

# The influence of the magnetic field in the Thomson Scattering

Herbert Weidner

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email: [herbertweidner@gmx.de](mailto:herbertweidner@gmx.de)

**Summary:** If an electromagnetic wave hits on electrons, so far only their reaction to the electric field component was calculated. By classical physics it can be shown that the electrons do not only perform the well-known transverse movement, but also a movement in the longitudinal direction at twice the frequency. This leads to an energy loss of the original electromagnetic wave without change of direction which was unknown up to now and has nothing to do with elastic or inelastic collision. The "lost" energy is radiated in two very different frequency ranges and the relative energy loss increases with decreasing frequency of the primary wave. An experimental confirmation of this phenomenon could influence the debate about "tired light".

Keywords: Thomson scattering, wave packet, transverse wave, tired light, polarization

## 1 Introduction

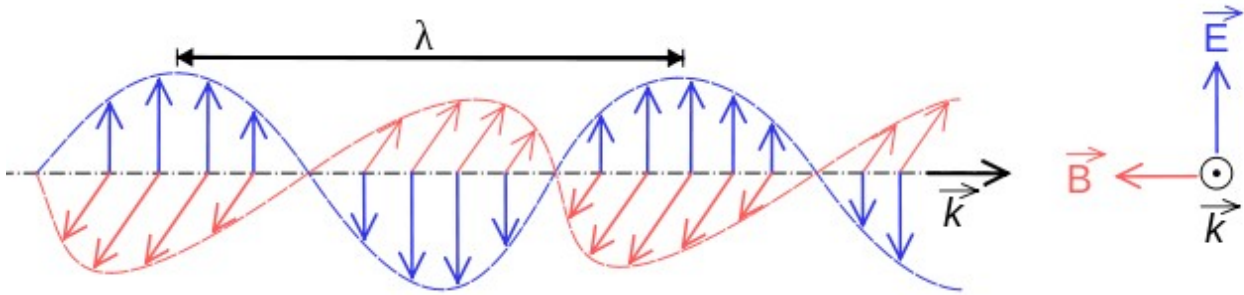
The usual explanation of the Thomson scattering is: "The electric field of the incident electromagnetic wave accelerates an unbound electron, causing it, in turn, to emit radiation at the same frequency. The magnetic field of the original waveform can be neglected." The effect of the Lorentz force is ignored, even though the electron moves in the variable magnetic field of the wave. This may not be the whole truth. Effects such as energy loss caused thereby are examined in the following work.

## 2 Basic assumptions of the far field

In the far field, an electromagnetic wave is defined by the direction of propagation, power density  $S$ , energy density  $w$  and frequency  $\omega$ . The values of the periodical electric and magnetic fields of the wave can be calculated with the Poynting formulae.

$E = \sqrt{Z_0 \cdot S}$  ,  $B = \mu_0 \sqrt{S / Z_0}$  and  $S = c \cdot w$  , where  $Z_0 = 377 \Omega$  is the characteristic impedance of free space. The maximum values are obtained by multiplying by a factor  $\sqrt{2}$  .

In the far field, the electric and magnetic field vectors  $\underline{E}$  and  $\underline{B}$  and the wave vector  $\underline{k}$  form a right-handed system, which is drawn in the picture for a linearly polarized wave.



The picture shows a short section of the wave train. Rightmost is a snapshot of the vectors  $\underline{E}$  and  $\underline{B}$ . The wave vector  $\underline{k}$  points to the observer.

If unbound, quiescent electrons are in the area of these fields, these are stimulated to oscillate and radiate energy. In the following two cases are discussed:

1. The electrons are part of an extended, thin plasma, wherein the mutual distance is considerably greater than the wavelength. The movement of the positive ions, which are also present, can be neglected because of their much bigger mass. However, they cause a resetting force on the electrons, so that they tend to oscillate at the plasma frequency. In the interstellar medium,  $\omega_p$  is very low with only a few kilohertz and is neglected below.
2. The wave can excite the electron shells of densely packed atoms to vibrate around their equilibrium positions. Even in transparent materials there are resonance effects which mostly lie in the UV region and are subsequently neglected. In solids, the distances between the "electron clouds" are significantly smaller than the wavelength, resulting in interference effects, which directs the radiation into certain directions. Those effects are well known in antenna technology.

### 3 Model of a wave packet

Every electromagnetic wave is limited in time, has a beginning and an end. Questions like coherence length, shape and volume of the continuous range in which the main part of the transferred energy is to be found are left to a follow-up article. Subsequently, only the consequences of the limited length of a wave packet are examined. Connections with "wave packets" of other branches of physics are not discussed.

First, an envelope is modeled, which is zero beyond a certain interval (Slowly varying envelope approximation). Some continuous functions  $C(t)$ , such as a Gaussian wave packet with cut slopes have unphysical discontinuities and are therefore less suitable. The frequently expressed assumption, the envelope is a decaying exponential function (damped wave) as one can hear after striking a bell, lacks any basis. In no atomic model one finds evidence of a lossy exchange of two forms of energy.

