence of purely mathematical interest (i.e., apart from its role as a constant of physics).

VII. QUARK AND LEPTON MIXING ANGLES

Moreover, the following four quantities seen in Eq. (5.2)

$$\begin{array}{c}
10\\
3\\
10
\end{array}
\qquad
\begin{array}{c}
1\\
3 \times 29999.932...\\
10\\
\hline
29999.932...
\end{array}$$

which can be reproduced from the sines squared of the quark and lepton angles, are also of purely mathematical interest, independent of their role in physics. Specifically, values such as 10/3, 10, etc. can be produced from the quark and lepton mixing angles L12, L13, L23, Q12, Q13, Q23, as follows

$$\begin{array}{ccc}
10/3 \approx & 1/\sin^2 L12 \\
1/3 \times 29\,999.932 \approx & \sin^2 Q13 \\
& 10 \approx \sin^2 L23 \times 1/\sin^2 Q12 \\
& 1/29\,999.932 \approx \sin^2 Q23 \times & \sin^2 L13
\end{array} \right} (7.1)$$

where a mathematical model conforming to the above relations, and predicting mixing angles (in degrees) of $33.210\,911$, $8.034\,394$, and 45 for leptons—and $12.920\,966$, $0.190\,986$, and $2.367\,442$ for quarks—is detailed in [5]; there, matrix algebra is used to impose three constraints on mixing, just one of which is independent of the four constraints imposed by Eq. (7.1). It is this additional constraint, along with the further constraint that $L23=45^{\circ}$, that allows the four constraints of Eq. (7.1) to produce the six mixing angles predicted above, which are all within the limits of experimental error.

VIII. WEINBERG ANGLE

Consider the similarity of Eq. (2.1) to Eq. (3.1): One cannot help but notice that the constants W and Z appear on similar footing mathematically. Hence, one might expect W and Z to appear on similar footing phenomenologically. In this section we show that W/Z can be used to fit economically the ratio of the W- and Z-boson masses.

At the outset, the variable names W and Z were chosen in anticipation of their use in a formula reproducing the ratio of the W- and Z-boson masses—that is, M_W/M_Z . That there appears to be some relationship between W/Z and M_W/M_Z is shown by

$$\frac{W}{Z} \approx \left(\frac{M_W}{M_Z}\right)^2$$

$$\approx \cos^2 \theta_W \quad , \tag{8.1}$$

where W = 108 and Z = 137.036, as before, and θ_W is the simplest of the Weinberg angle's definitions [8, 9].

Using the precisely-measured mass $M_Z=91.1876\pm0.0021$ GeV, we can calculate the value of M_W with the

aid of Eq. (8.1), giving $M_W \approx 80.951$ GeV. Experimentally, $M_W = 80.385 \pm 0.015$ GeV [10]. This calculated M_W differs from experiment by 1 part in 142 and is out of range of experiment; but uncertainty over the best definition of θ_W suggests that a modified, but still valid, Eq. (8.1) might give a better fit [8, 9].

In summary, the above mathematics fits the fine structure constant, the Weinberg angle, and (with the aid of two additional constraints) the six mixing angles L12, L13, L23, Q12, Q13, Q23. This says something for the congruity and efficiency of the above "cubic" framework in modeling fundamental constants. Moreover, as will be shown in the next two sections, the cubic equation responsible for all this possesses two unusually compact alternative solutions.

IX. SECOND SOLUTION

The second solution to Eq. (2.1) requires defining

$$\cos \theta_C = \sqrt{\frac{W}{Z}} \quad , \tag{9.1}$$

where $0 \le \theta_C < \pi/2$, so that

$$\sin \theta_C = \sqrt{1 - \frac{W}{Z}} \quad . \tag{9.2}$$

It can then be shown that Eq. (3.3) can be restated

$$v = \sqrt[3]{\frac{1 + \sin \theta_C}{1 - \sin \theta_C}} \quad , \tag{9.3}$$

so that Eq. (3.4) gives

$$w = \sqrt[3]{\frac{1 + \sin \theta_C}{1 - \sin \theta_C}} + \sqrt[3]{\frac{1 - \sin \theta_C}{1 + \sin \theta_C}} - 1 \quad . \tag{9.4}$$

