

# Hamiltonian Instability and the Classical-to-Quantum Transition

Ervin Goldfain

Ronin Institute, Montclair, New Jersey 07043, USA

E-mail [ervin.goldfain@ronininstitute.org](mailto:ervin.goldfain@ronininstitute.org)

## Abstract

The mechanism of *Arnold diffusion* (AD) describes the dynamic instability of nearly-integrable Hamiltonian systems. Here we argue that AD leads to action quantization for classical systems having an infinite number of degrees of freedom. Planck's constant emerges as long-time value of the action differential applied to large ensembles of oscillators in near equilibrium conditions.

**Key words:** Arnold diffusion, Hamiltonian chaos, Planck constant, action quantization, classical to quantum transition.

The formalism of *action-angle variables* applies to generic Hamiltonian systems undergoing periodic motion, such as harmonic oscillation or Kepler

rotation. It consists of replacing the generalized coordinates and momenta using the transformation [1-3]

$$(q, p) \rightarrow (\theta, I) \quad (1)$$

Action-angle variables are canonically conjugate and introduced through the generating function [3]

$$S(q, I) = \int_q p(q, H) dq \quad (2)$$

such that

$$I = \frac{1}{2\pi} \int_c p(q, H) dq = I(H) \quad (3a)$$

$$\theta = \frac{\partial S(q, I)}{\partial I} \quad (3b)$$

Since  $I$  is a cyclic variable, the corresponding drift of (2) per each period of  $I$  amounts to [1]

$$\Delta S = 2\pi I \quad (4)$$

Consider a nearly-integrable periodic system with  $n$  degrees of freedom defined by the Hamiltonian [4-6]

$$H = H_0(I) + \varepsilon H_1(I, \theta, \varepsilon) \quad (5)$$

where  $0 \leq \varepsilon \leq \varepsilon_0$  is a small perturbation parameter and  $H_0(I)$  is the unperturbed Hamiltonian, taken to be fully integrable in the limit  $\varepsilon = 0$ . The frequency of the unperturbed motion is determined by

$$\omega(I) = \frac{\partial H_0}{\partial I} \quad (6)$$

For  $\varepsilon \leq \varepsilon_0 \ll 1$ , the equations of motion read

$$\dot{I} = -\frac{\partial H(\theta, I)}{\partial \theta} = 0 \quad (7)$$

$$\dot{\theta} = \frac{\partial H(\theta, I)}{\partial I} = \omega(I) \quad (8)$$

in which  $I, \theta \in \mathbf{R}^n$ . The solution of (7)-(8) lies on invariant  $n$ -tori residing in the phase-space of dimension  $\mathbf{R}^{2n}$ . For  $n \leq 2$ , all solutions are stable since 2-tori confine trajectories on a 3-dimensional energy surface. This is no longer the case for  $n \geq 3$  where, according to the *Arnold diffusion conjecture*, the

action of nearly-integrable systems changes by  $O(1)$  over a sufficiently long time. In particular, assuming that

$$|H_0| < c_1, |H_1| < c_1 \quad (9)$$

where  $c_1$  is a positive constant, and taking the unperturbed Hamiltonian to represent a quasi-convex function of the action variable, the following condition holds [6, 8]

$$\delta I = |I(t) - I(0)| < C_1 \varepsilon^{1/2n} \quad (10)$$

over sufficiently long-times satisfying

$$0 \leq t \leq \exp\left(\frac{C_2^{-1}}{\varepsilon^{1/2n}}\right) \quad (11)$$

In (10) and (11),  $C_1, C_2$  are also positive constants. By (4), the corresponding drift in action is given by

$$\delta(\Delta S) = 2\pi \delta I < O(C_1 \varepsilon^{1/2n}) \quad (12)$$

Normalizing (12) to  $C_1$  confirms that the drift in action is of  $O(1)$ , which naturally replicates the process of *action quantization* for  $n \gg 1$ . It follows

that Planck's constant may be mapped to the long-time value of (12) applied to large ensembles of oscillators in near equilibrium conditions. Stated differently, the transition from classical to quantum behavior is expected to occur when

$$0 \leq t \leq \exp\left(\frac{C^{-1}}{\varepsilon^{1/2n}}\right); n \gg 1 \quad (13)$$

leading to

$$\boxed{\delta(\Delta S) = 2\pi\delta I = O(1)} \quad (14)$$

These findings fall in line with the approach taken in [7].

## **References**

1. Landau L. D., Lifschitz E. M., Mechanics, Butterworth-Heinenann, 3<sup>rd</sup> Edition, 1976.
2. Corben H. C., Stehle P., Classical Mechanics, Dover Publications, 1966.
3. Zaslavsky, G. M., Hamiltonian Chaos and Fractional Dynamics, Oxford, 2006.

4. <http://authors.library.caltech.edu/13561/1/KALbams08.pdf>
5. <https://webusers.imj-prg.fr/~pierre.lochak/textes/compendium.pdf>
6. <https://www.researchgate.net/publication/10731460>
7. <https://www.researchgate.net/publication/343403813>
8. [https://en.wikipedia.org/wiki/Nekhoroshev\\_estimates](https://en.wikipedia.org/wiki/Nekhoroshev_estimates)