Regardless of the exact course, this function  $C(t)$  is a slow amplitude modulation using the modulation frequency  $W$ . Slowly means  $W \ll \omega$ . Each modulation generates so-called side band frequencies in the vicinity of  $\omega$  whose amplitudes typically decrease with increasing frequency separation. The side-band frequencies are  $\omega \pm W ; \omega \pm 2 W ; \omega \pm 3 W ; \dots$ , the entire frequency range is called *natural line width*.

Numerical verifications show that modifications of the shape of the envelope have little effect, they change the results of this work only marginally. Therefore, the well-known von-Hann-Window  $C=(1-\cos(W \cdot t))/2$  with the acceptable range  $0 \leq W \cdot t=2 \pi$  is chosen as the envelope of  $E$  and

$B$ . The short formula simplifies the following calculations because the energy of the wave is transferred in the immediate vicinity of only the three closely spaced spectral frequencies

$\omega; \omega + W; \omega - W$ . With this choice, the bandwidth is  $\Delta\omega \approx 2W$  and the line width is

$\Delta\lambda \approx \frac{4\pi W c}{\omega^2} = \frac{\lambda^2 W}{c\pi}$ . Only those wave trains will be considered, which include at least 100 oscillations, so it is assumed  $W \ll \omega$ . This assumption is likely to be met for most spectral lines. For example, the natural line width of the sodium D-line amounts to about 10 MHz and therefore the length of a wave packet is approximately  $10^7$  oscillation periods. The limiting case  $W \rightarrow 0$  describes an non-physical wave of constant amplitude and infinite extent and is not covered here.

As soon as a free, unbound electron gets in the sphere of influence of a wave packet, it is affected during the period of time  $0 \leq W \cdot t < 2\pi$ . Before and after the electron is at rest. The biggest part of the energy, about 92.44 %, is transported in the central region of the envelope between the time

$W t_1 = \pi/2$  and  $W t_2 = 3\pi/2$ . This interval is called the *coherence length*  $L \approx \frac{\pi c}{W}$  with the

above definitions. The extensions are energetically negligible. If the wave moves in the dispersion-free space, the coherence length is unchanged and there is no wave packet spreading.

The  $C(t)$  function describes only the dimension of the wave packet in parallel to the wave vector  $\underline{k}$ . The extent transversely thereto, that is the three-dimensional structure and the volume are treated in a separate article in relation to the Fresnel zones.

## 4 Movement of the electron in the y-direction

The fields of an electromagnetic wave exert a force on each charged particle. For the subsequent calculations at the place of the electron the following coordinates and functions are specified (see right part of the image above):

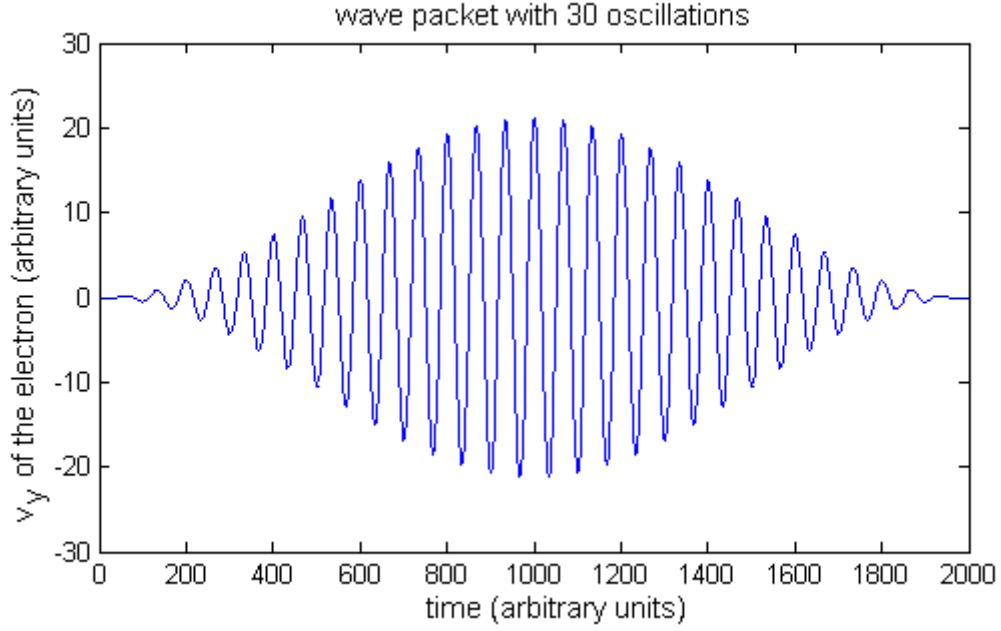
- The wave vector  $\underline{k}$  of an electromagnetic wave is parallel to the z axis and directed to the observer.
- The x-axis points to the right, the magnetic field at the location of the electron is  $B_x = -C(t) \cdot B_{max} \cdot \cos(\omega t) = \Re(-C(t) \cdot B_{max} \cdot e^{i\omega t})$
- The y-axis points upwards, the electric field at the location of the electron is  $E_y = C(t) \cdot E_{max} \cdot \cos(\omega t) = \Re(C(t) \cdot E_{max} \cdot e^{i\omega t})$

To simplify the calculations and to obtain compact results, these are mostly carried out using complex numbers. Only the real components are measurable.

