

Cable Capacitance Attack against the KLJN Secure Key Exchange

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Abstract: The Kirchhoff-law-Johnson-(like)-noise (KLJN) is an unconditionally secure key exchange system based on the laws of classical statistical physics. Similarly to quantum key distribution, in practical situations, due to the non-idealities of the building elements, there is a small information leak, which can be mitigated by privacy amplification or other techniques. In this paper, a professional cable and circuit simulator LTSPICE is used to validate the information leak due to one of the non-idealities in KLJN, the parasitic capacitance. Simulation results show that privacy amplification can efficiently mitigate the leak while a capacitor killer can fully eliminate it.

Keywords: KLJN; cable capacitance attack; capacitor killer; secure key exchange; unconditional security; privacy amplification.

1. Introduction

The Kirchhoff-law-Johnson-(like)-noise (KLJN) based unconditional secure key exchange system [1-4] was first introduced in 2005. Before KLJN, it was assumed that only Quantum Key Distribution (QKD) [5] could offer this level of security. However, QKD's fundamental security has been debated by experts in the field [6-12]. Furthermore its practical realizations, including all commercial quantum communicators, have been fully cracked by hacking, that is utilizing non-ideal features of the hardware building elements [13-26]. While countermeasures were later proposed to overcome these attacks, when the idea of a new attack is unknown by the communicating parties, the eavesdropper can fully utilize such an attack because no counter-measures have been developed and implemented yet. [27-30].

Naturally, there have been efforts to challenge KLJN's security [31-42]. Studies have consistently shown that both the ideal and the practical KLJN versions remain unconditionally secure [4,34-42] despite facing various attacks and related information leaks associated with the non-idealities of components in the system. The impacts of the attacks on the practical KLJN system have been weak with the Eavesdropper's (Eve) probability of successful guessing of bits approaching zero [3,34-38].

We will show that one of the most effective attacks against the practical KLJN system is called a cable capacitance attack. It was first mentioned in 2006 [36], but it has never been realized. Subsequently, in 2008, a solution was suggested to eliminate the attack by adding a capacitor killer arrangement [39]. In this paper, we use a professional cable and circuit simulator LTSPICE by Linear Technology to simulate a practical realization of the KLJN system and to evaluate the cable capacitance attack. Solutions to mitigate this attack, such as the capacitor killer arrangement [39], and privacy amplification [43] are also evaluated.

2. The KLJN secure key exchange system

2.1 The KLJN protocol

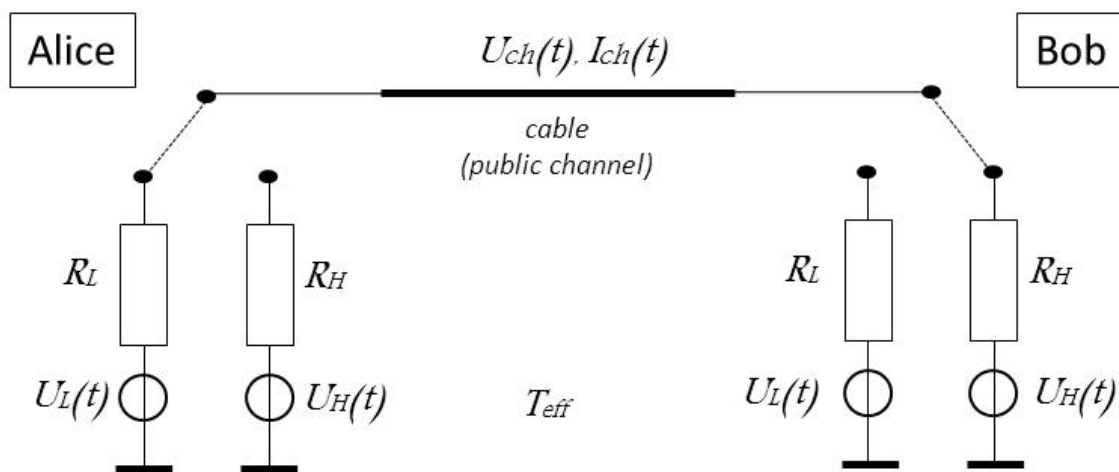


Figure 1. The core of the KLJN secure key exchange system [2]. R_L and R_H are the resistor values. $U_L(t)$ and $U_H(t)$ are the thermal noise voltages of R_L and R_H , respectively. $U_{ch}(t)$ and $I_{ch}(t)$ are the channel noise voltage and current, respectively.

