

# Realization of quasi-quanta via the forced contraction of loops

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## Abstract

The contraction of a loop on a string in the orthogonal time direction is contemplated. Its relationship to a certain mathematical concept, forcing notions, is examined. In addition, we evaluate local systems on the worldline of a particle traveling in the positive timelike direction.

## 1 Contracting loops

**Definition 1** *A slice,  $s$ , of a manifold  $M$ , is a stationary frame  $\mathfrak{F}$  whose subobjects are potentials of a gauge field.*

Let  $\mathfrak{S}$  be a collection of slices of a manifold  $\mathbb{R}^n$ ,  $n$  even, and let there be a path  $\rho : \inf(\mathfrak{S}) \rightarrow \sup(\mathfrak{S})$ . The product,  $\rho \cdot \mathfrak{S} = \{p, g\}$ , defines an equivalence class of diffeotopic smooth manifolds spanning a lightcone  $\mathbb{L}$  of events.

For every segment of the vector  $\vec{\rho}$ , define a *portable loop*,  $p_\ell$ .

**Definition 2** *A portable loop is a neighborhood about a fixed point  $p_i$  along the path spanned by  $\vec{\rho}$ , such that  $p_\ell^{-1} \circ p_\ell$  is the identity on  $p_i$ .*

Suppose the region enclosed by some  $p_\ell$  contains non-trivial topological (i.e., physical) data, such as a particle. Then, we say that the loop is not contractible to the point  $\hat{q}$ . We impose the following:

**Axiom 1** *For all non-contractible portable loops, there exists either (or both) of the following:*

- *A forcing notion,  $\Vdash$ , such that  $\hat{q} \Vdash p_\ell^{contr}$*
- *A frame  $\mathfrak{F}^+$  such that  $p_\ell^{contr} \in \mathfrak{F}^+$*

These essentially (although not equivalently) amount to designating a map  $t^+ : p_\ell \rightarrow \{*\}$ ; in other words, the successor function is applied to the dimension of time (in the second case), which *forces* the suspension of a quasi-quantum  $\hat{q}$ , thus transforming the non-trivial space it (virtually) occupies into a generic free variable.

Let us imagine that  $p_\ell$  has the following properties, where  $\mathcal{U}$  stands for the neighborhood it encloses:

1. Pressure increases over time as  $p_\ell \longrightarrow \{*\}$
2. Temperature increases within  $\mathcal{U}$  as it shrinks
3. The average number of molecules in  $\mathcal{U}$  decreases as  $p_\ell$  contracts

Then,

**Theorem 1** *Boltzmann's constant,  $k_B$ , remains constant as  $p_\ell \longrightarrow \{*\}$ .*

**Proof** We have

$$k_B = \frac{PV}{TN}$$

Assuming the decrease in the number of molecules in  $\mathcal{U}$  is proportional to the increase in temperature, the differences cancel out; assuming that pressure increases as the volume of  $\mathcal{U}$  decreases cancels out the terms in the numerator. Thus,  $k_B$  is constant under the map  $p_\ell \longrightarrow \{*\}$ .

## 1.1 Realization

In this paper, we envision that there is a certain *semitotic* propensity of quasi-quanta (virtual particles) to become actualized (topologically realized) as a result of the satisfaction of Axiom 1. Thus, the promotion of a wavefunction to a particle can be interpreted either as a class-theoretic (mathematical) operad, or as physical kinematics occurring across time in  $\mathbb{L}$ . The equation:

$$\mathbb{I} \times_h \hat{q}_{pot} \longrightarrow q = \mathfrak{F}^+ = \hat{q} \Vdash p_i^{contr}$$

relates the two paradigms to the production of a particle from an operator  $h$  mediating between an interval  $\mathbb{I}$  and the potential energy of a quasi-quanta. We can think of this as a sort of *crossed module*,  $\mathfrak{m}$ , acting on the group  $\mathfrak{g}_p$  of generators for the Lie group of the particle's neighborhood along a worldline. We remark here that the world-line,  $\mathfrak{W}$ , is a special case of the path  $\rho$  defined above, which has been restricted to timelike distinct sections of  $\mathbb{L}$ .