A negatively charged electron is accelerated downward, and the speed follows from the formula  $a_y = q_e E_y / m_e$  by integration:

$$v_y = \frac{q_e E_{max}}{2 m_e \omega (\omega^2 - W^2)} \left( e^{i\omega t} (\omega W \sin(W t) + i \omega^2 \cos(W t) - i \omega^2 + i W^2) - i W^2 \right) \quad (1)$$

The integration constant is chosen such that at the start of the electromagnetic wave at the time  $W \cdot t = 0$  the condition  $v_y = 0$  is satisfied. The picture gives an idea of the time course of  $v_y$ .



After a further integration, the real component gives the y-coordinate of the electron

$$y = \frac{q_e E_{max}}{2 m_e \omega^2} (\cos(\omega t) (\cos(W t) - 1)) \quad (2)$$

This solution in the y-direction does not contain physical news, if the effects of the surrounding plasma and the movement of the electron in the z-direction are ignored. This is discussed further below.

## 5 Radiation due to y-acceleration

The electron is accelerated by the incoming wave packet (primary wave)

$$a_y = \frac{q_e E_{max}}{m_e} \frac{1 - \cos(W \cdot t)}{2} \cos(\omega t) \quad (3)$$

and emits energy. The value is obtained by calculating the power (Larmor formula) and subsequent integration over time during the duration of the wave packet  $0 \leq W \cdot t < 2\pi$ .

$$P_y = \frac{q_e^4 E_{max}^2}{24 \pi \epsilon_0 c^3 m_e^2} (1 - \cos(W \cdot t))^2 (\cos(\omega t))^2 \quad (4)$$

$$\int P_y dt = \frac{q_e^4 E_{max}^2}{24 \pi \epsilon_0 c^3 m_e^2} \frac{3 W t - 4 \sin(W t) + \cos(W t) \sin(W t)}{4 W} \quad (5)$$

$$A_y = \frac{q_e^4 E_{max}^2}{24 \pi \epsilon_0 c^3 m_e^2} \frac{3 \pi}{2 W} = \frac{q_e^4 \mu_0 c S_y}{24 \pi \epsilon_0 c^3 m_e^2} \frac{3 \pi}{2 W} = \frac{q_e^4 \mu_0^2 \cdot S_y}{8 m_e^2 W} \quad (6)$$

$S_y$  is the power density of the primary wave. This classical approximation is not valid for big power densities.

## 6 The cross section of the electron

The incident wave packet with the power density  $S_y$  is limited to the time period  $0 \leq W \cdot t < 2\pi$ . Since the envelopes of the electric and magnetic components are modelled with the function

$$C = \frac{1 - \cos(W \cdot t)}{2}, \text{ one gets } S_y(t) = S_0 \frac{(1 - \cos(W \cdot t))^2}{4}. \text{ The integration over the period of the}$$

wave packet results in the energy  $Q$ , which is distributed over the total cross-sectional area of the electron.

$$Q = \int_0^{2\pi/W} S_0 \frac{(1 - \cos(W \cdot t))^2}{4} dt = \frac{3\pi S_y}{4W} \quad (7)$$

A free electron can not store energy and has to re-radiate the amount it has previously absorbed from the wave package. Thus its actual absorption cross section  $\sigma$  can be calculated:

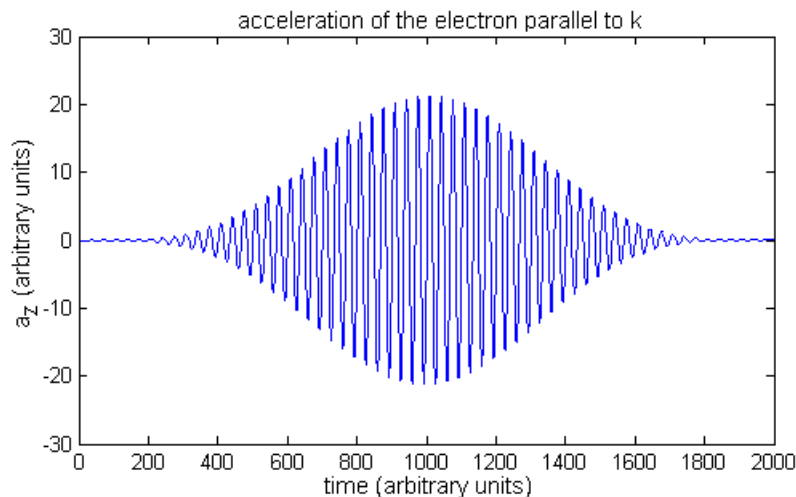
$$\sigma_{\text{Elektron}} = \frac{A_y}{Q} = \frac{q_e^4 \mu_0^2 S_y}{8m_e^2 W} \frac{4W}{3\pi S_y} = \frac{q_e^4 \mu_0^2}{6\pi m_e^2} = 6.625 \cdot 10^{-29} m^2$$

This result is identical with the classically calculated cross section of the electron.