The KLJN secure key exchange system [1-4,38-54] utilizes Kirchoff's Loop Law and the Fluctuation-Dissipation Theorem. The core KLJN system is illustrated in Fig. 1 [2]. It is composed of a cable as an information channel, and two identical pairs of resistors (R_L and R_H) possessed by a sender (Alice) and a receiver (Bob) respectively, where R_L represents the Low key bit (0) and R_H represents High key bit (1) at Alice's side (at Bob's side the opposite), with $R_L \neq R_H$ and $R_L < R_H$, and two switches at each end connecting the selected resistors to the cable.

At the beginning of each bit exchange period (BEP), Alice and Bob randomly select R_L or R_H at their ends and connect the corresponding resistor to the cable. The Gaussian voltage noise generators represent the Johnson noises of the resistors (denoted as $U_L(t)$ and $U_H(t)$ for R_L and R_H respectively), which deliver band-limited white noise with publicly agreed bandwidth B_{noise} and a publicly agreed effective temperature T_{eff} [40]. The noises are statistically independent from each other and from the noises in the previous BEP [4].

Within each of the BEP, Alice and Bob measure the mean-square channel noise voltages $\langle U_{ch}^2(t) \rangle$ and/or the channel noise currents $\langle I_{ch}^2(t) \rangle$ in the cable. The BEP has to be properly chosen to provide sufficient time for getting a good statistics of the mean-square noise voltages and currents but not enough time for Eve to utilize possible information leaks due to hardware non-idealities. According to Johnson's noise formula and Kirchhoff's Loop Law, it is given that:

$$\langle U_{ch}^2(t) \rangle = 4kT_{eff} \frac{R_A R_B}{R_A + R_B} B_{noise} \quad (1)$$

$$\langle I_{ch}^2(t) \rangle = 4kT_{eff} \frac{1}{R_A + R_B} B_{noise} \quad (2)$$

where k represents the Boltzmann's constant ($1.38 \times 10^{-23} \text{ J/K}$), T_{eff} is the effective temperature (typically $T_{eff} \geq 10^9 \text{ K}$), R_A and R_B are the resistors selected by Alice and Bob respectively in Ohm, and B_{noise} is the noise bandwidth.

Based on equation 1 or 2, by measuring $\langle U_{ch}^2(t) \rangle$ and/or $\langle I_{ch}^2(t) \rangle$, and by having their own resistors' value, Alice and Bob can figure out which of the resistors is used by the other party and hence they can identify the bit (0 or 1) sent from each other.

With the cable being public, an eavesdropper (Eve) can also measure the channel noise voltages and currents to obtain the total loop resistance in the cable. If Alice and Bob use the same resistance values, $R_L R_L$ or $R_H R_H$, the exchanged bit is non-secure and is discarded [2]. Conversely the combinations $R_L R_H$ and $R_H R_L$ are secure, and Eve cannot differentiate between the two alternatives. This is because the mean-square resultant noises are identical for $R_L R_H$ and $R_H R_L$ provided that the cable is ideal and short. Eve knows that Alice and Bob has exchanged a secure bit, but she does not know who is using R_L and who is using R_H .

In reality, the cable is non-ideal. Thus Eve can exploit the non-idealities of the cable, such as parasitic resistance, parasitic inductance and parasitic capacitance to attack the KLJN system. We will discuss parasitic capacitance in next section, in detail.