Motion along  $\mathfrak{W}$  obeys the following Leibniz rule:

$$(\mathfrak{m}i + \mathfrak{m}j)k = \mathfrak{m}ik + \mathfrak{m}jk = \partial\omega^{-1}\mathfrak{m}ijk$$

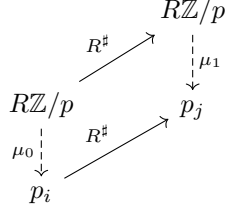
where  $\omega$  is a differential form of dimension equal to a particular Lagrangian submanifold of  $M$ . Thus, the function

$$f(\mathfrak{m}) = \mathfrak{u} \xrightarrow{\theta} \mathfrak{u}' \subset_{\mathbb{N}} \mathfrak{g}_p$$

is transitive for all smooth paths which non-trivially intersect covers of neighborhoods over  $p_i \in M$ . Classically,  $f(\mathfrak{m})$  defines a formally flat function

acting on trivial transport fibers. In our case, each segment  $T_x(p_i)$  tangent to the moment of a particle yields a foliation along a boundary  $\partial\mathcal{U}$  which projects to a singularity  $\mathfrak{b} \in \mathbb{L}$ , where a measurement either does or does not occur.

We may choose to enrich each copy of  $\partial\mathcal{U}$  with a connected ring of polynomials modulo a certain prime,  $p$ , giving us  $R\mathbb{Z}/p$ . Correspondingly, transport of the particle  $p_i$  is described as a *tilting*,  $R^\sharp$  which forms an  $\aleph$ -cell about a wave packet.



In the above parallelogram,  $p_i, p_j$  are distinct particles which share the same wavefunction  $\Psi(p, t)$ , and  $\mu_0, \mu_1$  are measurements, which are, respectively, projection onto the first and final coordinates of a local system. Directedness of the arrow  $p_i \rightarrow p_j$  denotes the irreversibility of time due to the second law of thermodynamics.

**Definition 3** A local system,  $LocSys(h)$ , is a closed, portable monoid equipped with a counting operad on lines.

Local systems naturally come with a bundle,  $Bun_G$ , which induces a simplicial stratification over a Hausdorff convex neighborhood of a manifold  $M$ . A local-system is  $G$ -equivariant with respect to reordering (shuffling) of place values, and is uniquely determined (up to isomorphism) by a collection of paths  $\vec{P}G$  out of any given point  $p$ . Thus, the identity of a local system is given by:

$$LocSys_{Id} = \int_0^{2\pi} \frac{\partial p_i}{di} \Omega G'$$

where  $\Omega G'$  is space of loops of any other Lie Group. This is essentially the Yoneda lemma for Markov blankets.

Let  $\pi_\eta$  be a map of fibers over  $LocSys(M)$ . We denote by  $Spec_\eta$  the spectral sequence:

$$\Pi : U(1) \rightarrow \eta_{ij} \rightarrow \eta_{jk} \rightarrow \eta_{ki} \rightarrow U(1)$$

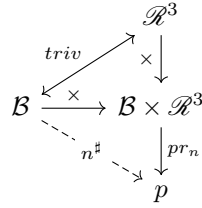
which is smooth. Denote the composition  $\Pi \circ^n \Pi$  by  $Nec_n(\Pi)$ . One has that the canonical 2-morphisms,  $\overset{b}{a} : (a, b) \rightrightarrows (c, d)$  are stable under the stack  $\mathcal{X}_{Top}$ , and the isofibrations  $[\overset{b}{a}]$  are arbitrarily productive. This means that we can take the quotient  $Nec_n(\Pi)/q$  and obtain a Hermitian Koszul complex,  $\mathcal{K}_{osz}$ , which preserves holonomy. Write

$$\mathcal{K}_{osz} = (LocSys(i) \times LocSys(j)) \xrightarrow{can} \Pi_{\bar{\omega}}$$

**Definition 4** A Koszul complex is a global system whose interior consists of the disjoint union of the symmetric product of  $n$  local systems.

All neighborhoods  $\mathcal{U}$ , and smooth covers  $\{\mathcal{U}_i\}_{i \in I}$  essentially arise as rank two restrictions of Koszul complexes. That is to say, that for each stalk  $f$  of  $\mathcal{K}_{osz}$ , there exists an infinitesimal thickening on the points of  $f$  (call them  $\tilde{f}_i$ ), which are thin homotopies of rank two of one another, such that  $\pi_2^2(\tilde{f} \in f)$  yields a conformal pullback to a site  $\theta$  at which the functions  $f(f)$  converge asymptotically.

**Remark 1** For a flat bundle,  $\mathcal{B} \in \mathcal{R}^3$ , the collection of tangent spaces over each point  $x \in \mathcal{B}$  contains a space whose projection onto the  $n$ th coordinate,  $x \xrightarrow{n} \mathcal{R}^3$ , converges to a point  $p \in \mathcal{B} \times \mathcal{R}^3$ .