Substituting into Eq. (3.5) gives the second of three solutions to Eq. (2.1)

$$x = m \left(\sqrt[3]{\frac{1 + \sin \theta_C}{1 - \sin \theta_C}} + \sqrt[3]{\frac{1 - \sin \theta_C}{1 + \sin \theta_C}} \right) - 2m \quad .$$

$$(9.5)$$

Now, given the round numbers

$$\begin{cases} x = 1 - 1/30000 \\ n = 3 \end{cases}$$
 Key solution

Eq. (2.1) gives

$$\left. \begin{array}{l} \alpha \approx 1/Z \\ \theta_W \approx \theta_C \end{array} \right. \hspace{0.5cm} \left. \right\} \hspace{0.5cm} \text{Key approximations} \\$$

or, equivalently,

$$\left. \begin{array}{l} \alpha \approx 1/137.036\,000\,002 \\ \theta_W \approx 27.407\,157\,329^\circ \end{array} \right\} \ \ {\rm Key\ approximations} \label{eq:theta-weight}$$

so that $\sin^2\theta_C\approx 0.211\,886$. These values are again surprisingly close to those obtained from experiment. Moreover, the above θ_W and θ_C are similar in multiple ways: They are not only close numerically, but they are also similar in how they are defined. Each depends on ~ 137.036 for its value in the same way — with ~ 137.036 implicitly appearing as a value (in the form of electron charge) in the Weinberg angle's definition [8, 9]. In addition, when one works out the dimensional details of θ_W , Eq. (8.1) emerges naturally.

X. THIRD SOLUTION

The third solution to Eq. (2.1) is the most obvious. One merely rewrites Eq. (2.1) in the form of the general cubic equation, and then solves it using the general cubic's classical solution. The classical solution to

$$ax^3 + bx^2 + cx + d = 0 ag{10.1}$$

is

$$x = \sqrt[3]{\frac{q}{2} + \sqrt{\frac{q^2}{4} + \frac{p^3}{27}}} + \sqrt[3]{\frac{q}{2} - \sqrt{\frac{q^2}{4} + \frac{p^3}{27}}} - r \quad , \tag{10.2}$$

where

$$p = \frac{c}{a} - \frac{b^2}{3a^3}$$

$$q = -\frac{2b^3}{27a^3} + \frac{bc}{3a^2} - \frac{d}{a}$$

$$r = \frac{b}{3a} .$$
(10.3)

Though complicated, tellingly, these complexities largely vanish when a, b, c, and d derive from Eq. (2.1). So, when Eq. (2.1) is expanded into the general cubic equation we get these coefficients

$$\begin{cases}
 a = 1 & b = 6m \\
 c = 9m^2 & d = m(4m^2 - 3Z)
 \end{cases}$$
(10.4)

in terms of m and Z. Substituting the coefficients of Eq. (10.4) into Eq. (10.3) allows simplifying Eq. (10.2) to get

$$x = \sqrt[3]{t + \sqrt{t^2 - m^6}} + \sqrt[3]{t - \sqrt{t^2 - m^6}} - 2m$$
, (10.5)

where

$$t = m^3 - \frac{d}{2}$$

= $m (1.5Z - m^2)$. (10.6)

This is the third of the three solutions given for Eq. (2.1), a solution that is notably economical.

XI. SOLUTION ECONOMY

The complexity of the above three solutions can be objectively assessed by comparing them against the exceptionally-simple classical solution to the *depressed cubic*, a natural off-the-shelf benchmark. This will help clarify just how "basic" Eq. (2.1) really is.