2.2 Cable capacitance attack [36]

In this paper, we discuss coaxial cables because, in this case, the cable capacitance attack can fully be eliminated. However, the attack works with any cable. Coaxial cables include two conductors: the inner wire, which is used as the KLJN channel, and the outer shield which is grounded (for the ground, see Fig. 1. There is a non-zero capacitance between these two conductors that leads to the capacitive currents. This capacitive current is where the information leak comes from. Part of the channel noise current is diverted by the parasitic capacitance, which causes a greater current at the side of the lower resistance. This gives Eve a chance to guess the value of the resistor at each side with probability of success greater than 0.5.

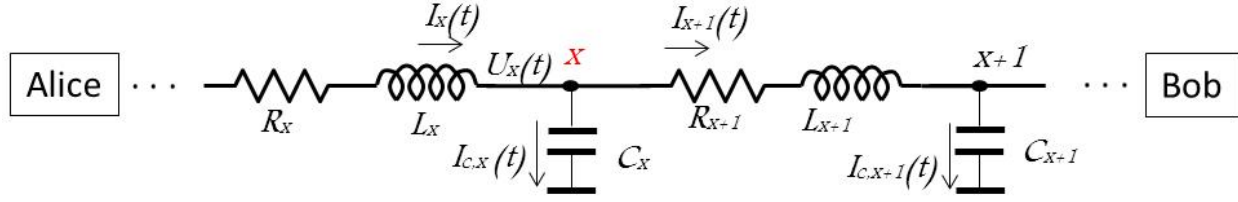


Figure 2. Cable model and cable capacitance attack

Fig. 2, shows the distributed elements model of coaxial cable. According to the Kirchoff's current law, at position x , the channel noise current $I_x(t)$ is the sum of the capacitive current $I_{c,x}(t)$ through the parasitic capacitor element C_x , and the Ohmic channel noise current $I_{x+I}(t)$. This is written as

$$I_x(t) = I_{c,x}(t) + I_{x+I}(t). \quad (3)$$

The capacitive current $I_{c,x}(t)$ is proportional to the time derivative of the channel noise voltage $U_x(t)$ and it is given by

$$I_{c,x}(t) = C_x \cdot \frac{dU_x(t)}{dt}. \quad (4)$$

We define a cross-correlation $\rho(x)$ at position x as the product of the channel noise current and the time derivative of the channel noise voltage:

$$\rho(x) = \left\langle I_x(t) \cdot \frac{dU_x(t)}{dt} \right\rangle_{\tau}, \quad (5)$$

where $\langle \rangle_{\tau}$ means finite time (τ) average. The cross-correlation ρ represents information leak [34].

3. Realization of the attack

We utilized the cable and a circuit simulator LTSPICE by Linear Technology to emulate the practical KLJN system. The RG58 coaxial cable from the library was used. Throughout the simulations, we assumed that Alice randomly selected R_L of 1 kOhm and Bob randomly selected R_H of 9 kOhm, which connected to the cable as illustrated in Fig. 3.

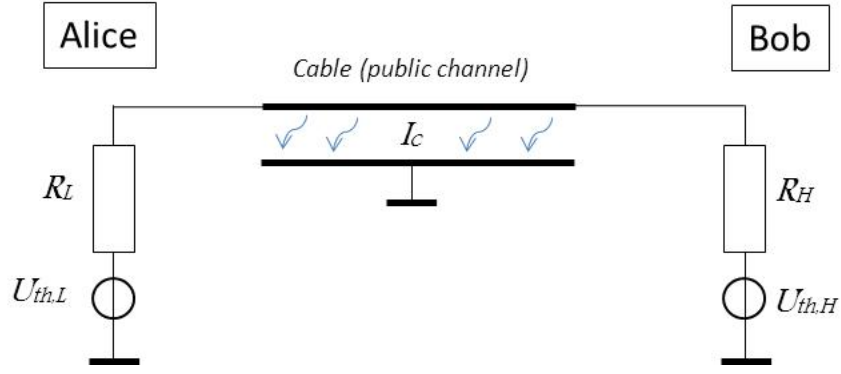


Figure 3. The simulated KLJN secure key exchange system with the capacitive current I_c . The generator voltages $U_{th,L}$ and $U_{th,H}$ are the root-mean-square (rms) Johnson noise amplitudes of R_L and R_H , respectively.

