The Universal Quantum Fluid

Fran De Aquino
Maranhao State University, Physics Department, S.Luis/MA, Brazil.
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The quantization of gravity showed that the matter is also quantized, and that there is an elementary quantum of matter, indivisible, whose mass is \( \pm 3.9 \times 10^{-73} \text{kg} \). This means that any body is formed by a whole number of these particles (quantization). It is shown here that these elementary quanta of matter should fill all the space in the Universe forming a Quantum Fluid continuous and stationary. In addition, it is also explained why the Michelson-Morley experiment was not able to detect this Universal Quantum Fluid.

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1. Introduction

Until the end of the century XX, several attempts to quantize gravity were made. However, all of them resulted fruitless \([1, 2]\). In the beginning of this century, it was clearly noticed that there was something unsatisfactory about the whole notion of quantization and that the quantization process had many ambiguities. Then, a new approach has been proposed starting from the generalization of the action function\(^*\). The result has been the derivation of a theoretical background, which finally led to the so-called quantization of gravity and of matter \([3]\). The quantization of matter shows that there is an elementary quantum of matter whose mass is \( \pm 3.9 \times 10^{-73} \text{kg} \). This means that there are no particles in the Universe with masses smaller than this, and that any body is formed by a whole number of these particles. Here, it will be shown that these elementary quanta of matter should fill all the space in the Universe, forming a quantum fluid continuous and stationary. In addition, it is also explained why the Michelson-Morley experiment found no evidence of the existence of the universal fluid \([4]\). A modified Michelson-Morley experiment is proposed in order to observe the displacement of the interference bands.

2. The Universal Quantum Fluid

The quantization of gravity showed that the matter is also quantized, and that there is an elementary quantum of matter, indivisible, whose mass is \( \pm 3.9 \times 10^{-73} \text{kg} \)\[^{[3]}\].

Considering that the inertial mass of the Observable Universe is \( M_U = c^3/2H_0G \approx 10^{53} \text{kg} \), and that its volume is

\[
V_U = \frac{4}{3} \pi R_U^3 = \frac{4}{3} \pi \left(\frac{c}{H_0}\right)^3 \approx 10^{79} \text{m}^3,
\]

where \( H_0 = 1.75 \times 10^{-18} \text{s}^{-1} \) is the Hubble constant, we can conclude that the number of these particles in the Observable Universe is

\[
n_U = \frac{M_U}{m_0(\text{min})} \approx 10^{125} \text{particles}
\]

By dividing this number by \( V_U \), we get

\[
n_U/V_U \approx 10^{46} \text{particles/m}^3
\]

Obviously, the dimensions of the elementary quantum of matter depend on its state of compression. In free space, for example, its volume is \( V_U/n_U \). Consequently, its “radius” is \( R_U/\sqrt[3]{n_U} \approx 10^{-15} \text{m} \).

If \( N \) particles with diameter \( \phi \) fill all space of \( 1 \text{m}^3 \) then \( N \phi^3 = 1 \). Thus, if \( \phi \approx 10^{-15} \text{m} \) then the number of particles, with this diameter, necessary to fill all \( 1 \text{m}^3 \) is \( N \approx 10^{45} \text{particles} \). Since the number of elementary quantum of matter in the Universe is \( n_U/V_U \approx 10^{46} \text{particles/m}^3 \) we can conclude that these particles fill all space in the Universe, forming a Quantum Fluid continuous and stationary, the density of which is

\[
\rho_{\text{CUF}} = \frac{n_U m_0(\text{min})}{V_U} \approx 10^{-27} \text{kg/m}^3
\]

Note that this density is smaller than the
density of the Intergalactic Medium \( \rho_{\text{IGM}} \approx 10^{-26} \text{ kg} / \text{m}^3 \).

The density of the Universal Quantum Fluid is clearly not uniform along the Universe, since it can be strongly compressed in several regions (galaxies, stars, blackholes, planets, etc). At the normal state (free space), the mentioned fluid is invisible. However, at supercompressed state, it can become visible by giving origin to the known matter, since matter, as we have seen, is quantized and consequently, formed by an integer number of elementary quantum of matter with mass \( m_{0(\text{min})} \). Inside the proton, for example, there are

\[
n_p = m_p / m_{0(\text{min})} \approx 10^{25}
\]

elementary quanta of matter at supercompressed state, with volume \( V_{\text{proton}} = \frac{m_p}{\sqrt{n_p}} \approx 10^{-30} \text{m}^3 \).

