

# DYSON ANALYSIS OF GRAVITON PRODUCTION, LIGO, AND THE GERTSENSHTEIN EFFECT

Andrew Beckwith  
Chongqing University department of physics,  
Emails: [abeckwith@uh.edu](mailto:abeckwith@uh.edu), [rwill9955b@gmail.com](mailto:rwill9955b@gmail.com)  
Chongqing, PRC, 400044

## **Abstract**

In a 2013 paper, Freeman Dyson presented thought experiments challenging the detectability of gravitons via LIGO interferometry and via the Gertsenshtein effect. Dyson assumed a distance of several light years would be required for detection of the interaction between gravitational waves (GWs) and tenuous B fields and photons, making gravitons experimentally unverifiable. In this paper, we present contrary theoretical evidence for detectability of near-field interaction of gravitons, photons, and a magnetic field. Our first example of 100% probability of the Gertsenshtein effect working is due to a GW generated by a Tokamak with a interaction of GW, B field, and photons, in a volume on the order of a few cubic meters. We furthermore outline how Dyson was giving misleading information on the efficiency of LIGO, which is inimical to that research initiative on gravitational wave still very pertinent to essential gravitational physics, via the interferometer arrays, which has been noted and corrected

**Keywords:** Gertsenshtein effect; Tokamak; GW, LIGO

## 1 Introduction

First, we mention an error in Dyson's argument against LIGO, in which he incorrectly rendered the value of gravitational constant  $G$ , times 1 solar mass, divided by the speed of light, squared as equal to about  $10^{-33}$  centimeters. The correct value is 1.5 kilometers. LIGO is extremely important for gravitational wave physics, and we mention more of this in the end of the introduction. Also, Dyson in [7] derived criteria as to the probability one could obtain physical phenomenon theoretically modeled by the Gertsenshtein effect [8]. The Gertsenshtein effect [8] is the coupling of magnetic fields, gravitons, and photons. In the Dyson treatment [7] of the Gertsenshtein effect [8], Dyson hypothesized distances up to many light years for an interaction of magnetic fields, gravitons and photons, for experimental signals which could be detected on the Earth's surface. This assumed geometry of many light years distance lead to the predicted Gertsenshtein effect [8] unable to allow for graviton detection. In contrast to this assumed vast distances for the Gertsenshtein effect in reference [7], the author has devised via Tokamak generation of gravity waves[3], which lead to an interaction length of meters for the magnetic field, gravitons, and photons. The reduced length is due to the magnetic field which the gravitons interact with, being inside the detector itself, thereby insuring a 100 % probability for the Gertsenshtein effect occurring. This is commensurate with predictions given in reference [14]. The Tokamak example brings up an important point, that even if one wants to measure gravitational waves and detect gravitons from the early universe, that in the 3DSR model for GW detection [22], the Gertsenshtein effect for gravitons, magnetic field, and photons is within the small 3 dimensional geometry of the detector, with an enormous magnetic field. Having the Gertsenshtein effect [8] in such a small volume dramatically raises the

likelihood of detection of gravitons, via resultant photons being picked up by the 3DSR device.

It is extremely relevant and very important to keep in mind that interferometric based gravitational wave detectors like LIGO as well as its numerous updates in India and Europe are important also to set up the falsifiable experimental inputs needed to really test gravity theories, as it has been shown very strongly by Dr. Corda in 2009 [5] , and that Dr. Dyson has severely criticized a valuable addition to what is needed to understand gravitational physics. Noting Dr. Dyson's mistake about LIGO [7] is essential for vetting all needed operational tools needed to understand gravity. While Dyson is important, the record needs to be set straight as has been done in this manuscript. We start off with the LIGO example, since the Dyson analysis is unfair to instrumentation which is still very important to understanding gravitational waves, and by extension gravitons. That Dyson arguably has not treated LIGO with the seriousness which it deserves and a dimensional analysis as to the fault Dyson alleged as to the array is brought up, which applies incorrect use of dimensional analysis to argue against continual usage of instrumentation which would be needed as secondary back up to the vetting of the Graviton hypothesis, which will be needed later. I.e. if Gravitons are the particle component of the Gravity wave, as assumed in quantum mechanics, bunches of gravitons will be necessarily giving experimental traces which LIGO could pick out or at least indicate should be investigated via either LIGO or its presumed more refined successors. After this example is gone into, we will delve into the matter of the Gertscheshtein effect [7,8] which is serious.

## **2 Looking at the problem of LIGO , and reviewing Dyson's claims**

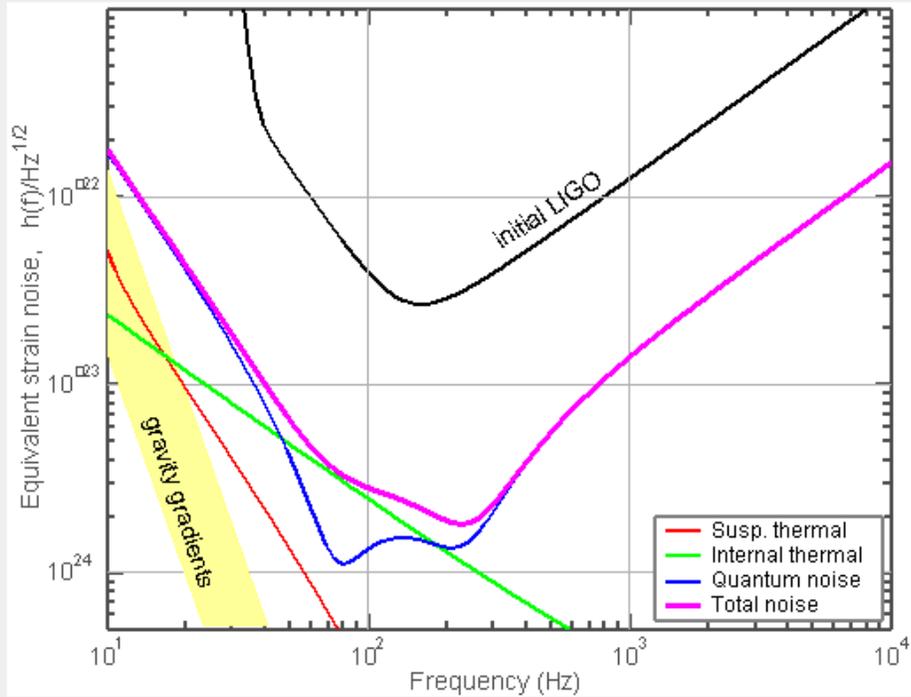
From [1] there is the following important claim, as to Gravitational wave detection which centers about how black holes, in collision may be essential for the identification of gravitational waves. We render the quote as follows, as being