3.1 Generating the noise

For the simulations, we generated the Gaussian band-limited white noises. According to Johnson's noise formula, the rms thermal noise voltage U_{th} on an open-circuit resistance R in thermal equilibrium at temperature T_{eff} is

$$U_{th} = \sqrt{4kT_{eff}RB_{noise}}. \quad (6)$$

As the mean value is zero, the rms noise voltages are the same as their standard deviations (denoted as σ_L and σ_H for $U_{th,L}$ and $U_{th,H}$ respectively). Thus

$$U_{th,L}/U_{th,H} = \sigma_L/\sigma_H = \sqrt{R_L/R_H}, \quad (7)$$

where $\sqrt{R_L/R_H} = 1/9$, thus $\sigma_L/\sigma_H = 1/3$. For the simulations, the rms thermal noise voltages of R_L and R_H were chosen 1V and 3V, respectively.

The amplitude density function and the cumulative distribution of the noise voltage of R_L is shown in Fig. 4. The graphs' shape indicated that the thermal noise generated was in good quality, especially as shown by the straight line in the normal distribution plot in Fig. 4(b). For the sake of simplicity, the plots of R_H are not shown as they looked the same as that of R_L but with standard deviation = 3V.

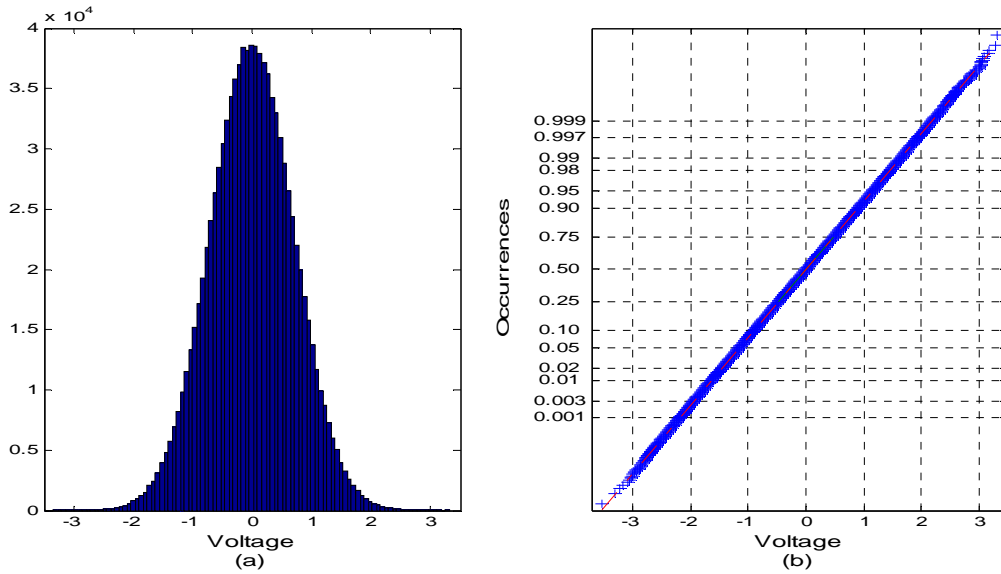


Figure 4. Thermal noise voltage of R_L with 10^6 samples. (a) amplitude density (histogram); (b) cumulative distribution as normal probability plot where straight line indicates exactly Gaussian distribution.