## 7 Motion of the electron in the z-direction

The just described motion of the electron is in the y-direction. The magnetic field of the wave is oriented in anti-parallel to the x axis. The Lorentz force  $F_z = q_e v_y \times B_x$  ties together both and provides for the fact that the electron is accelerated in the direction of the positive z axis, so in parallel with  $\underline{k}$ . When using formula (1) results for

$$a_z = \frac{q_e^2 E_{\max} B_{\max} (1 - \cos(W \cdot t)) e^{i\omega t}}{-4m_e^2 \omega (\omega^2 - W^2)} \left( e^{i\omega t} (i\omega^2 \cos(Wt) - i\omega^2 + \omega W \sin(Wt) + iW^2) - iW^2 \right) \quad (8)$$



When compared with the preceding image you'll notice a temporal compression and the frequency doubling. Their cause is no inherent "non-linearity" of the electron, but the fact that in the far field of every electromagnetic wave the electric and magnetic fields reverse polarity at the same time. The integration of equation (8) provides the speed, whose real part is

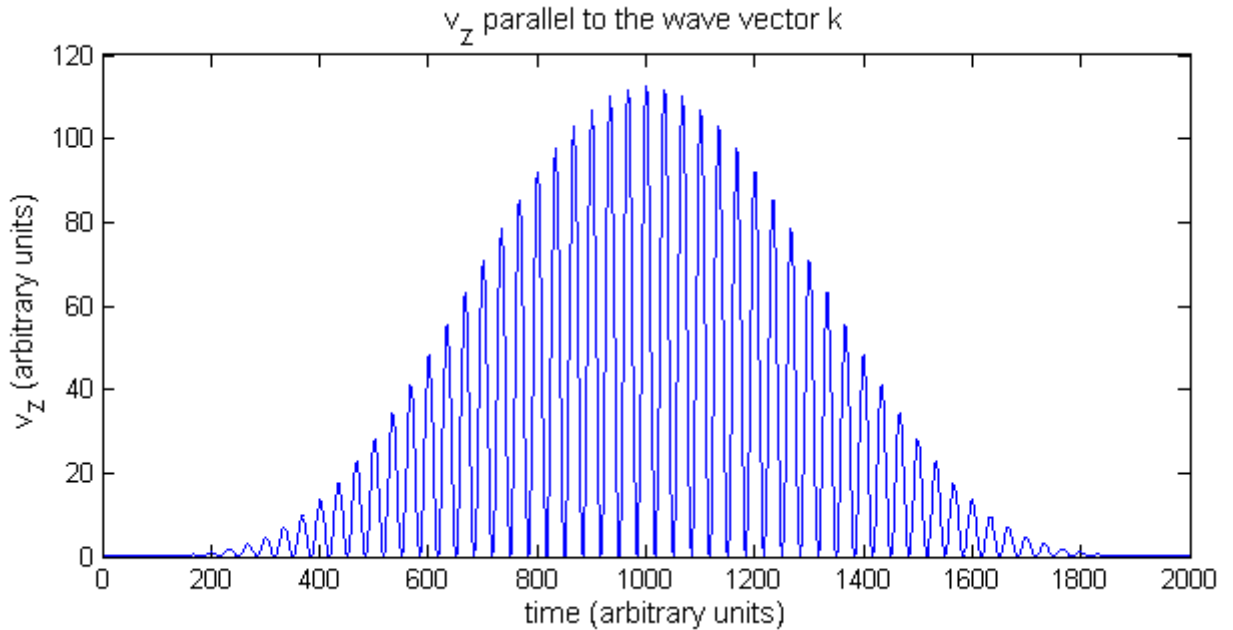
$$v_z = \frac{q_e^2 \mu_0 S_y}{m_e^2} \frac{R_1 \cdot R_2}{4 \omega^6} \quad \text{with} \quad (9)$$

$$R_1 = \cos(\omega t) \cos(W t) \omega^2 - W^2 - \cos(\omega t) (\sin(W t) W + \omega) \omega$$

$$R_2 = \cos(\omega t) \cos(W t) \omega^2 - W^2 + \cos(\omega t) (\sin(W t) W - \omega) \omega$$

The integration constant has been chosen so that at time  $t = 0$  the initial condition  $v_z = 0$  is satisfied. The following figure shows that the velocity  $v_z$  is never negative, and that therefore the unbound electron will only move away from the light source. The same result is reached by elementary numerical integration (see Appendix). It is noteworthy that the initial and final velocity are equal, so that the electron is not permanently accelerated. Therefore, this displacement of the electron consumes no energy and is not a consequence of the radiation pressure.

Evidently, the waveform shown in the image is the superposition of a cosine pulse and a wave packet with the frequency  $2\omega$ . Accordingly, the Fourier analysis of  $v_z$  provides two well-defined frequency ranges: A very broad band with a focus on  $W$  and a narrow band with the centre frequency  $2\omega$ . In both these frequency bands energy is emitted at the expense of the primary wave.



By integration over the period  $0 \leq W \cdot t < 2\pi$ , the distance travelled by the electron is calculated during the passage of the wave packet

$$z = \frac{q_e^2 \mu_0 S_y}{m_e^2} \frac{W^3 \pi}{2 \omega^6} \quad (10)$$

This displacement in z-direction can be observed only in unbound electrons, because their mass is sufficiently small. If the primary wave hits atoms and stimulates their electron shells to vibrate, the far higher total mass of the atom must be inserted into the formula. Therefore,  $z$  is reduced by at least

a factor of  $\left(\frac{m_{proton}}{m_{electron}}\right)^2 = 1836^2 \approx 3.4 \cdot 10^6$  and thus negligible.