a presumed primary source of Gravitational waves which may be useful in detection of gravitons, should gravitons exist.

Quote:

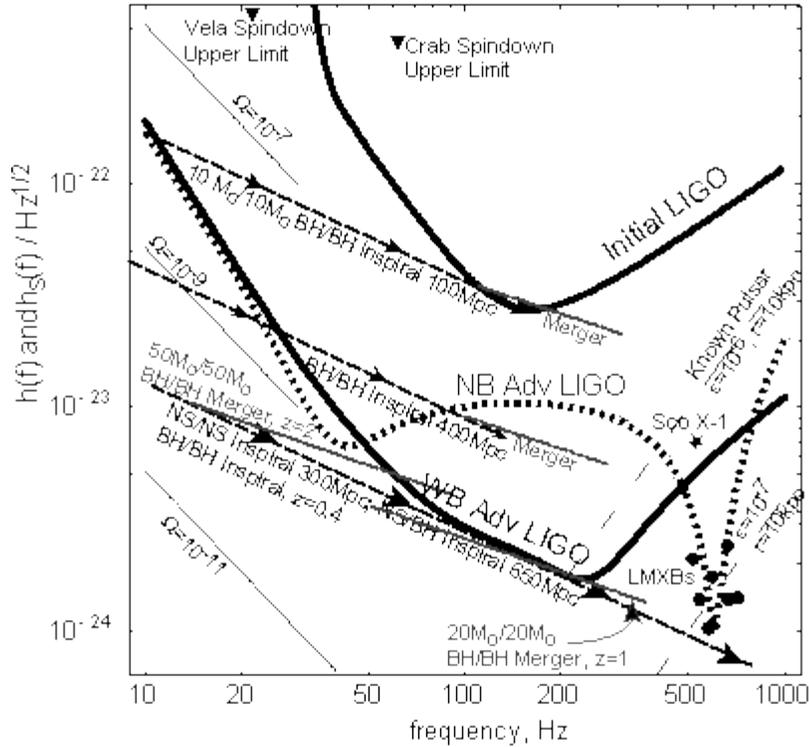
- *BH+BH mergers and ringdowns:* When rapidly spinning BH's collide, they should trigger large-amplitude, nonlinear oscillations of curved spacetime around their merging horizons. Little is known about the dynamics of spacetime under these extreme circumstances; we can learn about it by comparing LIGO's observations of the emitted waves with supercomputer simulations. Advanced LIGO can detect the merger waves from BH binaries with total mass as great as 2000 solar mass to cosmological redshifts as large as  $z=2$ .

In making this claim, in [1] the LIGO detector has frequency versus strain noise considerations. These are shown in reference [1] where there is the following diagram given in their document which we reproduce below as



**Figure 1** Noise Anatomy of Advanced LIGO. This model of the noise performance is based on the LIGO current requirements set, and represents the principal contributors of the noise and the least-squares sum of those components expressed as an equivalent gravitational wave strain.

Futhermore, [1] leads to the following descriptions of detectability, namely as given in its document the following diagram [1] which is reproduced below as



**Figure 2** The estimated signal strengths  $h(f)$  from various sources (thin lines, filled circles and star) compared with the noise  $h(f)$  (heavy lines) of three interferometers: initial LIGO, Advanced LIGO in a wideband (WB) mode, and Advanced LIGO narrowbanded (NB) at 600 Hz. See text for explanations of sources. The signal strength  $h_s(f)$  is defined in such a way that, wherever a signal point or curve lies above the interferometer's noise curve, the signal, coming from a random direction on the sky and with a random orientation, is detectable with a false alarm probability of less than one per cent using currently understood data analysis algorithms.

The signal strength of LIGO as given by [2] depends upon

$$h \sim \frac{GM}{c^2} \times \frac{1}{r} \times \left[ \frac{v}{c} \right]^2 \quad (1)$$

Here,  $r$  is the distance of this gravitational generation from the detector, and  $v/c$  is the ratio of say objects within the gravitational detector, and the speed of light. Usually,  $v/c$  is much less than 1, so (1) is particularly relevant to the problem of inspiraling black holes falling into each other, and so, now with this, we should review what Dyson had to say about gravitons, and GW, as well as LIGO.