3.2 Comparing a lumped and the distributed element models at different wavelengths

The quasi-static condition is required for the security of the KLJN system [2, 34]. That means

$$L_{ch} \ll \lambda = c/B_{noise} \quad \text{or} \quad \gamma = \lambda/L_{ch} \gg 1, \quad (8)$$

where L_{ch} is the cable length, λ is the shortest wavelength at the highest frequency component in the noise bandwidth B_{noise} , c is the propagation velocity in the cable, and γ is the ratio of the wavelength to the cable length. The γ should be much larger than 1 in order to fulfill the “no-wave” condition (i.e., quasi-static electrodynamics) [34]. But, how large should γ be for the KLJN system?

Fig. 5(a) and 5(b) shows a simple lumped element model and the distributed model of the RG58 coaxial cable respectively. Three simulations were run to compare the resultant voltage waveforms at Alice’s side, at three different noise bandwidths B_{noise} (250 kHz, 25 kHz, 0.25 kHz) on these 2 models. The cable length was set at 1000 meters, based on equation 8, the three corresponding wavelengths (λ) were 800 m, 8 km, and 800 km, while the corresponding γ ratios were 0.8, 8 and 800. Other parameters such as the component values of the models used in the simulations are also shown at Fig. 5.

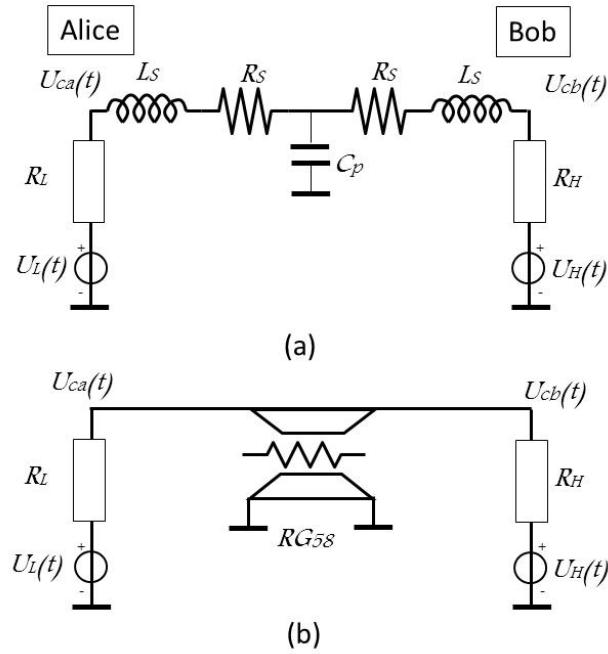


Figure 5. The RG58 coaxial cable models (1000 m length) with R_L (1 kOhm) and R_H (9 kOhm).

(a) The lumped element model: the component values: $R_S = 10.5 \text{ Ohm}$, $L_S = 125 \mu\text{H}$, $C_P = 100 \text{ nF}$.

(b) The distributed model had the following parameters: $R=0.021 \text{ Ohm/meter}$, $L=250 \text{ nH/meter}$, $C=100 \text{ pF/meter}$. The characteristic impedance of the cable is 50 Ohm. The propagation velocity c in the cable is $2 \times 10^8 \text{ meter/sec}$.

Fig. 6 shows the simulation results in which $U_{ca,lump}$ and $U_{ca,cont}$ denoted the voltage timefunctions of the lumped and distributed element models, respectively. In Fig. 6(a), the two waveforms were significantly different for the shortest wavelength with $\gamma = 0.8$. As the waves can only be simulated in the distributed model but not in the lumped element model, the two timefunctions look substantially different.

In Fig. 6(b), with $\gamma=8$, the two waveforms are very similar whereas in Fig 6(c), at $\gamma=800$, the two waveforms are indistinguishable. Both cases are fine for the KLJN operation thus we can conclude that for $\gamma \geq 8$, the lumped element simulations are satisfactory.

For our resistor values $R_L = 1 \text{ kOhm}$ and $R_H = 9 \text{ kOhm}$, the channel bandwidth of the practical KLJN system is 1.76 kHz and 17.6 kHz for a 1000 and a 100 meters cable, respectively. In order to avoid the noise bandwidth getting truncated by the system, we utilized noise bandwidth $B_{noise} = 0.25 \text{ kHz}$ ($\gamma = 800$) for the following simulations.