The average velocity component of the electron in parallel with the wave vector  $\underline{k}$  and the z-axis is

$$v_z(\text{mean}) = \frac{q_e^2 \mu_0 S_y}{m_e^2} \frac{W^4}{4 \omega^6}$$

The wave packet loses no net energy by the displacement of the electron, because the electron subsequently is at rest. In the period  $0 \leq W \cdot t < \pi$  radiation energy is temporarily converted into kinetic energy and transforms back in the period  $\pi \leq W \cdot t < 2\pi$ . Overall, there is no momentum transfer from the wave packet to the electron.

These two consecutive accelerated movements cause the low frequency radiation, which is discussed below. The z-velocity obviously includes two different frequency components, which are discussed separately.

## 8 High-frequency emission due to the z-acceleration ( $2\omega$ )

The incoming wave packet accelerates the electron by the interaction of electric and magnetic field component in the z-direction at the frequency  $2\omega$ , in parallel with the wave vector  $\underline{k}$ . First, the power is calculated with the Larmor formula using equation (8)

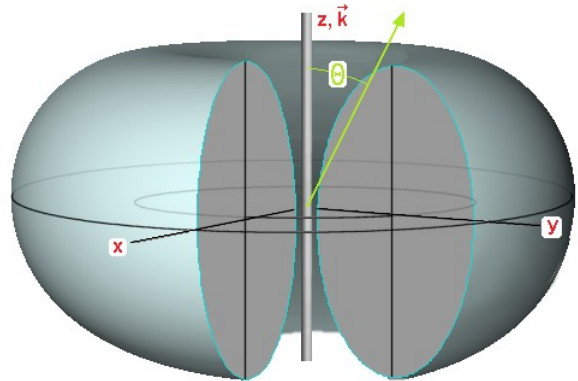
$$P_z = \frac{q_e |a_z|^2}{6 \pi \epsilon_0 c^3} = \frac{q_e^6 \mu_0^3 S_y^2}{24 \pi c m_e^4} \frac{(\sin(\omega t))^4 (\cos(Wt) - 1)^4}{\omega^2}$$

The integration over the period  $0 \leq W \cdot t < 2\pi$  results the value of the radiated energy:

$$A_z(\text{fast}) = \frac{q_e^6 \mu_0^3 S_y^2}{24 \pi c m_e^4 \omega^2} \frac{105 \pi}{32 W} = \frac{35 q_e^6 \mu_0^3 S_y^2}{256 c m_e^4 \omega^2 W} \quad (11)$$

It should be noted that the power  $P_z$  is generated and emitted during a very short interval.

The radiation pattern of the secondary radiation of a single electron is a torus, as we know from a dipole antenna. Whose axis is the  $\underline{k}$ -vector of the primary wave and whose maximum radiation lies in the xy-plane ( $\Theta \approx \pi/2$ ). The intensity of the secondary wave in the z-direction ( $\Theta = 0$ ), that is parallel to the  $\underline{k}$ -vector, is zero, because no longitudinal electromagnetic waves exist. Nothing is changed when the primary wave is circularly polarized.



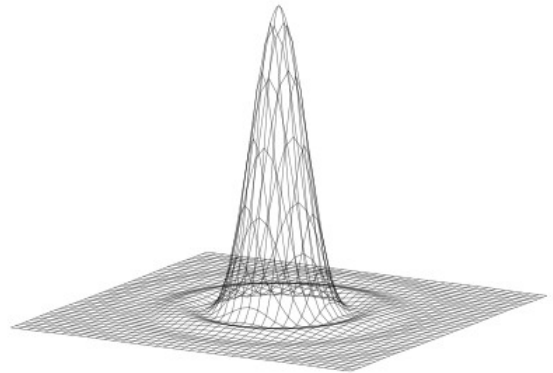
Because of the rotational symmetry of the torus, the radiation energy can *not* create a pulse component in the x-y-plane, and therefore also causes *no change* in the direction of the  $\underline{k}$ -vector. This could reignite the discussion of the hypothesis "tired light". Although the wave packet loses energy, there is no scattering in the physical sense, which is

of course always connected with a change in direction of the  $\mathbf{k}$ -vector. The energy loss of the primary wave described above produces *no image blur* - not for astronomical observations and not during irradiation of compact transparent material such as lenses. If only a few electrons vibrate, the energy loss will be very small and difficult to detect. Whether it can be ignored in any case, is not the subject of this physical derivation and deserves a separate investigation.

## 9 Radiation pattern with a small mutual spacing of the electrons

The torus is strongly distorted when the wave packet of the primary wave stimulates not only a single electron to oscillate, but many who are in close proximity. In solids, the mutual distance of the electrons and electron shells is much smaller than the wavelength, which simplifies the description.