Right before Dyson's [7], in his section 4, there is a statement that the frequency range for a single graviton to kick an electron out of a single atom, which is  $10^{15}$  Hertz [7]. We will later on comment this estimate [7] as a way to obtain a graviton-photon interaction and also refer to Dyson's claim just before his section 5, about thermal graviton generators, that the absorption cross section of ordinary matter (for a graviton) is  $10^{-41}$  square centimeters per gram. For LIGO, the frequency range is about  $10^2$  Hz for two black holes inspiraling into each other, not  $10^{15}$  Hertz, so the option of having a single graviton displace an electron from an atom, is zero. Which leads us to consider the relation given by Dyson, as his [7] formula (10), namely an upper bound to a minimum separation between two objects, say in a LIGO grid, is given by

$$\frac{GM}{c^2} > D \quad (2)$$

If  $M$  is the mass of the sun, then the L.H.S. of (2) is 1.482 times  $10^3$  meters, i.e. roughly 1.5 kilometers, or approximately a mile. Assume that then we wish to compare (1) with (2) with a value of  $V/c \sim 10^{-3}$ , we obtain that two inspiraling black holes with a strain value of  $h \sim 10^{-22}$  are about 1000 light years from Earth, for two black holes, combined mass of about one solar mass.

This example in itself, plus Dyson's mathematics should alert the reader, that Dyson, while undoubtedly brilliant in terms of his field theory work and research as up to the 1970s, is not parsing the problem of graviton detection correctly.

Having said this, the next step will be to review what could be done as far as looking at the early universe, as a source of GW, while moving beyond the mistakes we just outlined. In doing so, we assume that if our analysis is complete, we may be able to investigate early universe conditions, via considering if an improvement over the Gertsenshtein effect is possible. We then go to the matter of small object geometry

### **3 Probability for the Gertsenshtein effect, as described by Dyson for the Tokamak GW experiment**

We will briefly report upon Dyson's well written summary results, passing by necessity to the part on the likelihood of the Gertsenshtein effect occurring in a laboratory setting [7]. In doing so we put in specific limits as to frequency and the magnetic field, since in our work the objective will be to have at least theoretically a 100% chance of photon-graviton interaction [7] which is the heart of what Dyson reported in his research findings. What we find, is that with a frequency of about 10 to the 9<sup>th</sup> Hertz and a magnetic field of 10 to the 9<sup>th</sup> Gauss that there is nearly 100% chance of the Gertsenshtein effect being observed, within the confines of the Tokamak experiment as outlined in [3, 16].

In general relativity the metric  $g_{ab}(\mathbf{x}, t)$  is a set of numbers associated with each point which gives the distance to neighboring points. I.e. general relativity is a classical theory. By necessity, perturbations from flat Euclidian space, are usually configured as ripples in ‘flat space’, which are the imprint of gravitational waves in space-time. Our paper is to first of all give the probability of a pairing of photons to gravitons linkage, the Gertentshtein effect, as to how the signatures of a perturbation to the metric  $g_{ab}(\mathbf{x}, t)$  is linkable to photons and vice versa. The Gertentshtein effect is linked to how there is a linkage, signal wise, between gravitons and photons, and we are concerned as to what is a threshold as to insure that GW may be matched to the photons used by Dr. Li and others [7] to signify GW in a detector . To do so let us look at the Dyson criteria as a minimum threshold for the Gertentshtein effect happening [7], namely

$$D \cdot B^2 \cdot \omega \leq 10^{43} \quad (3)$$

The propagation distance is given by  $D$ , the magnetic field by  $B$ , and the frequency of gravitational radiation is given by  $\omega$ . We assume that the gravitational frequency is commensurate with the gravitational frequency of gravitons, i.e. that they are, averaged out one and the same thing. In doing so, making use of [7] we suppose on the basis of analysis that  $D$  is of the order of 10 to the 2<sup>nd</sup> power, since  $D$  is usually measured in centimeter, and by [7] we are thinking of about a 1 meter If  $B$  is of the order of 10 to the 9<sup>th</sup> Gauss Hertz, as deemed likely by [3] , then we have that if the GW frequency ,  $\omega$  is likewise about 10 to the 9<sup>th</sup> Hertz , that (1) is easy to satisfy. Note that if one has a vastly extended value for  $D$ , say 10 to the 13<sup>th</sup> centimeters that the inequality of (3) does not hold, so that by definition, as explained by Dyson that in a lot of cases, not relevant to [3] , that (3) is not valid, hence there would be no interexchange between gravitons and photons, and hence, if applied to the Dr. Li detector [13, 22] no way to measure gravitons by their photonic signature. Fortunately, as given by [3] this extended version of  $D$ , say 10 to the 13<sup>th</sup> centimeters does not hold.

And that then (1) holds. If so then, the probability of the Gertentshtein effect is presentable as, approximately,

$$P \leq (10^{36}/B^2 \cdot \omega^2) \propto 10^{36}/10^{18} \cdot 10^{18} \sim 1 \equiv 100\% \quad (4)$$

Summing up (4) is that the chosen values, namely if  $D$  is of the order of 10 to the 2<sup>nd</sup> power,  $B$  is of the order of 10 to the 9<sup>th</sup> power Gauss, and  $\omega$  is likewise about 10 to the 9<sup>th</sup> Hertz leads to approximately 100% chance of seeing Gertentshtein effects in the planned Tokamak experiment in [3]. In making this prediction as to (4), we can say that the left hand side, leading up to the evaluation of  $P$  with a numerator equal to 10 to the 36<sup>th</sup> power will be about unity for the values of  $B$  detector fields in Gauss ( magnetic field) or the generated gravitational field frequency  $\omega$  from the Tokamak, making an enormous magnetic field in the GW detector itself mandatory, which would necessitate a huge cryogenics effort, with commensurate machinery. Keep in mind that the GW detector is, as given in [3] about five meters above the Tokamak [3] , i.e. presumably the one in Hefei, PRC [16] .