Therefore, the solidification of the matter is just a transitory state of this Universal Quantum Fluid, which can turn back into the primitive state when the cohesion conditions disappear.

Due to the cohesion state of the elementary quanta of matter in the Universal Quantum Fluid, any amount of linear momentum transferred to any elementary quantum of matter propagates totally to the neighboring and so on, in such way that, during the propagation of the momentum, the elementary quanta of matter do not move, in the same way as the intermediate spheres in Newton’s pendulum (the well-known device that demonstrates conservation of momentum and energy) \[5, 6\]. Thus, whether it is a photon that transfers its momentum to the elementary quanta of matter, then the momentum variation due to the incident photon is \( \Delta p = h/\lambda \), where \( \lambda \) is its wavelength. As we have seen, the diameter of the elementary quantum of matter is \( \Delta x \approx 10^{-15} \text{m} \). According to the Uncertainty Principle the variation \( \Delta p \) can only be detected if \( \Delta p \Delta x \geq h \). In order to satisfy this condition we must have \( \lambda \leq 2\pi \Delta x \approx 10^{-14} \text{m} \). This means that momentum variations, in the elementary quanta of matter, caused by photons with wavelength \( \lambda > 10^{-14} \text{m} \) cannot be detected. That is to say that the propagation of these photons through the Universal Quantum Fluid is equivalent to its propagation in the free space. In practice, it works as if there was not the Universal Quantum Fluid. This conclusion is highly important, because it can easily explain why in the historical Michelson-Morley experiment there was no displacement of the interference bands namely because the wavelength of the light used in the Michelson-Morley experiment was \( \lambda = 5 \times 10^{-7} \text{m} \) fact that led Michelson to conclude that the hypothesis of a stationary ether was incorrect. Posteriorly, several experiments \[7-13\] have been carried out in order to check the Michelson-Morley experiment, but the results basically were the same obtained by Michelson.

Thus, actually there was no displacement of the interference bands in the Michelson-Morley experiment because the wavelength used in the experiment was \( \lambda = 5 \times 10^{-7} \text{m} \), which is a value clearly much greater than \( 10^{-14} \text{m} \), and therefore, does not satisfy the condition \( \lambda \leq 2\pi \Delta x \approx 10^{-14} \text{m} \) derived from the Uncertainty Principle. The substitution of light used in the Michelson-Morley experiment by radiation with \( \lambda \leq 10^{-14} \text{m} \) is clearly impracticable. However, the Michelson-Morley experiment can be partially modified so as to yield the displacement of the interference bands. The idea is based on the generalized expression for the momentum obtained recently\[3\], which is given by

\[
\rho = M_g V
\]

where \( M_g = m_g / \sqrt{1 - V^2 / c^2} \) is the relativistic gravitational mass of the particle and \( V \) its velocity; \( m_g = \chi m_0 \) the general expression of the correlation between the gravitational and inertial mass; \( \chi \) is the correlation factor\[3\]. Thus, we can write

\[
\frac{m_g}{\sqrt{1 - V^2 / c^2}} = \frac{\chi m_0}{\sqrt{1 - V^2 / c^2}}
\]

Therefore, we get
\[ M_g = \chi M_i \]  
(6)

The Relativistic Mechanics tells us that
\[ p = \frac{U V}{c^2} \]  
(7)

where \( U \) is the total energy of the particle. This expression is valid for any velocity \( V \) of the particle, including \( V = c \).

By comparing Eq. (7) with Eq. (4) we obtain
\[ U = M_g c^2 \]  
(8)

It is a well-known experimental fact that
\[ M_i c^2 = hf \]  
(9)

Therefore, by substituting Eq. (9) and Eq. (6) into Eq. (4), gives
\[ p = \frac{\sqrt{\frac{V}{c}} \frac{h}{\lambda}} \]  
(10)

Note that this expression is valid for any velocity \( V \) of the particle. In the particular case of \( V = c \), it reduces to
\[ p = \frac{\sqrt{\frac{h}{\lambda}}}{\lambda} \]  
(11)

By comparing Eq. (10) with Eq. (7), we obtain
\[ U = \chi hf \]  
(12)

Note that only for \( \chi = 1 \) Eq. (11) and Eq. (12) are reduced to the well-known expressions of DeBroglie \( (q = \frac{h}{\lambda}) \) and Einstein \( (U = hf) \).

