

Fine Structure Constant $\alpha \sim 1/137.036$ and Blackbody Radiation Constant $\alpha_R \sim 1/157.555$

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Abstract

The fine structure constant $\alpha = e^2/\hbar c \approx 1/137.036$ and the blackbody radiation constant $\alpha_R = e^2(a_R/k_B^4)^{1/3} \approx 1/157.555$ are linked by prime numbers. The blackbody radiation constant is a new method to measure the fine structure constant. It also links the fine structure constant to the Boltzmann constant.

Planck and Einstein noted respectively in 1905 and 1909 that $e^2/c \sim h$ have the same order and dimension.[1, 2] This was before the introduction of the fine structure constant $\alpha = e^2/\hbar c$ by Sommerfeld in 1916.[3] Therefore, the search for a mathematical relationship between $e^2/c \sim h$ was started from the blackbody radiation. The Stefan-Boltzmann law states that the radiative flux density or irradiance is $J = \sigma T^4$ [erg · cm⁻² · s⁻¹] in CGS units or [W · m⁻²] in SI units. From the Planck law, the Stefan-Boltzmann constant $\sigma = 5.670400(40) \times 10^{-5}$ [erg · cm⁻²K⁻⁴s⁻¹] is

$$\begin{aligned} \sigma &= 2\pi \int_0^\infty \frac{x^3 dx}{e^x - 1} \frac{ck_B^4}{(hc)^3} = \frac{2\pi^5}{15} \frac{ck_B^4}{(hc)^3} \\ &= 2\pi\Gamma(4)\zeta(4) \frac{ck_B^4}{(hc)^3} = \frac{4^2\pi^5}{5!} \frac{ck_B^4}{(hc)^3} \end{aligned} \quad (1)$$

The Stefan-Boltzmann law can be expressed as the volume energy density of the blackbody $\varepsilon_T = \mathbf{a}_R T^4$ [erg · cm⁻³], where the radiation density constant \mathbf{a}_R is linked to the Stefan-Boltzmann constant

$$\mathbf{a}_R = \frac{4\sigma}{c} = \frac{4^3\pi^5}{5!} \frac{k_B^4}{(hc)^3} \quad (2)$$

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In 1914, Lewis and Adams noticed that the dimension of the radiation density constant divided by the 4th power Boltzmann constant \mathbf{a}_R/k_B^4 is (energy \times length)⁻³, while \mathbf{e}^2 is (energy \times length). However, they obtained an incorrect result equivalent to $\alpha^{-1} = \hbar c/\mathbf{e}^2 = 32\pi (\pi^5/5!)^{1/3} = 137.348$. [4] In 1915, Allen rewrote it as $\alpha = \mathbf{e}^2/\hbar c = (15/\pi^2)^{1/3}/(4\pi)^2$. [5]

In CGS units, $\mathbf{e}^2 = (4.80320427(12) \times 10^{-10})^2$ [erg·cm], $\mathbf{a}_R = 7.56576738 \times 10^{-15}$ [erg·cm⁻³K⁻⁴], and $k_B^4 = (1.3806504(24) \times 10^{-16})^4$ [erg⁴K⁻⁴]. We get the experimental dimensionless constant [6]

$$\begin{aligned}\alpha_R &= \mathbf{e}^2 \left(\frac{\mathbf{a}_R}{k_B^4} \right)^{1/3} = \frac{1}{\mathbf{157.5548787}} \\ &= \mathbf{0.00634699482}\end{aligned}\quad (3)$$

This is the blackbody radiation constant α_R , which is on the same order of the fine structure constant $\alpha = \mathbf{e}^2/\hbar c$

$$\begin{aligned}\alpha_R &= \frac{2}{\pi} \left(\frac{\pi^5}{5!} \right)^{1/3} \alpha = \left(\frac{\pi^2}{15} \right)^{1/3} \alpha \\ &= \left(\frac{\Gamma(4)\zeta(4)}{\pi^2} \right)^{1/3} = 0.8697668 \cdot \alpha\end{aligned}\quad (4)$$