Note, that , ironically, Dyson gets much smaller values of(2) than the above, by postulating GW frequency inputs as to the value of  $\omega$  about 10 to the 20<sup>th</sup> Hertz, i.e. our value of  $\omega$  is likewise about 10 to the 9<sup>th</sup> Hertz, much lower. If one has such a high frequency, as given by Dyson, the of course, (4) would then be close to zero for the probability of the Gertentshtein effect happening. I.e. our analysis indicates that a medium high GW frequency, presumably close to 10 to the 9<sup>th</sup> Hertz, and  $D$  10 to the 2<sup>nd</sup> power, presenting satisfaction of both (3) and (4). Note the main point though, for large values of  $D$ , (3) will not hold, making (4) not relevant, and that means in terms of the Dyson analysis, that far away objects generating gravitons will not be detectable. Via the Gertentshtein effect.

There is no such limitation due to a failure of (3) in the Tokamak GW generation setup [3] since then, for Tokamaks,  $D$  is very small. But if  $D$  is large in the case of a lot of astrophysical applications, then almost certainly one never gets to (4) since the Gertsenshtein effect is ruled out. We assume, next that refinements as to the Gertsenshtein effect are in the works, as given by [15,17,18] and next work out a protocol as to the next topic, i.e. early universe shift in space-time geometry leading to GW signals. We will briefly mention what the GW signals are, which are probably accessible if the Gertsenshtein effect is improved upon. Note we will review, briefly, what was given by Weinberg [21] as a black body analysis as to the feasibility of GW/ graviton production via an analysis similar to the black body radiation protocols, and show that the above mentioned figures as to GW/graviton production

#### **4 Brief review of graviton production for massless gravitons, using Weinberg Black body analogy**

From the book written by [21].For frequencies, between  $\omega$  and  $\omega+d\omega$ : the number of gravitons is given by Weinberg[21], page 287 Formula 10.89 as

$$n(\omega)d\omega = \frac{\omega^2}{\left[ \pi^2 \cdot \left( \exp \left[ \frac{\hbar \cdot \omega}{k_B T} \right] - 1 \right) \right]} \cdot d\omega \quad (5)$$

Integrate this, between two band widths of frequency for the graviton, or for a very narrow graviton frequency width  $\bar{\omega}$ , the following approximation is acceptable as a modification of (3) as from [21]

$$n(\bar{\omega}) \approx \frac{\bar{\omega}^2}{\left[ \pi^2 \cdot \left( \exp \left[ \frac{\hbar \cdot \bar{\omega}}{k_B T} \right] - 1 \right) \right]} \quad (6)$$

Note that  $k_B$  above is for the Boltzmann constant, and that T can be set by ANYTHING one wants to have it set by, and the upshot, is that for frequencies approximately as  $\bar{\omega}$  approximately of about  $10^9$  Hz, and with a temperature, of a Plasma as of about 100 KeV, then for (6) figure that one is going to have (4) per unit meter, cubed, in volume, which would lead to variations of easily  $10^4$  in magnitude from a baseline starting point of say 1 graviton per cubic centimeter, per second. This becomes important when comparing this graviton number, per cubic centimeter against the purported graviton flux number appearing in the case of the Earth given as a Graviton detector which appears in this paper. I.e. see (8) below. The contrast with (6) is stunning.

**5 Why the work by Dyson is not pertinent to long distance approximations as done in his manuscript if the main magnetic field for the Gertsenshtein effect occurs within a detector?**

On the face of it, the way the question as to if the Gertsenshtein effect [8] occurs outside a gravitational wave detector appears to be contrived. We assert this is not a contrived question, since the planned detector has a magnetic field many times stronger than what would be expected by conditions on the Earth surface, with Gertsenshtein effects occurring due to the Earth's comparatively very minor magnetic field not playing a role. As given by [8] there is a well defined physical process for graviton-magnetic field interactions which would lead to a photon

cascade, enough so, so that large D values, as given above to the tune of many kilometers in length are not advisable or necessary. Needless to say, if one does not believe that the Gertsenshtein effect is not mainly restricted within a GW detector, there are still serious problems with the Dyson formulation.

Review of (3) and (4) above come up with the datum that satisfying (3) is necessary for implementation of (4), i.e. (4) in full generality would likely read as [7]

$$P \sim \sin^2 \sqrt{(10^{36}/B^2 \cdot \omega^2)} \quad (7)$$

The main absurdity of this formulation is that usually, in interstellar space that one has low B field magnitudes, and low GW frequency values, i.e.  $\omega$  as low as 100 Hz. Or as high as  $\omega \sim 10^9 - 10^{10} \text{ Hz}$  i.e. in that sense, the Dyson examples chosen as of implementation of (3) and (4) go off the rails, with it being extraordinarily easy for enormous values of  $(10^{36}/B^2 \cdot \omega^2)$  in many situations. I.e. Dyson picked the values of B and also the picked value of  $\omega \sim 10^{20} \text{ Hz}$  is chosen for the purpose of making  $P \sim \sin^2 \sqrt{(10^{36}/B^2 \cdot \omega^2)} \propto 10^{36}/B^2 \cdot \omega^2 \ll 1$ , i.e. Dyson [7] cherry picked the numbers to make the probability for the Gertsenshtein effect as almost non existent, even if (1) were satisfied. But show me an example where one would have  $\omega \sim 10^{20} \text{ Hz}$  in interstellar space? This is important since  $\omega \sim 10^{20} \text{ Hz}$  is not feasible to entertain in most examples, and if one is looking at GW detectors, as has been done in [3] one is visualizing  $\omega \sim 10^9 - 10^{10} \text{ Hz}$  in the high end of the GW frequency values, as is given in the Tokamak example in Section II. I.e. Dyson's analysis [7] of

$P \sim \sin^2 \sqrt{(10^{36}/B^2 \cdot \omega^2)} \propto 10^{36}/B^2 \cdot \omega^2 \ll 1$  was arbitrarily picked to kill the possibility of a reading of the Gertsenshtein effect [8].