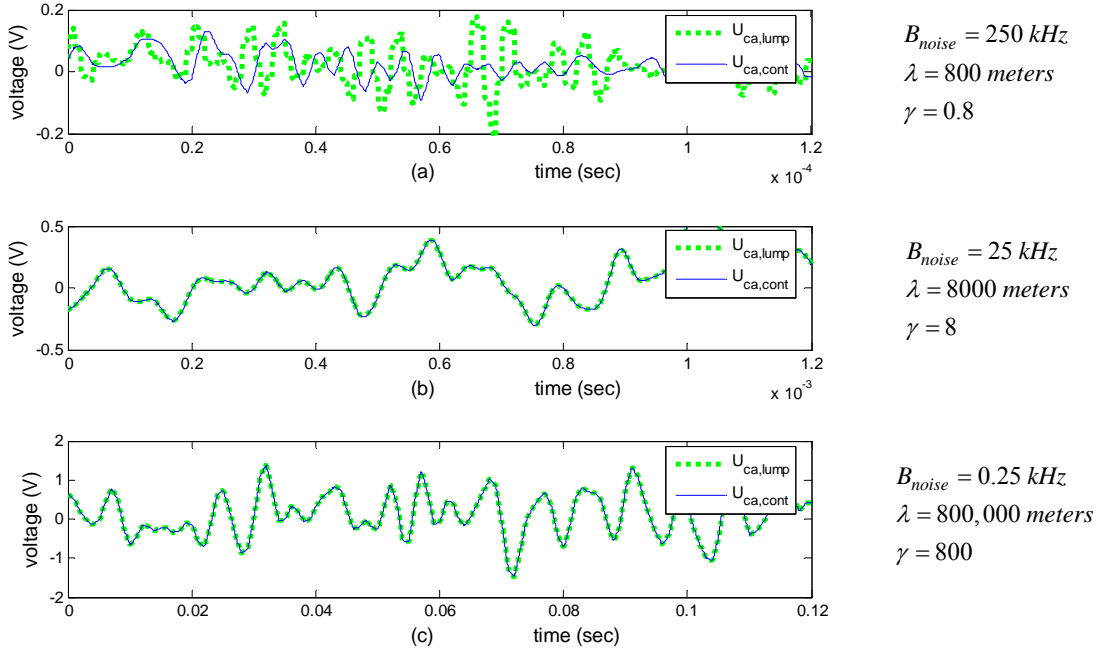


Figure 6. The voltage waveforms, $U_{ca,lump}$ and $U_{ca,cont}$, at Alice's side for the lumped and distributed element models, respectively, for a 1000 meters cable, at 3 different ratios: (a) $\gamma=0.8$; (b) $\gamma=8$; (c) $\gamma=800$.

3.3 The attack protocol

In this section, we will discuss how the information leak caused by the parasitic capacitance, that is, Eve's success probability of guessing the secure key bit.

