

Quasi photons

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I. MODEL

$$\begin{aligned}
S_0 &= \int dx \left(-\frac{1}{4} F_{\mu\nu} f(-\square) F^{\mu\nu} - \frac{\xi}{2} \partial A f(-\square) \partial A \right) \\
&= \frac{1}{2} \int dx [-\partial_\mu A_\nu f(-\square) \partial^\mu A^\nu + \partial_\mu A_\nu f(-\square) \partial^\nu A^\mu + \xi A_\mu f(-\square) \partial^\mu \partial A] \\
&= \frac{1}{2} \int dx A_\mu [g^{\mu\nu} \square - (1 - \xi) \partial^\mu \partial^\nu] f(-\square) A_\nu \\
&= \frac{1}{2} A D_0^{-1} A
\end{aligned} \tag{1}$$

Feynman gauge: $\xi = 1$,

$$D_0^{-1}(p) = -g^{\mu\nu} p^2 \tag{2}$$

$$\Pi^{++\mu\nu}(q) = -g^{\mu\nu} \ell^2 (q^2)^2 \tag{3}$$

$$f(p^2) = 1 + \ell^2 (p^2)^2 \tag{4}$$

$$D(k) = \frac{1}{k^2 + i\epsilon} - \frac{1}{k^2 + \frac{1}{\ell^2} - i\epsilon} = \frac{1}{k^2 + i\epsilon} \frac{1}{1 + \ell^2(k^2 - i\epsilon)} = \frac{1}{k^2 + \ell^2(k^2)^2 + i\epsilon} \tag{5}$$

II. FIELD OPERATOR

A. Harmonic oscillator with negative norm

Linear space $H = H_+ \oplus H_-$, $\eta H_\sigma = \sigma H_\sigma$

Matrix elements: $\langle m|A|n\rangle = (m, \eta A n)$

Adjoint: $\langle m|\bar{A}|n\rangle = \langle n|A|m\rangle^*$, $\bar{A} = \sigma_A A$, $\sigma = \pm 1$

$A|\lambda\rangle = \lambda|\lambda\rangle$, $A|\rho\rangle = \rho|\rho\rangle$, $(\lambda - \sigma_A \rho^*) \langle \rho|\lambda\rangle = 0$:

real spectrum for skew-adjoint operators with non-orthogonal, degenerate eigenvectors

$\bar{q}_\sigma = \sigma q_\sigma$, $\bar{p}_\sigma = \sigma p_\sigma$, $[q_\sigma, p_\sigma] = i$

$a_\sigma = (m\omega q_\sigma + ip_\sigma)/\sqrt{2m\omega}$, $q_\sigma = (a_\sigma + \sigma \bar{a}_\sigma)/\sqrt{2m\omega}$, $p_\sigma = (a_\sigma - \sigma \bar{a}_\sigma)/\sqrt{2}i$, $[a_\sigma, \bar{a}_\sigma] = \sigma$

$H_\sigma = \sigma(\frac{1}{2m} p_\sigma^2 + \frac{m\omega^2}{2} q_\sigma^2) = \sigma\omega(\bar{a}_\sigma a_\sigma + \frac{1}{2})$

$b = a_+$ or \bar{a}_- , $\bar{b} = \bar{a}_+$ or a_- , $[b, \bar{b}] = 1$

basis: $\bar{b}b|\lambda\rangle = \lambda|\lambda\rangle, \dots, b^n|\lambda\rangle, \dots, b|\lambda\rangle, |\lambda\rangle, \bar{b}|\lambda\rangle, \dots, \bar{b}^n|\lambda\rangle, \dots$

$\bar{b}b$ eigenvalues: $\dots, \lambda - n, \dots, \lambda - 1, \lambda, \lambda + 1, \dots, \lambda + n, \dots$

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$\sigma = +1$: stops at left, $\lambda \geq 0$, $\text{sign}\langle \lambda + 1 | \lambda + 1 \rangle = \text{sign}\langle \lambda | \lambda \rangle$
 $\sigma = -1$: stops at right, $\lambda \leq -1$, $\text{sign}\langle \lambda - 1 | \lambda - 1 \rangle = -\text{sign}\langle \lambda | \lambda \rangle$
 $H_\sigma | \lambda + \sigma n \rangle = E_\sigma(n) | \lambda + \sigma n \rangle$, $E_\sigma(n) = n + \frac{1}{2} + \sigma \lambda$
 Coordinate eigenvalue: $\bar{q}_\sigma = \sigma q_\sigma$

$$\mathbb{1} = \int dq |\sigma q\rangle \langle q| = \int dp |\sigma p\rangle \langle p| \quad (6)$$

$$H = H_+ + H_-,$$

$$H_\sigma = \sigma \left(\frac{p_\sigma^2}{2m_\sigma} + \frac{m_\sigma \omega_\sigma^2}{2} q_\sigma^2 \right) = \sigma \omega_\sigma \left(\bar{a}_\sigma a_\sigma + \frac{1}{2} \right) \quad (7)$$

$$\begin{aligned} \langle q_f, -q'_f | e^{-itH} | q_i, q'_i \rangle &= \int D[p] D[p'] D[q] D[q'] e^{i \int dt [p\dot{q} + p'\dot{q}' - H(q, q', p, p')]} \\ &= \int D[q] D[q'] e^{i \int dt L(q, q', p, p')} \end{aligned} \quad (8)$$

$$H(q, q', p, p') = \frac{\langle q, -q' | H | p, p' \rangle}{\langle q, -q' | p, p' \rangle} = \frac{\langle q | H_+ | p \rangle}{\langle q | p \rangle} + \frac{\langle -q' | H_- | p' \rangle}{\langle -q' | p' \rangle} = \frac{p_+^2}{2m_+} + \frac{m_+ \omega_+^2}{2} q_+^2 - \frac{p_-^2}{2m_-} - \frac{m_- \omega_-^2}{2} q_-^2 \quad (9)$$

$$L(q, q', p, p') = \frac{\dot{q}_+^2}{2m_+} - \frac{m_+ \omega_+^2}{2} q_+^2 - \frac{\dot{q}_-^2}{2m_-} + \frac{m_- \omega_-^2}{2} q_-^2 \quad (10)$$

convergence is achieved by $\omega_\sigma \rightarrow \omega_\sigma - \sigma i \epsilon$