We close this section as to a review of [7] by stating there is a choice as to where the Gertsenshtein effect should occur in terms of space-time interactions for proper utilization of a Device physics analysis of where gravitons and B fields interact, and that the large D values Dyson in [7] postulates, are not relevant to the case where the Gertsenshtein effect occurs, mainly inside a GW detector. This concludes our analysis of Dyson's failure to properly set up the benchmarks as to analysis of where the Gertsenshtein effect really occurs. So then, we conclude with this statement, and then move to the deficiencies as to Dyson's assertion as to the Earth as a graviton detector, which is section 5 below.

## **6 Dyson's analysis of the Earth as a GW detector; incomplete physics and why**

We now review the particulars of Dyson's analysis of the Earth as a GW detector [7]. In doing so we are using the same numbers, and our break down of the results show that Dyson is making some assumptions here, which need to be seriously reviewed. In debt with the methodology of finding out what is germane in his analysis to research. To begin with, Dysons, formulae as given in reference [7] which Dyson in his reference calls formula (23) has a next flux of Gravitons hitting the surface of the Earth from the Sun

$$F(\text{flux}) = 4 \times 10^{-4} \text{ Gravitons per cm, squared, per second} \quad (8)$$

In this example, of (8) above, using Dyson's numbers, he claims that only 1 graviton out of  $10$  to the  $32^{\text{nd}}$  power of gravitons can be detected by the Earth's surface, assuming a graviton has about a kilovolt of energy i.e. this is, in its heart a situation where Dyson [7] is assuming an absorption cross section  $10$  to the minus  $41^{\text{st}}$  power per square centimeter per gram for the Earth, and an absurdly low collision rate. If this were true we are neglecting the Gertsenshtein interaction, since we are assuming no magnetic interface with incoming gravitons. This is only justifiable if there is a hard sphere collision between incoming 'gravitons' and ordinary matter. The analysis is incomplete and unnecessary since Dyson has set up a research meme where the Gertsenshtein [7,8] interaction regime stretching kilometers in duration with no fidelity as to the fact that the interaction space between gravitons and a magnetic field is within a GW detector, and does not stretch kilometers in duration away from the GW detector. Having said, that, there is an even more significant error as to Graviton detection and GW in the Dyson analysis of the LIGO device, which is to be brought up next.

## **7 Using the helpful part of the Dyson analysis**

What we have done is to ascertain that the Gertsenshtein interaction is valuable in near field device physics geometry. We have in Section 2, where the Dyson analysis can set and fix appropriate GW and graviton frequency values, and magnetic field values, so the Gertsenshtein interaction is certain to occur

Note that (7) and (8) theoretically could in themselves, if one assumed strain values of  $h \sim 10^{-30}$ , lead to very early universe detection. No one, however, posits that such sensitivity low values could be remotely detectable with conceived of, or extrapolated laser interferometer technology. Also, even in the matter of BHs, entropy speculations, leading to, that the ‘entropy’ of a BH is given by, where  $M$  is the mass of the BH,  $L_p$  Planck length, and  $A_{hor}$  is the area of the Event Horizon of a black hole., and we state the entropy as [10]

$$S = 4\pi M^2 = \frac{1}{4} \cdot \left( \frac{A_{hor}}{L_p^2} \right) \quad (9)$$

Here, in [10] we have that in its ( reference [10] ) formulae ( 24), that its main result is about the differential of the area of an event Horizon which is given as, if there is a Brane theory connection to the formation of BHs, with  $N$  the number of dimensions, say up to 10, that what is known as super-radiance , ie. bouncing of incoming radiation off the event horizon is a consequence, of the following derivation, namely if

$$dA_{hor} = \frac{8\pi r_H}{B} dM_{BH} \cdot \left( 1 - \frac{1}{\omega} \sum_{j=1}^{N/2} m_j \cdot \Omega_j \right) \quad (10)$$

If  $dM < 0$ , then the quantity  $\left( 1 - \frac{1}{\omega} \sum_{j=1}^{N/2} m_j \cdot \Omega_j \right) < 0$ , where the quantum numbers  $m_j > 0$  and  $\Omega_j = \frac{a_j}{a_j^2 + r_H^2}$  as frequency of BH arising due to the  $j$ th

component of BH angular Momentum  $J_j$  as correlated to event horizons of the BH . Such an analysis would have profound effects upon the Dyson analysis of the probability of Graviton detection, where the phenomenon of super-radiance could play a major role as far as GW and gravitons emitted by BHs, especially in the case of inspiraling black holes [10] collapsing upon each other.  $a = \sqrt{x^2 + y^2}$  can go to zero, and also  $r_H = M_{BH} - \sqrt{M_{BH}^2 - a^2}$  . Corresponding to BHs with, or without spin, which would affect GW and graviton production.

Having said, that we should examine what could happen if we have a refinement of the Gertsenshtein effect, and its aftermath. Especially as to early universe astronomy

### **8 Generalization to larger cosmological problems. i.e. what if refinements of the Gertsenshtein effect occur, and allow early universe GW astronomy?**

The simplest way to consider what may be involved in alterations of geometry is seen in the fact that in pre-Octonion space time regime (which is pre-Planckian), one would have (Crowell, 2005 [6] )

$$|x_j, x_i| \neq 0 \text{ under ANY circumstances, pre Planckian} \quad (11)$$

Whereas in the Octonion gravity space time regime where one would have(12) below hold that for enormous temperature increases (Crowell, 2005)[6]

$$[x_j, x_i] = i \cdot [\Theta_{ji}] \xrightarrow{Temp \rightarrow \infty} 0 \quad (12)$$

Here,

$$\Theta_{ji} \sim \Lambda_{NC}^{-2} \sim [\Lambda_{4-Dim}]^{-2} \propto 1/[T^{2\beta}] \xrightarrow{T \rightarrow \infty} 0 \quad (13)$$

Specifically (12) transformed to (13) will undergo physical geometry changes which show up in  $\delta_0$ . The space-time shift from pre Planck to the Planck epoch has gravity wave background radiation containing the imprint of the very earliest event. Next, is to consider what happens if Quantum (**Octonion** geometry) conditions hold. The supposition as given by in [12 ]

“Considering all these recent developments, it is plausible that quantum mechanics and gravity has information as a common ingredient, and information is the key to explain the strange connection between two”.