Therefore, $\alpha_R \neq \alpha$, both α and α_R are experimental results incapable of producing the α math formula. Physically, the fine structure constant α is obtained from the atomic *discrete* spectra, while the blackbody radiation constant α_R is obtained from the thermal radiation of a 3D cavity in the *continuous* spectra. However, their relationship can be given by the Riemann zeta-function or by the modification of Euler's product formula (1737)

$$\begin{aligned}\frac{\alpha_R^3}{\alpha^3} &= \frac{\pi^2}{15} = \frac{\zeta(4)}{\zeta(2)} = \frac{\sum_{n=1}^{\infty} \frac{1}{n^4}}{\sum_{n=1}^{\infty} \frac{1}{n^2}} = \frac{\prod_p \left(1 - \frac{1}{p^2}\right)}{\prod_p \left(1 - \frac{1}{p^4}\right)} = \prod_p \left(\frac{p^2}{p^2 + 1} \right) \\ &= \frac{2^2}{(2^2 + 1)} \frac{3^2}{(3^2 + 1)} \frac{5^2}{(5^2 + 1)} \frac{7^2}{(7^2 + 1)} \cdots = \frac{4}{5} \cdot \frac{9}{10} \cdot \frac{25}{26} \cdot \frac{49}{50} \cdot \frac{121}{122} \cdots\end{aligned}\quad (5)$$

where the Euler product extends over all the *prime* numbers. In other words, the fine structure constant and the blackbody radiation constant can be linked by the prime numbers. From (5), the Stefan-Boltzmann law as the volume energy density of the blackbody ε_T is related to the fine structure constant α and the oscillating charge \mathbf{e}^2 with the different resonating frequencies in a cavity

$$\begin{aligned}
\varepsilon_T &= \mathbf{a}_R T^4 = \frac{4\sigma}{c} T^4 = \frac{4^3 \pi^5}{5!} \frac{k_B^4}{(hc)^3} T^4 \\
&= \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 k_B^4 T^4 = \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4
\end{aligned} \tag{6}$$

and the radiative flux density is

$$J = \sigma T^4 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 \frac{c}{4} k_B^4 T^4 = \frac{c}{4} \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4 \tag{7}$$

and the total brightness of a blackbody is

$$B = \frac{J}{\pi} = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 \frac{c}{4\pi} k_B^4 T^4 = \frac{c}{4\pi} \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4 \tag{8}$$

and the inner wall pressure of the blackbody cavity is

$$P = \frac{4\sigma}{3c} T^4 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 \frac{1}{3} k_B^4 T^4 = \frac{1}{3} \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4 \tag{9}$$

According to the Bose-Einstein model of photon-gas,[7] the free energy of the thermodynamics is

$$F = -PV = -\frac{4\sigma}{3c} VT^4 = -\frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 \frac{V}{3} k_B^4 T^4 = -\frac{V}{3} \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4 \tag{10}$$

and the total radiation energy is

$$E = -3F = 3PV = \frac{4\sigma}{c} VT^4 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 V k_B^4 T^4 = V \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^4 \tag{11}$$

where the photon gas $E = 3PV$ is the same as the extreme relativistic electron gas, and the entropy is

$$S = -\frac{\partial F}{\partial T} = \frac{16\sigma}{3c} VT^3 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 \frac{4V}{3} k_B^4 T^3 = \frac{4V}{3} \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^3 \tag{12}$$

and the specific heat of the radiation is

$$C_V = \left(\frac{\partial E}{\partial T} \right)_V = \frac{16\sigma}{c} VT^3 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{\mathbf{e}^2} \right)^3 4V k_B^4 T^3 = 4V \left(\frac{\alpha_R}{\mathbf{e}^2} \right)^3 k_B^4 T^3 \tag{13}$$

Landau assumed that the volume V in (10)~(13) must be sufficiently large in order to change from *discrete* to *continuous* spectra. Experimentally, solids or dense-gas have the continuous spectra, and hot low-density gas emits the discrete atomic spectra. The pattern of Planck spectra is given by $f(x) = x^3/(e^x - 1)$

where photon $h\nu$ is hidden in $x = h\nu/k_B T$. The *photon* integral in (1) is equal to a dimensionless constant (Fig. 1)