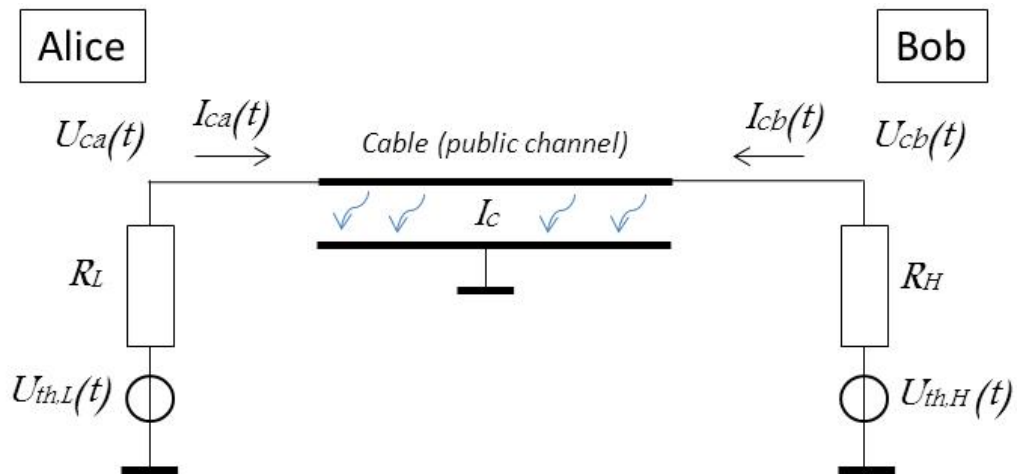


Figure 7. The simulated model with $R_L = 1 \text{ k}\Omega$ and $R_H = 9 \text{ k}\Omega$. U_{ca} , I_{ca} , U_{cb} and I_{cb} are the voltages and currents at Alice's and Bob's ends, respectively.

The cross-correlations ρ_{ia} and ρ_{ib} are:

$$\rho_{ia} = \left\langle I_{ca}(t) \cdot \frac{dU_{ca}(t)}{dt} \right\rangle_{\tau}, \quad (9)$$

$$\rho_{ib} = \left\langle I_{cb}(t) \cdot \frac{dU_{cb}(t)}{dt} \right\rangle_{\tau}, \quad (10)$$

where $U_{ca}(t)$, $I_{ca}(t)$, $U_{cb}(t)$ and $I_{cb}(t)$ are the voltages and currents at Alice's and Bob's ends, respectively, see Fig. 7. the time average $\langle \rangle_{\tau}$ is taken over the bit exchange period τ . We take the difference between the 2 cross-correlations, $\rho_i = \rho_{ia} - \rho_{ib}$ ($i = 1, \dots, N$) and to evaluate Eve's statistics, we decide as:

$$\begin{aligned} \text{If } \rho_i > 0 & \text{ then } q_i = 1 \quad (\text{Eve guessed the bit correctly}) \\ \text{If } \rho_i < 0 & \text{ then } q_i = 0 \quad (\text{Eve guessed the bit wrongly}) \end{aligned}, \quad (11)$$

When N approaches infinity, then the probability of Eve's successful guessing of the bits is equal to the expected value of q and

$$\langle q_i \rangle_N = p_E = 0.5 + \varepsilon, \text{ where } 0 \leq \varepsilon < 0.5. \quad (12)$$

where the non-zero ε represents an information leak. When $\varepsilon = 0$ the KLJN key exchange system is perfectly secure. We found that the higher the difference between the resistances, the bandwidth, or the parasitic capacitance (longer the cable), the higher the leak.

3.4 Simulation results of the cable capacitance attack

We simulated 6 different attack scenarios with these parameters: $R_L = 1$ kOhm, $R_H = 9$ kOhm, noise bandwidth $B_{noise} = 0.25$ KHz, sampling period $t_s = 1$ msec, for 3 different sample numbers (correlation times) per bit (20, 50, 100), at 2 different cable lengths (100 and 1000 meters). At each scenario, the key was 1000 bits long.

The simulation results are shown in Table 1. At 20 samples per bit (50 bits per second), with a 100 meters cable, Eve's success rate was 50.9%. However, when the cable length was increased to 1000 meters with the other parameters unchanged, Eve's success rate became to 62.2%.

Table 1. Attack simulation results - Eve's success rate p_E (%) with 1000 bits key length

Samples per bit	Bits per second	100 meters cable	1000 meters cable
20	50	50.9%	62.2%
50	20	52.1%	69.7%
100	10	52.6%	76.9%

If the samples per bit were increased to 50 and 100, Eve's success rate increased accordingly as shown in Table 1. In the most effective attack case, with 100 samples per bit and 1000 m cable, Eve success rate was 76.9%.