When quantum geometry holds, as seen by (14) , GW information is loaded into the octonion space time regime, and then transmitted to the present via relic GW which identified via the phase shift in GW as measured in a GW detector. This phase shift is  $\delta_0$ . The following flow chart is a bridge between the two regimes of (Crowell, 2005)[6] the case where the commutators for QM hold and then again to where the commutators for QM do not hold at all.

$$\begin{aligned} [x_j, p_i] &\neq -\beta \cdot (l_{Planck} / l) \cdot \hbar T_{ijk} x_k \\ \xrightarrow{\text{Transition-to-Planckian-regime}} [x_j, p_i] &= -\beta \cdot (l_{Planck} / l) \cdot \hbar T_{ijk} x_k \quad (14) \end{aligned} \quad (14)$$

(14) above represents the transition from pre-Planckian to Planckian geometry.

Also questions relating to how pre and post Planckian geometries evolve can be answered by a comparison of how entropy, in flat space geometry is linked with quantum mechanics (Lee, 2010)[12] . Once (14) happens, Beckwith hopes to look at the signals in phase shift  $\delta_0$

$$\left[ x_j, p_i \right] = -\beta \cdot (l_{Planck} / l) \cdot \hbar T_{ijk} x_k$$

$\xrightarrow{\text{Transition-to-release-of-relic-Gravitational-waves-in-flat-space}}$

*Planckian – Era – Generated – GW* (15)

Lee's paper [12] gives the details of information theory transfer of information from initially curved space geometry to flat space. When one gets to flat space, then, by (15) one then has a release of relic GW. The readers are referred to appendix A summarizing the relevant aspects of [12] ( Lee, 2010) in connecting space time geometry (initially curved space, of low initial degrees of freedom) to Rindler geometry for the flat space regime occurring when degrees of freedom approach a maxima, initially from  $t > 0s$  up to about  $t < 1s$  as outlined in an argument given below in (16). One of the primary results is reconciling the difference in degrees of freedom versus a discussion of dimensions. Also, as (16) occurs, there will be a buildup in the number of degrees of freedom, from a very low initial level to a higher one, as in the Gaussian mapping [4,13].

$$x_{i+1} = \exp[-\tilde{\alpha} \cdot x_i^2] + \tilde{\beta} \quad (16)$$

The feed in of temperature from a low level, to a higher level is in the pre Planckian to Planckian thermal energy input as by [4]

$$E_{thermal} \approx \frac{1}{2} k_B T_{temperature} \propto [\Omega_0 \tilde{T}] \sim \tilde{\beta} \quad (17)$$

(17) would have low numbers of degrees of freedom, with an eventual Gauss mapping up to 100 to 1000 degrees of freedom, as described by (Kolb and Turner, 1990)[ 11 ] .

It is important to note that the above proposed phase transition is speculative, but it could lead to another source of GW and maybe even Graviton production which with suitable analysis, would lead to more experimental opportunities for astrophysics investigations

Briefly put, this (17) could lead to the other development, namely that in research work as given by [15] (Li, and Yang, 2009), the following case for amplitude

$$A_{\otimes} = A_{\oplus} = \tilde{A} \quad (18)$$

Furthermore, first order perturbative terms of an E&M field have its components written as (Li, and Yang, 2009)[15]

$$\tilde{F}_{0_2}^{(1)} = i\tilde{F}_{0_1}^{(1)} \quad (19)$$

Secondly, there is a way to represent the” number” of transverse first order perturbative photon flux density as given in an earth bound high frequency GW detector [15] .

$$n_r^{(1)} = \frac{c}{2\mu_0\hbar\omega_e} \text{Re}\{ \} \quad (20)$$

$$\{ \} = i(\exp[-i\theta]) \cdot \tilde{F}_{0_1}^{(1)*} \cdot \left[ \frac{i}{\omega_e} \cdot \left( \frac{\partial\Psi_x}{\partial y} - \frac{\partial\Psi_y}{\partial x} \right) \right] \quad (21)$$

Here the quantity  $\frac{i}{\omega_e} \cdot \left( \frac{\partial\Psi_x}{\partial y} - \frac{\partial\Psi_y}{\partial x} \right)$  represents the z component of the magnetic field of a Gaussian beam used in an EM cavity to detect GW. We introduce the quantity Q, the quality factor of the detector cavity set up to observe GW, and  $\tilde{A}$ , the experimental GW amplitude. In the simplest case,  $\hat{B}_y^{(0)}$  is a static magnetic field. Then  $\tilde{F}_{0_2}^{(1)} = i\tilde{F}_{0_1}^{(1)}$  leads to by [15]

$$\tilde{F}_{0_1}^{(1)} = i2\tilde{A}\hat{B}_y^{(0)}Q \cdot \left[ \sin\left[ \frac{n\pi z}{b} \right] \right] \cdot \exp[i(-\omega_g t + \delta_0)] \quad (22)$$