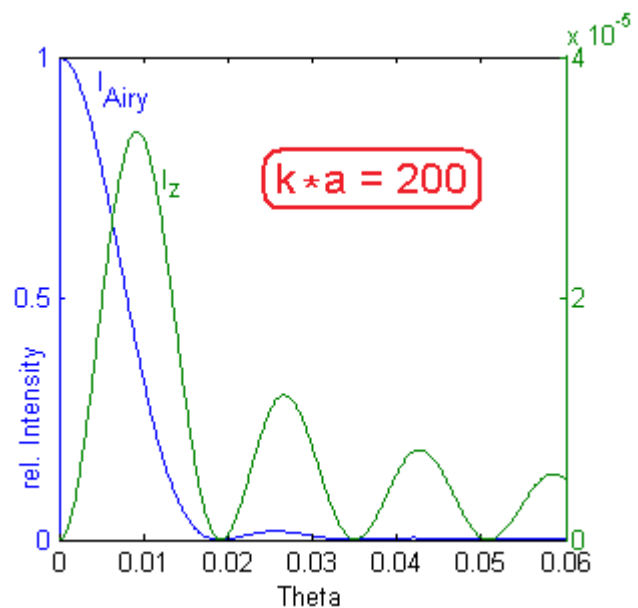
If the electrons form a circular area having a diameter  $2a$  significantly exceeding the wavelength and the polarization is parallel to the surface, the radiation focuses close to the surface normal ( $\Theta = 0$ ) and the intensity function is the known Airy disk, as shown in the image. The surface normal of the circle indicates the direction of maximum intensity. With increasing angular deviation from the normal, the intensity decreases rapidly and, in particular in the plane of the circular surface, the radiation disappears due to destructive interference.



Since the polarization of the secondary wave is parallel to the surface normal, any radiation in the direction of the normal is prohibited. Therefore, a compromise between the conflicting radiation patterns torus and Airy directivity pattern must be found.

If a glass cylinder of radius  $a$  is illuminated with light in parallel to the axis of symmetry, the plane wave fronts of the primary electromagnetic wave are parallel to the x-y-plane of the radiation. All electrons in any x-y-plane are forced to synchronous oscillations. As shown above, the velocity in conjunction with the magnetic field of the primary wave generates a very small movement of all electron clouds in the z-direction. Destructive interference ensures that in the x-y-plane ( $\Theta \approx \pi / 2$ ) no secondary wave is emitted.

The compromise is possible, if the radiation patterns torus and Airy disc are mathematically formulated and multiplied together:



a) Electrons oscillating in the z-direction do not radiate energy into the z-direction ( $\Theta = 0$ ) because longitudinal electromagnetic waves do not exist. For the radiated field strength, the formula is

$$E \sim \sin(\Theta)$$

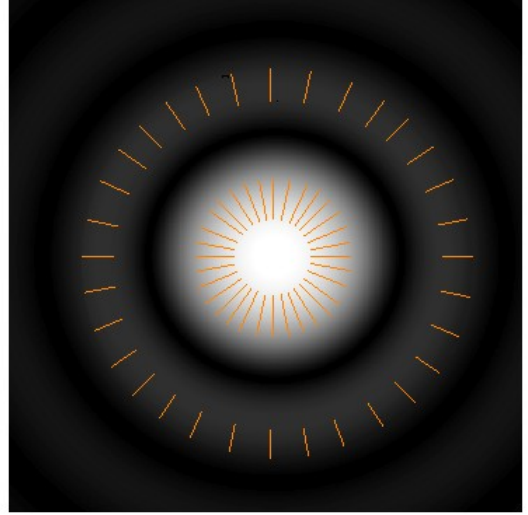


b) For constructive interference applies  $E \sim \frac{2J_1(k a \sin(\Theta))}{k a \sin(\Theta)}$  as for the diffraction at a circular aperture, where  $J_1$  is the Bessel function of the first kind and  $a$  is the radius of the glass cylinder.

The intensity of the radiation generated by the synchronous movement of the electrons in the z-direction, is the square of the product

$$I_z \sim (E_z)^2 \sim \left( \frac{2J_1(k a \sin(\Theta))}{k a} \right)^2 \quad (12)$$

As the picture shows, the radiation power of the secondary radiation disappears for  $\Theta = 0$  and has a maximum for  $1,842 = k a \sin(\Theta)$ . After the light has passed through the glass cylinder, an analysis should show, that the strong primary can be separated from the considerably weaker secondary wave due to different polarization and angle of observation.



The primary wave does not change polarization and frequency, the intensity profile corresponds after a sufficient distance to the well known Airy disk with a strong central peak, the angle of the first intensity minimum (innermost dark ring) can be calculated with the well-known formula  $(\sin \Theta)_{Airy} = 1,22 \frac{\lambda}{2a}$ . In the image,

the intensity of the primary wave is indicated by the area brightness.

The highest intensity of the secondary wave ( $2\omega$ ) is observed on concentric circular rings. For half the opening angle of the innermost circle, the formula is  $(\sin \Theta)_{2\omega} = 0,586 \frac{\lambda}{2a}$ . In the picture orange lines indicate two annular areas in which the secondary radiation should be measurable. The main difference, however, is to be expected in terms of polarization. The orientation of the lines shows the position-dependent linear polarization of the secondary wave. This is independent of the polarization of the primary wave, and upon whether it is linear or circular. Illustratively stated, the orange lines are short sections of the z-axis, seen from the particular view. In this section, the electron motion in parallel to the z-axis is discussed.

An analysis in the opposite direction of the primary wave can facilitate the detection of the secondary wave, because there the intensity of the reflected primary wave can be significantly reduced by a suitable coating of both ends of the cylinder.