4. Defense against the attack

4.1 Capacitor killer

The parasitic capacitance of the RG58 coaxial cable can be eliminated by providing the same voltage on the outer shield of the cable as on the inner wire [39]. This can be done by a unity-gain voltage buffer, which is called capacitor killer. A capacitor killer is connected between the inner wire and the outer shield to make the electric potentials between them equal, as shown in Fig. 8. There is no capacitive current from the inner wire to the outer shield thus, the attack is nullified.

We simulated the capacitor killer on the most effective attack scenario (i.e., when Eve success rate was 76.9%). The simulation results showed that after capacitor killer was added, Eve success rate was reduced from 76.9% to 50.1%. This indicated that the capacitor killer is very effective in eliminating the leak due to the parasitic capacitance.

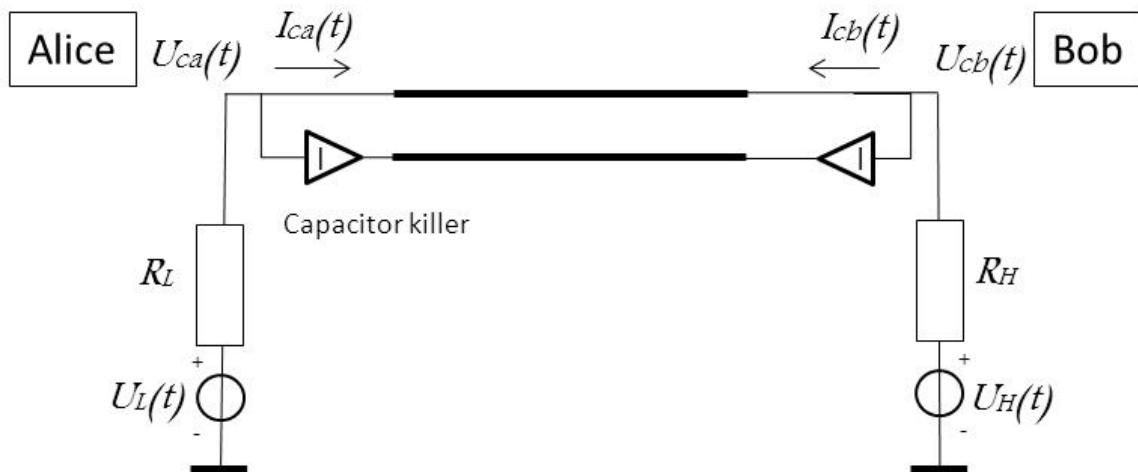


Figure 8. The KLJN system with the capacitor killer. A unity-gain voltage buffer is driving the outer shield, which is not grounded at this time.

4.2 Privacy Amplification

Another method to secure the key exchange and to reduce information leak is by utilizing privacy amplification. Privacy amplification is a software-based method to improve security on a secure key exchange [43]. The simplest and most secure concept is when Alice and Bob XOR the pairs of the key bits to have a half as long key with reduced ϵ . We simulated the effect of this technique by utilizing the most effective attack scenario. Simulation results showed that by XOR-ing once, Eve's success rate was reduced from 76.9% to 64.2%, and was further reduced to 54.4% by XOR-ing twice. Though privacy amplification is not as effective as a capacitor killer in reducing the information leak, it can be utilized to reduce any information leak due to the low bit error probability of the KLJN scheme [50-52].

Conclusions

By utilizing the LTSPICE simulator we have validated the cable capacitance attack. The results have shown that Eve success rate was 76.9% for a 1000 meters RG58 coaxial cable, with 100 samples per bit, at 10 bits/sec. To reduce the information leak due to the cable capacitance attack, capacitor killer and privacy amplification techniques have been utilized and simulated. The capacitor killer can fully eliminate the information leak by reducing Eve's success rate from 76.9% to 50.1%, while a privacy amplification technique can also arbitrarily reduce Eve's success rate but the price is slowing down. The simulations presented have demonstrated that the capacitor killer and privacy amplification are very successful against the cable capacitance attack, and that the unconditional security of a practical KLJN key exchange system is maintained against this attack, too.

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