The formula  $E_{thermal} \approx \frac{1}{2}k_B T_{temperature} \propto \tilde{\beta}$  [13] is a feed into  $\omega_g$  provided time  $t \propto$  Planck time, and set (22) with  $\omega_g \sim \omega_g$  by setting up  $E_{thermal} \approx \frac{1}{2}k_B T_{temperature} \approx \tilde{\beta}$ . In other words, for relic GW production, a interrelationship between  $\tilde{\alpha}$  and  $E_{thermal} \approx \frac{1}{2}k_B T_{temperature} \propto \tilde{\beta}$  for increases in degrees of freedom. This is a different perspective than what is normally used in analyzing what happens in a transition between initial Planck time  $\sim 10^{-44}$  seconds, and cosmological evolution up to  $10^{-30}$  seconds. The next discussion is on research done by [14], as to identifying traces of massive gravitons

## 9. Recasting the problem of detecting Gravitons in a detector for “massive” Gravitons

We now turn to the problem of detection. The following discussion is based upon with the work of Dr. Li, Dr. Beckwith, and other physics researchers in Chongqing University [14]. What (Li et al, 2003) have shown in 2003 [14] which Beckwith made an extension is to obtain a way to present first order perturbative electromagnetic power flux, i.e.  $T^{(i)uv}$  in terms of a non zero four dimensional graviton rest mass, in a detector, in the presence of uniform magnetic field [14]. What if we have curved spacetime with an energy momentum tensor of the electromagnetic fields in GW fields as given by [14] ?

$$T^{uv} = \frac{1}{\mu_0} \cdot \left[ -F_{\alpha}^{\mu} F^{\nu\alpha} + \frac{1}{4} \cdot g^{\mu\nu} F_{\alpha\beta} F^{\alpha\beta} \right] \quad (23)$$

By reference [14] we state that  $F_{\mu\nu} = F_{\mu\nu}^{(0)} + \tilde{F}_{\mu\nu}^{(1)}$ , with  $|\tilde{F}_{\mu\nu}^{(1)}| \ll |F_{\mu\nu}^{(0)}|$  will lead to

$$T^{uv} = T^{uv(0)} + T^{uv(1)} + T^{uv(2)} \quad (24)$$

The 1<sup>st</sup> term to the right side of (24) is the energy – momentum tensor of the back ground electro magnetic field, and the 2<sup>nd</sup> term to the right hand side of (24) is the first order perturbation of an electro magnetic field due to the presence of gravitational waves helps constrain (25) just below. This discussion as to section 8 is admittedly very preliminary, but it could be a way forward as to beginning to use the concept of a ‘current’ as in a GW/graviton detector, which with much more detail could take into account early universe phase transitions which occur at the beginning of the inflationary era. Secondly in conjunction with reference [20], it may remove problems associated with heavy gravity.

$$J_{effective} \cong n_{count} \cdot m_{4-D-Graviton} \quad (25)$$

As stated, [9,11]  $m_{4-D-Graviton} \sim 10^{-65} \text{ grams}$ , while  $n_{count}$  is the number of gravitons which may be in the detector sample. What Beckwith and Li intend to do is to isolate out an  $T^{uv(1)}$  assuming a non zero graviton rest mass. . I.e. use  $\tilde{\beta} \cong |F|$  and make a linkage with  $T^{00(1)}$ . The term  $T^{00(1)}$  isolated out from  $T^{uv(1)}$ . The point is that detected GW

## 10. Conclusion.

This paper raises questions as to the appropriateness of the Dyson analysis, in particular the Dyson dismissal of LIGO is based upon an incomplete rendering of a distance, D, as less than Planck Length, which we disprove by elementary analysis of the left hand side of (2) which with one solar mass is 1.48 kilometers, 1 mile, in value, as opposed to the Dyson sub Planck length. It is worth noting

that LIGO has kilometer long interferometer arms, and plenty of space, as to the obtaining GW and/or Graviton itself in instrumentations.

Dyson [7] also insisted upon evaluation of the Gertscheshtein effect in terms of light year distances as to light and magnetic field interactions, thereby concluding with virtually non existent Graviton interaction with instrumentation.

Next is an extensive discussion of the errors in the application of the Gertscheshtein interaction . Bluntly put, we think Dyson picked numbers for the magnitude of the interacting fields affected by the Gertscheshtein interaction, and did it probably to destroy fundamental research into gravitational physics. The author does not know why Dyson would be motivated to do such a reducto absurdum attack upon the feasibility of inquiry into foundations into Gravitational physics , but the resulting analysis in this text highlights some of the problems.

What we will bring up in closing is that the Gertscheshtein interaction is not necessarily the last word in effective graviton-magnetic field interactions and that improvements are in the offing which could enhance the role of GW detection. To do so, we can make an estimate that from a very simplistic viewpoint, that the view point of what is called the Li effect [13,15,22] involves a magnetic field of the same frequency, direction and appropriate phase of the gravitational wave field.

. For one thing, as given in the early part of the manuscript, what Dyson hypothesized for the probability of Gertscheshtein interaction for measurable gravitons as to a Tokamak generation of GW is appropriate and may be , for sufficiently large strain values of  $h \sim 10^{-25}$ , may be detected with advanced instrumentation. The problem is this. What Dyson postulates as to the probability of a Gertscheshtein interaction between Gravitons and a magnetic field is no issue in that situation. I.e. a very strong magnetic field would be inside the detector itself.

The Tokamak discussion is the opposite situation from the vast distances Dyson postulated photons traveled versus intervening galactic magnetic fields, as

then producing gravitons, is actually the reverse of the situation expected and modeled by Dr. Li and others [13 ,15, 22] I.e. the Gertenshtein effect is for within a DETECTOR device, and Dyson's calculations [7] as to light year distance of traveling of photons through magnetic fields is the reverse of the situation which was designed by the American and Chinese teams using 3DSR technology[22].