Equations (10) and (12) show, for example, that any real particle (material particles, real photons, etc) that penetrates a region (with density \( \rho \), conductivity \( \sigma \) and relative permeability \( \mu_r \)), where there is an electromagnetic field \( (E, B) \), will have its momentum \( p \) and its energy \( U \) reduced by the factor \( \chi \), given by[3]:
\[ \chi = \frac{m_k}{m_0} = \left\{ 1 - 2 \left[ 1 + \left( \begin{array}{c} \frac{\Delta p}{m_k c} \end{array} \right)^2 \right] \right\} = \left\{ 1 - 2 \left[ 1 + 1.758 \times 10^{-2} \left( \frac{\mu_0 \sigma^2}{\rho_f^2 \beta^3} \right) B_{rms}^2 \right] \right\} \]  
(13)

where \( B_{rms} \) is the rms value of the magnetic field \( B \).

The remaining amount of momentum and energy, respectively given by
\[ (1 - \chi) \left( \frac{V}{c} \right) \frac{h}{\lambda} \]  
and \[ (1 - \chi) hf \], are transferred to the imaginary particle associated to the real particle\(^\dagger\) (material particles or real photons) that penetrated the mentioned region.

It was previously shown that, when the gravitational mass of a particle is reduced to a range between \( +0.159 M_i \) to \( -0.159 M_i \), i.e., when \( \chi < 0.159 \), it becomes imaginary\[3\], i.e., the gravitational and the inertial masses of the particle becomes imaginary. Consequently, the particle disappears from our ordinary space-time. It goes to the Imaginary Universe. On the other hand, when the gravitational mass of the particle becomes greater than \( +0.159 M_i \), or less than \( -0.159 M_i \), i.e., when \( \chi > 0.159 \), the particle return to our Universe.

Figure 1 (a) clarifies the phenomenon of reduction of the momentum for \( \chi > 0.159 \), and Figure 1 (b) shows the effect in the case of \( \chi < 0.159 \). In this case, the particles become imaginary and, consequently, they go to the imaginary space-time when they penetrate the electric field \( E \). However, the electric field \( E \) stays in the real space-time. Consequently, the particles return immediately to the real space-time in order to return soon after to the imaginary space-time, due to the action of the electric field \( E \). Since the particles are moving at a direction, they appear and disappear while they are crossing the region, up to collide with the plate (See Fig.1) with a momentum, \( p_* = \chi \left( \frac{V}{c} \right) \frac{h}{\lambda} \), in the case of a material particle, and \( p_* = \chi \frac{h}{\lambda} \), in the case of a photon.

If this photon transfers its momentum to elementary quanta of matter \( (\Delta x \approx 10^{-15} m) \), then the momentum variation due to the incident photon is \( \Delta P = \chi \frac{h}{\lambda} \). According to the Uncertainty Principle the variation \( \Delta p \) can only be detected if \( \Delta p \Delta x \approx h \), i.e., if
\[ \lambda \leq 2\pi \chi \Delta x \]  
(14)

We conclude, then, that the interaction between the light used in the Michelson-
Morley experiment \((\lambda = 5 \times 10^{-7} \text{m})\) and the Universal Quantum Fluid just can be detected, and to produce of the displacement of the interference bands, if
\[
\chi \geq 8 \times 10^7 \quad (15)
\]
In order to satisfy this condition in the Michelson-Morley experiment, we must modify the medium where the experiment is performed (for example substituting the air by low-pressure Mercury plasma), and apply through it an electromagnetic field with frequency \(f\). Under these conditions, according to Eq. (13), the value of \(\chi\) will be given by
\[
\chi = \left\lvert 1 - \left(1 + 1.758 \times 10^{-27} \left( \frac{\mu_0 \sigma_3}{\rho^2 f^3} \right) e^4 B_{\text{rms}}^4 - 1 \right) \right\rvert \quad (16)
\]
If the low-pressure Mercury plasma is at \(P = 6 \times 10^3 \text{Torr} = 0.8 \text{Nm}^{-2}\) and \(T \approx 318.15 \text{K}\) [14], then the mass density, according to the well-known Equation of State, is
\[
\rho = \frac{P M_0}{R} \approx 6.067 \times 10^{-5} \text{kg.m}^{-3} \quad (17)
\]
where \(M_0 = 0.2006 \text{g.mol}^{-1}\) is the molecular mass of the Hg; \(Z \approx 1\) is the compressibility factor for the Hg plasma; \(R = 8.314 \text{joule.mol}^{-1} \cdot \text{K}^{-1}\) is the gases universal constant.