The necklace,  $Nec_p(\mathcal{B})$  gives the set of all maximal chains

$$\Delta \times p^{-1} \longrightarrow max|\mathcal{B}|$$

## 2 Creation and Transformation

Prior to the assignment of meaning to a symbol (or better yet, a symbol to a meaning), the “meaning” to be specified remains in a superposition of possible states, which we denote by  $\heartsuit$ . The so-called “creation map” defined below, is an exit path

$$\mathcal{EP}_{\heartsuit} : (x = \{\}) \longrightarrow x$$

out of the empty set, into a proposition  $x$ .

**Definition 5** The creation map,  $Cr$  shall be written

$$Cr : \heartsuit \longrightarrow \{p\}$$

where  $p \sim \{*\}$  for some zero-dimensional manifold.

**Definition 6** The transform map,  $T$ , is given by

$$(x \rightarrow y) \leftrightarrow \exists x \vee y \rightarrow \exists y$$

where  $y$  is the top of a frame.

**Remark 2** The existential quantifier used here is classical (strong).  $T$  thus represents a map  $\exists^\bullet \rightarrow \exists(\sim)$ . The sequence

$$\heartsuit \xrightarrow{Cr} x \xrightarrow{T} y \xrightarrow{(Cr \circ T)^{-1}} \tilde{\heartsuit}$$

is equivalent to  $f^!(x, y) \circ f_!(x)$ . Further, the identity on a fixed object,  $x$ , is  $Cr = T^{-1}$ .

Let  $X$  and  $Y$  be subsets of  $Z$ . Then, let  $x \in X$  and  $y \in Y$ . We have  $x \in^X Z$  and  $y \in^Y Z$  defining a filtered inclusion relationship, such that the superscript  $\in^\bullet$  entails  $\exists \bullet \in \sim$ , where  $\sim$  is the least upper bound on all  $\bullet, \bullet'$ .

Each creation map corresponds to an actual measurement,  $\mu_x$  over the object  $x$ , and a transformation map represents a first order differential taken over  $\mu_x$ .

$$T = \mu_y - \mu_x = d\mu_x$$

## 2.1 Sampling Populations of Measures

Let there be a large number of measurements taken across a sample  $\mathfrak{s}$  of transform maps. We shall write

$$\frac{\Sigma(d\mu_x)}{\text{card}(\mu)} = \varphi(\mu)$$

to mean the average of the differences between each set of correlated pairs of points,  $x$  and  $y$ .

**Proposition 1** Let  $t^+ : [0, 1] \rightarrow [0, 1]$  be the time-step functor. We have  $\lim_{t^+} T = \varphi(\mu)$ .

**Proof** Our argument is proved by writing:

$$\varphi(\mu) = \lim_{\mu \rightarrow \infty} \frac{\Sigma(d\mu_x)}{\text{card}(\mu)} = T_\infty$$

due to the fact that  $t^+$  is essentially the functor  $\mu \rightarrow \mu + 1$ . As a result, the average transformation asymptotically approximates the universal average taken over an infinite population. *Q.E.D.*

## 2.2 Generalized transforms and their actualization

We were motivated to form a *process-based* definition of an object. In this pursuit, we have established the following:

**Definition 7** A transformation,  $T(x)$ , is a map  $x \rightarrow ?$  such that  $T^{-1}(x)$  is the identity on  $x$ .

One may be dismayed that this definition lacks a clear-cut physical interpretation. So, let us renew this definition, this time taking into account the wavefunction on a particle  $x$ :

$$T(x) = (\Psi(x) \longrightarrow ||x||) \vee (||x|| \longrightarrow ||y||)$$

Call the left-hand side of the disjunction the *actualization*. This is *one type of* creation map, but it is also implicated in the transformation process. The right hand side represents the ordinary transformation of observable eigenstates. Over a small period of time-evolution, quasi-quanta may enter or exit a given eigenstate, as parameterized by a given truth value  $\tau$ . The evolution of time,  $\tau \rightarrow \tau + 1$ , is a form of monodromy in the F-theory description, where  $\tau$  is the modulus of elliptic fibers of some locus  $Y$ . This can (and has) been used to model dynamics on intersecting seven-branes. We refer the reader to [1] for more information.

### 3 References

- [1] C. Cordova, *Decoupling gravity in F-theory*, (2011)