Dyson's analysis is in several specific cases not related to the actual situation of GW/ Graviton detection. As an example, Dyson states that  $10^{15}$  Hz for a graviton is required as to kicking an electron out of an atom [7], as though such a frequency is what would be expected of gravitons/GW. The fact is, that the Gertshenshtein effect does not need a frequency of  $10^{15}$  Hz due to GW / gravitons, to lead to detectable signals, in a detector.

At the same time, the directness of the questions asked by Dyson is welcome and the author acknowledges that until Dyson framed his article questions, that much of the GW/graviton issues were too incompletely rendered to permit an analysis of the relevant experimental issues.

### **Acknowledgement**

This work is supported in part by National Nature Science Foundation of China grant No. 11375279

### **References**

- [1] From the website maintained by David Shoemaker, Caltech, on Advanced Ligo [http://www.ligo.caltech.edu/advLIGO/scripts/ref\\_des.shtml](http://www.ligo.caltech.edu/advLIGO/scripts/ref_des.shtml)
- [2] From the website maintained by David Shoemaker , Caltech, on Gravitation <http://www.tapir.caltech.edu/~teviet/Waves/gwave.html>
- [3]Andrew Beckwith," Review of the Grischuk and Sachin Gravitational Wave Generator Via Tokamak Physics", <http://vixra.org/abs/1311.0132>
- [4]Andrew Beckwith, How to use the cosmological Schwinger principle for energy flux, entropy, and "atoms of space-time" to create a thermodynamic

- space–time and multiverse”. DICE 2010 (Pisa, Italy) contributed talk, 2011 *J. Phys.: Conf. Ser.* 306 012064 doi:10.1088/1742-6596/306/1/012064 ;  
<http://iopscience.iop.org/1742-6596/306/1/012064>
- [5] C. Corda "Interferometric detection of gravitational waves: the definitive test for General Relativity", *Int. J. Mod. Phys. D* 18, 2275 (2009)
- [6] L. Crowell, *Quantum Fluctuations of Space-time*, in World Scientific Series in Contemporary Chemical Physics, Volume 25, Singapore, Republic of Singapore, 2005
- [7] FREEMAN DYSON, *Int. J. Mod. Phys. A*, **28**, 1330041 (2013) [14 pages] DOI: 10.1142/S0217751X1330041X
- [8] M. E. Gertsenshtein, 1961. Wave Resonance of Light and Gravitational Waves", *JETP*, 41, 113-114, English translation in *Soviet Physics JETP*, 14, 84-85 (1962)
- [9] A. S. Goldhaber, M.M. Nieta "Photon and Graviton Mass Limits", <http://arxiv.org/pdf/0809.1003.pdf> Arxiv 0809.1003 V 5 [hep-th] 4 October 2013
- [10] E. Jung, S. Kim ,D.K. Park, "Conditions for the Super-radiance Mode in Higher Dimensional Rotating Black Holes with Angular Momentum Parameters", *Phys. Lett. B* 619, (2005), pp. 345-351, Arxiv:hep-th/0504139
- [11] E. Kolb, E. and S. Turner, *The Early Universe*, Westview Press, Chicago, USA, 1991
- [12] J.W. Lee, "Quantum Mechanics Emerges from Information Theory Applied to Causal Horizons", *Found Phys*, DOI: 10.1007/s10701-010-9514-3, pp 1-10 ( on line November 3, 2010)
- [13] F Y Li et al. *Eur. Phys. J. C* 56, 407-423 (2008)  
<http://www.gravwave.com/docs/Li-Baker%206-22-08.pdf>
- [14] F. Li, M. Tang, D. Shi. "Electromagnetic response of a Gaussian beam to high frequency relic gravitational waves in quintessential inflationary models", *PRD* 67, 104008 (2003), pp 1-17
- [15] F. Li and N. Yang, "Phase and Polarization State of High Frequency Gravitational waves", *Chin Phys. Lett.* Vol 236, No 5(2009), 050402, pp 1-4
- [16] H. Li., Y. Guo, et.al., "A long-pulse high-confinement plasma regime in the Experimental Advanced Superconducting Tokamak", *Nature Physics* (2013) doi:10.1038/nphys2795; Received 17 April 2013 Accepted 24 September 2013

Published online 17 November 2013

[17] M. Maggiore, *Gravitational Waves: Volume 1. Theory and Practice* Oxford University press, 2008, Oxford UK

[18] M. Meyers *Meyers – Perry Black holes*, pp 101-131 in *Black holes in higher dimensions* edited by G. Horowitz, Cambridge University Press, NYC, NY, USA, 2012

[19] T. Padmanabhan *Gravitation, Foundations and Frontiers* Cambridge University Press, NYC, NY, USA 2010

[20] P. Tinyakov, "Giving Mass to the Graviton", course 12, pp 471-498, of Session LXXXVI, *Particle Physics and Cosmology, the fabric of Space – Time*, editor F. Bernardeau, C. Grojean, J. Dalibard, 2006 ( Les Houches), Elsevier, Oxford, UK

[21] Steven Weinberg, *Gravitation and Cosmology, Principles and applications of the General theory of Relativity*, John Wiley Sons, 1972

[22] R. Clive Woods, Robert M L Baker, Fangyu Li, Gary V Stephenson, Eric W Davis, Andrew W Beckwith; A New Theoretical Technique for the Measurement of High-Frequency Relic Gravitational Waves, *Journal of Modern Physics*, SCIRP, DF PP. 498 - 518 DOI: 10.4236/jmp.2011.26060