The electrical conductivity of the Hg plasma, under the mentioned conditions, has already been calculated [15], and is given by
\[
\sigma \approx 3.419 \text{ S.m}^{-1} \quad (18)
\]
By substitution of the values of \(\rho\) and \(\sigma\) into Eq. (16) yields
\[
|\chi| = \left\lvert 1 - 2 \left[ 1 + 1.547 \times 10^{17} \frac{B_{\text{rms}}^4}{f^3} - 1 \right] \right\rvert \quad (19)
\]
By comparing with (15), we get
\[
\frac{B_{\text{rms}}^4}{f^3} \geq 0.01037 \quad (20)
\]
Thus, for \(f = 1 \text{Hz}\), the ELF magnetic field must have the following intensity:
\[
B_{\text{rms}} \geq 0.32 T \quad (21)
\]
This means that, if in the Michelson-Morley experiment the air is substituted by Hg plasma at \(6 \times 10^3 \text{Torr}\) and \(318.15 \text{K}\), and an ELF magnetic field with frequency \(f = 1 \text{Hz}\) and intensity \(B_{\text{rms}} \geq 0.32 T\) is applied through this plasma (Fig. 2), then the displacement of the interference bands should appear.

It is important to note that due to the Gravitational Shielding effect [3], the gravity above the magnetic field is given by \(g \geq 7.8 \times 10^7 \text{m.s}^{-2}\). This value, extends above the vacuum chamber for approximately 10 times its length. In order to eliminate this problem we can replace the ELF magnetic field, \(B_1\), shown in Fig. 2, by two ELF magnetic fields, \(B_1\) and \(B_2\), sharing the same frequency, \(f = 1 \text{Hz}\). The field, \(B_1\), is placed vertically through the region of the experimental set-up. The field, \(B_2\), is also placed vertically, just above \(B_1\) (See Fig. 3).

Thus, the gravity above \(B_2\) is given by \(\chi_1 \chi_2 g\) where \(\chi_1 = m_{g1}/m_{i1}\) and \(\chi_2 = m_{g2}/m_{i2}\) are respectively, the correlation factors in the Gravitational Shieldings 1 and 2, produced by the ELF magnetic fields \(B_1\) and \(B_2\), respectively. In order to become \(\chi_1 \chi_2 g = g\) we must make \(\chi_2 = 1/\chi_1 \approx 1/8 \times 10^7\). According to Eq. (19), this value can be obtained if \(B_{\text{rms}(2)} = 5.331481522 \times 10^{-5} T\) and \(B_{\text{rms}(1)} = 0.32 T\). Note that the value of \(B_{\text{rms}(2)}\) is less than the value of the Earth’s magnetic field \((B_0 \approx 6 \times 10^{-5} T)\). However, this is not a problem because the steel of the vacuum chamber works as a magnetic shielding, isolating the magnetic fields inside the vacuum chamber.

\[
B_{\text{rms}} \geq 0.32 T
\]
There are a type of neutrino, called "ghost" neutrino, predicted by General Relativity, with zero mass and zero momentum. In spite its momentum be zero, it is known that there are wave functions that describe these neutrinos and that prove that really they exist.

Fig. 1 – The correlation factor in the expression of the Momentum (a) Shows the momentum for $\chi > 0.159$. (b) Shows the effect when $\chi < 0.159$. Note that in both cases, the material particles collide with the cowl with the momentum $p_m = \chi \left( \frac{V}{c} \right) \left( \frac{h}{\lambda} \right)$, and the photons with $p_r = \chi (h/\lambda)$. 

* There are a type of neutrino, called “ghost” neutrino, predicted by General Relativity, with zero mass and zero momentum. In spite its momentum be zero, it is known that there are wave functions that describe these neutrinos and that prove that really they exist.
Fig. 2 - The modified Michelson-Morley experiment. The air is substituted by Hg plasma at $6 \times 10^3 \text{Torr}$ and 318.15 K, and an ELF magnetic field with frequency $f = 1 \text{Hz}$ and intensity $B_{\text{rms}} \geq 0.32 T$ is applied through this plasma, then the displacement of the interference bands should appear.
Fig. 3 – Cross-section of the vacuum chamber showing the magnetic fields $B_1$ and $B_2$.

$$B_1 = \mu_0 \left( \frac{N_1}{L_1} \right)_1; \quad B_2 = \mu_0 \left( \frac{N_2}{L_2} \right)_2$$
References