## 10 Low-frequency radiation as a result of z-acceleration

In the z-direction, the electron is accelerated with two different time constants. The previous section dealt with the radiation of frequency  $2\omega$ , subsequent the low-frequency component is discussed, which can be observed only in unbound electrons. A suitable low-pass filtering of equation (9) yields the mean velocity of the electrons:

$$v_z(slow) = \frac{q_e^2 \mu_0 S_y (1 - \cos(W t))^2}{m_e^2 8 \omega^2} \quad (13)$$

and the acceleration

$$a_z(\text{slow}) = \frac{q_e^2 \mu_0 S_y}{m_e^2} \frac{(1 - \cos(Wt)) \sin(Wt)}{4\omega} \quad (14)$$

One first calculates the power with the Larmor formula

$$P_z(\text{slow}) = \frac{(q_e \cdot a_z(\text{slow}))^2}{6\pi \epsilon_0 c^3} = \frac{q_e^6 \mu_0^3 S_y^2}{96\pi c m_e^4} \frac{((\cos(Wt) - 1) \sin(Wt))^2}{\omega^2}$$

and then integrates over the period of the wave packet. The radiated energy is

$$A_z(\text{slow}) = \frac{q_e^6 \mu_0^3 S_y^2}{96\pi c m_e^4 \omega^2} \frac{5\pi}{4W} = \frac{5q_e^6 \mu_0^3 S_y^2}{384 c m_e^4 \omega^2 W} \quad (15)$$

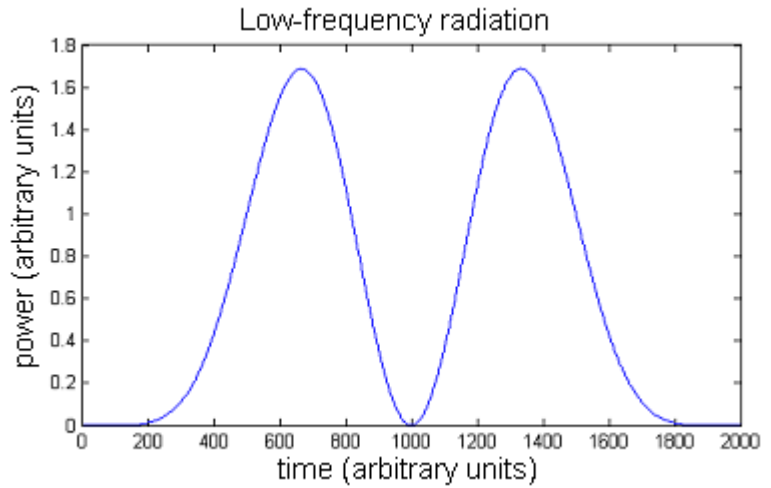
Since the movement of the electrons is considered in parallel to the  $\mathbf{k}$ -vector, the energy is mainly emitted in the x-y-plane.

The picture shows, that the low-frequency radiation is emitted not continuously during the period

$0 \leq W \cdot t < 2\pi$ , it focuses around the two time points

$$t_1 \approx \frac{2\pi}{3W} \quad \text{and} \quad t_2 \approx \frac{4\pi}{3W},$$

measured from the start of the primary wave. The difference depends on the contour of the envelope.



Since the electric field strength is proportional to  $a_z(\text{slow})$ , the mean radiated frequency corresponds to the variable factor  $\omega(\text{slow}) \approx (1 - \cos(Wt)) \sin(Wt)$ . Because the period of oscillation is hardly different from the duration of the wave packet, the frequency can be very imprecisely defined, the bandwidth of this secondary wave is very large and the signal is more like a noise. For the polarization, the same considerations apply as for the high-frequency component.

Example: From the wavelength  $\lambda_{\text{prim}} = 550 \text{ nm}$  and line width  $\Delta\lambda = 2 \text{ nm}$  follows

$W = 6,23 \cdot 10^{12} \text{ Hz}$ . If light passes through thin plasma, secondary radiation is a broadband spectrum in the frequency range  $f(\text{slow}) \approx 1,5 \cdot 10^{12} \text{ Hz}$  corresponding  $\lambda(\text{slow}) \approx 200 \mu\text{m}$ .

Interestingly, the cosmic background radiation is measured approximately in the same range. This radiation should be detectable perpendicular to the direction of propagation of the primary wave.

## 11 Specific energy loss as a result of z-acceleration

A free electron may not save energy and radiates all the energy, which it receives from the primary wave. Every time a wave packet excites a free electron to oscillate, it loses a tiny fraction of its energy, and then continues running with slightly reduced frequency. This extends the wavelength of

the primary wave by a tiny amount (redshift). For the three energy components, the following relationships apply:

$$A_y = \frac{q_e^4 \mu_0^2 \cdot S_y}{8 m_e^2 W} \quad (6)$$

$$\frac{A_z(\text{fast})}{A_z(\text{slow})} = \frac{21}{2} \quad A_z(\text{fast}) + A_z(\text{slow}) = \frac{115 q_e^6 \mu_0^3 S_y^2}{768 c m_e^4 \omega^2 W}$$

$$R = \frac{A_z(\text{fast}) + A_z(\text{slow})}{A_y} = \frac{115 q_e^6 \mu_0^3 S_y^2}{768 c m_e^4 \omega^2 W} \frac{8 m_e^2 W}{q_e^4 \mu_0^2 \cdot S_y} = \frac{115 q_e^2 \mu_0 S_y}{96 c m_e^2 \omega^2}$$

Of particular interest is the energy or power ratio  $R$ , when an electron is stimulated by the primary wave packet to execute short-term oscillations and therefore radiates secondary waves. The relative energy loss increases with decreasing frequency, and does not depend on the modulation frequency  $W$ , thus the duration of the wave packet. Because the energy difference is emitted with *no preferred direction*, the initial wave packet undergoes no change of direction, there is no scattering in the classic sense.