The depressed cubic is merely the general cubic equation without its quadratic term (i.e., b=0). Assuming the coefficient of the depressed cubic's leading term equals one (i.e., a=1), we have

$$x^3 + cx + d = 0 . (11.1)$$

Substituting this equation's coefficients into Eq. (10.3) gives

$$p = c$$

$$q = -d$$

$$r = 0$$
(11.2)

so that Eq. (10.2) gives this compact solution to Eq. (11.1)

$$x = \sqrt[3]{\frac{-d}{2} + \sqrt{\frac{d^2}{4} + \frac{c^3}{27}}} + \sqrt[3]{\frac{-d}{2} - \sqrt{\frac{d^2}{4} + \frac{c^3}{27}}},$$
(11.3)

which dates back to Cardano's Renaissance masterwork Ars Magna. This solution is roughly as complicated as the earlier second and third solutions, showing that Eq. (2.1) and two of its solutions are sufficiently fundamental to be of interest to mathematicians. But perhaps Eq. (2.1) should also engage the interest of physicists, given its apparent connections to physical quantities?

XII. SUMMARY AND CONCLUSION

To help shed light on this, consider that it was the proximity of α to $1/137.\overline{037}$ that earlier suggested that Eq. (2.1) might produce interesting mathematics—as it does. In the same way, θ_W helped the author find the second solution to Eq. (2.1)—the one using θ_C —with θ_W providing the primary clue that such a solution existed. The precise way that α and $(M_W/M_Z)^2$ mapped over to the purely mathematical constants 1/137.036 and 108/137.036 appeared to require that θ_W have a mathematical correlate among the cubic equation's solutions—as it does. But why should these purely mathematical results obtained from solutions to a special case of the cubic relate to empirical constants such as α and θ_W ?

Yet another way to consider the above solutions is in their historical context: In the 16th century the Italian mathematicians Scipione del Ferro, Niccolò Tartaglia, and Gerolamo Cardano did pioneering work on the solution to the cubic equation. Had they probed more deeply, a congruous set of numbers fitting

- the fine structure constant
- the Weinberg angle
- $1/\sin^2 L12$
- $\sin^2 Q13$
- $\sin^2 L23 \times 1/\sin^2 Q12$
- $\sin^2 Q23 \times \sin^2 L13$

might have formed an integral part of their work from the outset, with these numbers waiting several centuries till their eventual re-discovery as part of $20^{\rm th}$ century physics.

^[1] J. S. Markovitch, "A rewriting system applied to the simplest algebraic identities" (2012) http://www.vixra.org/abs/1211.0029.

^[2] H. Kragh, "Magic Number: A Partial History of the Fine-Structure Constant," Archive for History of Exact Sciences 57:5:395 (July, 2003).

^[3] P. J. Mohr, B. N. Taylor, and D. B. Newell (2011), "The 2010 CODATA Recommended Values of the Fundamental Physical Constants" (Web Version 6.0). This database was developed by J. Baker, M. Douma, and S. Kotochigova. Available: http://physics.nist.gov/constants [Friday, 22-Jul-2011 10:04:27 EDT]. National Institute of Standards and Technology, Gaithersburg, MD 20899.

^[4] J. S. Markovitch, "A theorem producing the fine structure constant inverse and the quark and lepton mixing

angles" (2012) http://www.vixra.org/abs/1203.0106.

^[5] J. S. Markovitch, "Underlying symmetry among the quark and lepton mixing angles (Five year update)" (2012) http://www.vixra.org/abs/1211.0049.

^[6] J. S. Markovitch, "The cubic equation and 137.036" (2014) http://www.vixra.org/abs/1402.0050.

^[7] J. S. Markovitch, "Some mathematics inspired by 137.036" (2013) http://www.vixra.org/abs/1312.0016.

^[8] A. Pais, Inward Bound: of matter and forces in the physical world (Oxford University Press, 1995), pp. 212, 591–600.

^[9] G. Passarino and M. J. G. Veltman, "On the Definition of the Weak Mixing Angle," Phys. Lett. B 237, 537 (1990).

^[10] J. Beringer et al. [Particle Data Group Collaboration], Phys. Rev. D 86, 010001 (2012).