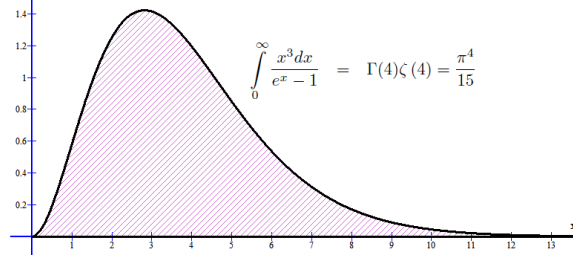


Figure 1: Photon integral is a dimensionless number $\Gamma(4)\zeta(4) = \frac{\pi^4}{15} = 6.4939394$

$$\begin{aligned} \int_0^{\infty} \frac{x^3 dx}{e^x - 1} &= \Gamma(4)\zeta(4) = \frac{\pi^4}{15} = 2 \cdot 3 \cdot \prod_p \left(\frac{p^4}{p^4 - 1} \right) \quad (14) \\ &= 2 \cdot 3 \cdot \frac{2^4}{(2^4 - 1)} \frac{3^4}{(3^4 - 1)} \frac{5^4}{(5^4 - 1)} \frac{7^4}{(7^4 - 1)} \cdots \\ &= \frac{3^3 5^2}{2^3 \cdot 13} \cdot \frac{7^4}{(7^4 - 1)} \frac{11^4}{(11^4 - 1)} \cdots = \frac{3^3 5^2}{2^3 \cdot 13} \prod_{p>5} \left(\frac{p^4}{p^4 - 1} \right) \end{aligned}$$

where the Euler product extends over all the *prime* numbers. The photon distribution integral (14) yields a zeta-function that is linked to the Euler prime products, therefore, there is no $h\nu$ in (6)~(13). (14) shows clearly how the fine structure constant α for the discrete spectra in (4) is converted to the blackbody radiation constant α_R for the continuous spectra by multiplying a dimensionless constant. (5) and (14) indicate that this dimensionless constant can be expressed as the Euler infinite prime number product.

In (6)~(13), the oscillators of the thermal electrons $\alpha/e^2 = 1/\hbar c$ or α_R/e^2 play a critical role in the electromagnetic coupling on a 3D surface (Fig. 2).

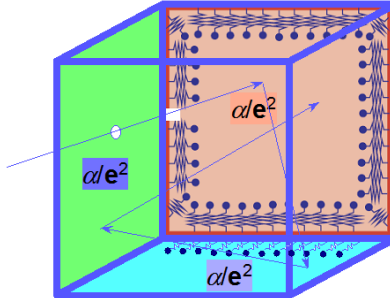


Figure 2: Blackbody radiation is related to α and e^2 in a 3D cavity.

This 3D box (or sphere) model does not necessarily have solid walls, the plasma gas layer of a star can have the same effect. This links the quantum theory to the classical theory of blackbody radiation with or without using the Planck constant

$$\begin{aligned} \alpha_R &= \left(\frac{\alpha_R}{e^2}\right)^3 k_B^4 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{\alpha}{e^2}\right)^3 k_B^4 = \frac{\zeta(4)}{\zeta(2)} \left(\frac{2\pi}{hc}\right)^3 k_B^4 \\ &= \left(\frac{\alpha}{e^2}\right)^3 k_B^4 \cdot \frac{4}{5} \cdot \frac{9}{10} \cdot \frac{25}{26} \cdot \frac{49}{50} \cdot \frac{121}{122} \cdot \frac{169}{170} \cdot \frac{289}{290} \cdot \frac{361}{362} \dots \end{aligned} \quad (15)$$

The Planck constant h with the revolutionary concept of energy quanta is a bridge between classical physics and quantum physics. Einstein's proposal of the light quanta $h\nu$ in 1905 was based on the Planck constant, however, Planck always had reservations due to the continuous spectra of blackbody radiation and the wave-particle duality. In 1951, Einstein said that, "All these fifty years of conscious brooding have brought me no nearer to the answer to the question, 'What are light quanta?'" [7] In QED, the photon is treated as a gauge boson, and the perturbation theory involves the finite power series in α . The discrete-continuous spectra is bridged by the Bose-Einstein distribution, and the *prime* sequences link the fine structure constant α to the blackbody radiation constant α_R . The blackbody radiation constant is a new method to measure the fine structure constant. It also links the fine structure constant to the Boltzmann constant.

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