Abstract

The “4/3 problem” of electrodynamics arose from the attempt to describe the mass of the electron as entirely electromagnetic in origin. Unfortunately, in the conventional treatment, there has been no acceptable solution offered. We present, here, a simple resolution to the ‘problem’ and show, by considering a previously overlooked part of the ‘electromagnetic’ field, that the mass of the electron is entirely ‘electromagnetic’ in origin.

1 Introduction

I would like to offer a resolution to the famous “4/3 problem” of electrodynamics. The theory of relativity implies that the momentum of the field of an electron must be the same as the rest energy of the field times $v/c^2$, where $v$ is the magnitude of the velocity of the electron. However, the momentum of the field, calculated from the Poynting vector, is $4/3$ times the energy of the field times $v/c^2$. Until now, this “4/3 problem” has gone unresolved (this non-relativistic, three-dimensional treatment is derived from my relativistic, four-dimensional equations.\footnote{See “New Transformation Equations and the Electric Field Four-Vector” at http://www.softcom.net/users/der555/newtransform.pdf}

I will show, here, that the mass of an electron is entirely ‘electromagnetic’ in origin. I’ll be using a derivation which closely parallels the one that Richard Feynman uses in “Lectures on Physics”, vol. 2, Sections 28-1 through 28-3. I’ve included, however, the necessary additions that Feynman didn’t consider. In order to conform with Feynman’s derivation, I’ve used SI units, here, in contrast to the Gaussian units used in the paper, above.

2 A Resolution to the Problem

The value that the mass $m_{\text{elec}}$, derived from the momentum of the field, must have to be considered entirely ‘electromagnetic’ in origin is the energy of the field, $U_{\text{elec}}$, divided by $c^2$, or

$$m_{\text{elec}} = \frac{U_{\text{elec}}}{c^2} \quad (1)$$
The value for $U_{\text{elec}}$, which Feynman calculated in Section 28-1, Eq. (28.2), is

$$U_{\text{elec}} = \frac{1}{2} \frac{e^2}{a}$$  \hspace{1cm} (2)

where $a$ is the lower limit of integration of the field energy density, and

$$e^2 = \frac{q^2}{4\pi \epsilon_0}$$  \hspace{1cm} (3)

where $q$ is the charge of the electron and $\epsilon_0$ is the permittivity constant. However, as I have shown elsewhere,\(^2\) the factor $1/2$ should not appear and the value for $U_{\text{elec}}$ should, instead, be

$$U_{\text{elec}} = \frac{e^2}{a}$$  \hspace{1cm} (4)

So in order to be considered entirely ‘electromagnetic’, in origin, our $m_{\text{elec}}$ needs to be

$$m_{\text{elec}} = \frac{e^2}{ac^2}$$  \hspace{1cm} (5)

Suppose that the electron is in uniform motion with velocity $v \ll c$. The momentum density of the field $g = \epsilon_0 E \times B$, where $E$ and $B$ are the conventional electric and magnetic field three-vectors, is directed obliquely to the line of motion for an arbitrary point $P$ at a distance $r$ from the center of the charge.\(^3\) The magnetic field is $B = v \times E/c^2$, which has the magnitude $(v/c^2)E \sin \theta$, where $\theta$ is the angle between $v$ and $E$. The momentum density $g$, then, has the magnitude

$$g = \frac{\epsilon_0 v}{c^2} E^2 \sin \theta$$  \hspace{1cm} (6)

The fields are symmetric about the line of motion, so when we integrate over space, the transverse components will sum to zero, giving a resultant momentum parallel to $v$. The component of $g$ in this direction is $g \sin \theta$ or, from (6)

$$g \sin \theta = \frac{\epsilon_0 v}{c^2} E^2 \sin^2 \theta$$  \hspace{1cm} (7)

However, the momentum due to $g \sin \theta$ alone, when integrated over all space and divided by $v$, as Feynman points out later, gives a value for $m_{\text{elec}}$ of

$$m_{\text{elec}} = \frac{4}{3} \frac{U_{\text{elec}}}{c^2}$$  \hspace{1cm} (8)

which is, clearly, not the same as the value from (1) required in order for the mass of the electron to be entirely ‘electromagnetic’ in origin. This is the “4/3 problem” problem.\(^4\)

I would like to consider, now, a contribution to the momentum density from $h = -\epsilon_0 E(\nabla \cdot A)$, where $A$ is the vector potential.\(^5\) Since $A = v \phi/c^2$, where $\phi$ is the static electric potential, we can also write $\nabla \cdot A$ as

$$\nabla \cdot A = \nabla \cdot \left( \frac{v \phi}{c^2} \right) = \frac{1}{c^2} v \cdot (\nabla \phi) = -\frac{1}{c^2} v \cdot E$$  \hspace{1cm} (9)

\(^2\)See “Energy Density Correction” at http://www.softcom.net/users/der555/enerdens.pdf

\(^3\)Refer to Fig. 28-1, page II-28-2 in Feynman’s ”Lectures on Physics”.

\(^4\)Actually, since the factor 1/2 in (2) has been removed, it should be called the ”2/3 problem”.

\(^5\)This is not an ad hoc addition. It is derived from the time component of my electric field four-vector, and is part of the momentum density components of my energy-momentum tensor in the paper in footnote 1.
The magnitude of $\mathbf{v} \cdot \mathbf{E}/c^2$ is $((v/c^2)E \cos \theta)$, so the magnitude of $\mathbf{h}$ is

$$h = \frac{\varepsilon_0 v}{c^2} E^2 \cos \theta \quad (10)$$

The component of $\mathbf{h}$ in the direction of $\mathbf{v}$ is $h \cos \theta$ or, from (10)

$$h \cos \theta = \frac{\varepsilon_0 v}{c^2} E^2 \cos^2 \theta \quad (11)$$

So the total momentum density is given by $g \sin \theta + h \cos \theta$. We now have to integrate the total momentum density over all space to find the total field momentum $p$.\(^6\) Feynman takes the volume element as $2\pi r^2 \sin \theta \, d\theta \, dr$, so that the total momentum is

$$p = \int_{\text{all space}} (g \sin \theta + h \cos \theta) 2\pi r^2 \sin \theta \, d\theta \, dr \quad (12)$$

or, using (7) and (11),

$$p = \int_{\text{all space}} \frac{\varepsilon_0 v}{c^2} E^2 (\sin^2 \theta + \cos^2 \theta) 2\pi r^2 \sin \theta \, d\theta \, dr \quad (13)$$

Since $\sin^2 \theta + \cos^2 \theta = 1$, this reduces to

$$p = \int_{\text{all space}} \frac{\varepsilon_0 v}{c^2} E^2 2\pi r^2 \sin \theta \, d\theta \, dr \quad (14)$$

The result of integrating this over all space, with the limits of $\theta$ being 0 and $\pi$, and the limits of $r$ being $a$ and $\infty$, is

$$p = \frac{e^2}{ac^2} v \quad (15)$$

Since $p = mv$, the mass of the electron $m_{\text{elec}}$ is

$$m_{\text{elec}} = \frac{e^2}{ac^2} \quad (16)$$

As you can see, this is exactly the same value as we got in (5). The value for the mass of the field derived from the energy of the field divided by $c^2$, and the value for the ‘electromagnetic’ mass derived from the momentum of the field are identical, meaning that the mass of the electron is entirely ‘electromagnetic’ in origin.

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\(^6\)Refer to Fig. 28-2, page II-28-2.