## 12 Concluding Remarks

The trigger for this investigation was a side note in [http://en.wikipedia.org/wiki/Thomson\\_scattering](http://en.wikipedia.org/wiki/Thomson_scattering) (version October 25, 2013). This is "... *the main cause of the acceleration of the particle will be due to the electric field component of the incident wave, and the magnetic field can be neglected.*"

Really? After only two hours programming work for a simulation using the Matlab program (see Appendix) amazing features of the electron motion in the z-direction showed up. They could not be removed by program changes (for example, better Integration procedures or smaller step size). The rest of the examination was - despite the use of symbolic algebra programs – the painstaking proof that the simulation is correct. Perhaps the side note will be adjusted to reality someday.

Although the results of this work may contribute to the discussion about interesting aspects of *tired light*, this was neither starting point nor target of the investigation.

## 13 Equipment:

Eugene Hecht: Optics

John David Jackson: Classical Electrodynamics

The Software packets „Mathematica“ , „Reduce“ and „Matlab“

OpenOffice Writer

perseverance and the infinite patience of my wife

## 14 Appendix

The following Matlab program is a simple simulation without any mathematical tricks. It generates some pictures to illustrate some results mentioned in this investigation. After copying in the Matlab editor window, it should be stopped after the selected rows (breakpoint) to look at the intermediate

results carefully. A change from the Hann Window to the Blackman Window shows that the shape of the envelope has little effect. It would be a great pleasure if someone would discover the correct form of the envelope.

```

%Thomson scattering; Author: Herbert Weidner
qe=-1.6e-19; me=9.11e-31; c=3e8; %basic units
cc=100; %points per oscillation
lambda=1e-6; %Meter
T=lambda/c; dt=T/cc;
S=1e6; %W/m² power density (Poynting)
%wave propagation z direction
Ey0=sqrt(377*S); %linear polarization; Z0=377 Ohm, points in y-direction
Bx0=12.566e-7*S/Ey0; %µ0, points in -x-direction
%z-direction to the observer, E x B = k (right-handed)
%a free electron is moved
c=30; %number of cycles
yz=zeros(2,c*cc); vyz=zeros(2,c*cc); %az=zeros(1,c*cc);
n=1:c*cc; %ma(y)=qE
%in the far field E and B are in-phase
ay=(hann(c*cc)).*sin(2*pi*n/cc)*qe*Ey0/me; %x-acceleration
Bx=-(hann(c*cc)).*sin(2*pi*n/cc)*Bx0; %- because EBk is right-handed
for n=2:c*cc-1 %compute vy and y
    vyz(1,n)=vyz(1,n-1)+dt*(ay(n-1)+ay(n))/2;
    yz(1,n)=yz(1,n-1)+dt*(vyz(1,n-1)+vyz(1,n))/2;
end
az=-qe*vyz(1,:).*Bx/me; %qv(y)B(x)=ma(z)
plot(vyz(1,:)); ylabel('y-velocity (m/s)');
title('wave packet, 30 oscillations'), xlabel('time')
%----Breakpoint in the following line -----
for n=2:c*cc-1 %compute vz and z
    vyz(2,n)=vyz(2,n-1)+dt*(az(1,n-1)+az(1,n))/2;
    yz(2,n)=yz(2,n-1)+dt*(vyz(2,n-1)+vyz(2,n))/2;
end
plot(vyz(2,:)); ylabel('z-velocity (m/s)');
title('velocity of the electron parallel to the k-Vektor')
xlabel('time (arbitrary units)')
n=1:size(vyz,2);
%----Breakpoint in the following line -----
plotyy(n,az,n,vyz(2,:));
ylabel('left: az, right: vz (m/s)');
title('a and v in z-direction'), xlabel('time')
%----Breakpoint in the following line -----
plotyy(n,vyz(1,:),n,vyz(2,:));
ylabel('left: vy, right: vz (m/s)');
title('velocity in y- and z-directions')
xlabel('time (arbitrary units)')
%----Breakpoint in the following line -----
plotyy(n,yz(1,:),n,yz(2,:));
ylabel('left: y, right: z in m');
title('coordinates of the electron in y- and z-direction')
xlabel('time (arbitrary units)')
%----Breakpoint in the following line -----
fprintf(1, 'vy/vz= %e\n',max(abs(vyz(2,:)))/max(vyz(1,:)))
%path of the electron in the y-z-plane
scatter(yz(1,:),yz(2,:),'.')
xlabel('deflection in y-direction (in m)')
ylabel('deflection in z-direction (in m)')
title('path of the electron in the y-z-plane')

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