B. Field operator

$$M^2 = \ell^{-2}$$

$$L = \frac{1}{2} \partial_\mu \phi_+ \partial^\mu \phi_+ - \frac{m^2}{2} \phi_+^2 - \frac{1}{2} \partial_\mu \phi_- \partial^\mu \phi_- - \frac{M^2}{2} \phi_-^2 \quad (11)$$

negative mas square in the negative norm subspace.

$$\begin{aligned} [a(\mathbf{p}), a^\dagger(\mathbf{p}')] &= (2\pi)^3 2\omega_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{p}') \\ [b(\mathbf{p}), b^\dagger(\mathbf{p}')] &= -(2\pi)^3 2\Omega_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{p}'), \\ [a(\mathbf{p}), b^\dagger(\mathbf{p}')] &= [a(\mathbf{p}), b(\mathbf{p}')] = 0 \end{aligned} \quad (12)$$

$$\omega_{\mathbf{p}} = \sqrt{m^2 + \mathbf{p}^2}, \quad \Omega_{\mathbf{p}} = \sqrt{\mathbf{p}^2 - M^2}, \quad a(\mathbf{p})|0\rangle = b^\dagger(\mathbf{p})|0\rangle = 0, \quad \phi(x) = \phi_+(x) + \phi_-(x)$$

$$\begin{aligned} \phi_+(x) &= \int \frac{d^4 k}{(2\pi)^4} 2\pi \delta(k^2 - m^2) a(k) e^{-ikx} = \int_{m, \mathbf{k}} [a(\mathbf{k}) e^{-ikx} + a^\dagger(\mathbf{k}) e^{ikx}] \\ \phi_-(x) &= \int \frac{d^4 k}{(2\pi)^4} 2\pi \delta(k^2 + M^2) b^\dagger(k) e^{-ikx} = \int_{M, \mathbf{k}} [b(\mathbf{k}) e^{-ikx} + b^\dagger(\mathbf{k}) e^{ikx}] \end{aligned} \quad (13)$$

$$\int_{m, \mathbf{k}} f_{\mathbf{k}} = \int \frac{d^3 k}{(2\pi)^3 2\omega_{\mathbf{k}}} f_{\omega_{\mathbf{k}}, \mathbf{k}}, \quad \int_{M, \mathbf{k}} k_{\mathbf{k}} = \int_{\mathbf{k}^2 > |M^2|} \frac{d^3 k}{(2\pi)^3 2\Omega_{\mathbf{k}}} f_{\Omega_{\mathbf{k}}, \mathbf{k}} \quad (14)$$

C. Path integral

$$\begin{aligned}
e^{iW[\hat{j}]} &= \text{Tr} T[e^{-i \int_{t_i}^{t_f} dt' [H(t') - j_+(t') \phi_+(t')]} |0\rangle \langle 0| \bar{T}[e^{i \int_{t_i}^{t_f} dt' [H(t') + j_-(t') \phi_-(t')]}]] \\
&= \int D[\hat{\phi}] e^{\frac{i}{2} \hat{\phi} \cdot \hat{D}^{-1} \cdot \hat{\phi} + i \hat{j} \cdot \hat{\phi}} = e^{-\frac{i}{2} \hat{j} \cdot \hat{D} \cdot \hat{j}}
\end{aligned} \tag{15}$$

$$\begin{aligned}
i \frac{\delta^2 W[\hat{j}]}{\delta i j_a^+ \delta i j_b^+} &= \sum_n \langle 0 | \bar{T}[e^{i \int_{t_i}^{t_f} dt' H_i(t')}] |n\rangle \langle n | T[\phi_a \phi_b e^{-i \int_{t_i}^{t_f} dt' H_i(t')}] |0\rangle = \langle 0 | T[\phi_a \phi_b] |0\rangle = i D_{ab}^{++} \\
i \frac{\delta^2 W[\hat{j}]}{\delta i j_a^- \delta i j_b^-} &= \sum_n \langle 0 | \bar{T}[\phi_a \phi_b] |n\rangle \langle n | 0\rangle = \langle 0 | T[\phi_b \phi_a] |0\rangle^* = i D_{ab}^{--} = -i D_{ba}^{++*} \\
i \frac{\delta^2 W[\hat{j}]}{\delta i j_a^- \delta i j_b^+} &= \sum_n \langle 0 | \phi_a |n\rangle \langle n | \phi_b |0\rangle = \langle 0 | \phi_a \phi_b |0\rangle = i D_{ab}^{-+} \\
i \frac{\delta^2 W[\hat{j}]}{\delta i j_a^+ \delta i j_b^-} &= \sum_n \langle 0 | \phi_b |n\rangle \langle n | \phi_a |0\rangle = \langle 0 | \phi_b \phi_a |0\rangle = i D_{ab}^{+-} = -i D_{ab}^{-+*}
\end{aligned} \tag{16}$$

$$\begin{aligned}
T[\phi_a \phi_b] + \bar{T}[\phi_a \phi_b] &= \phi_a \phi_b + \phi_b \phi_a \\
D - D^\dagger &= D^{+-} - D^{-+*}
\end{aligned} \tag{17}$$

$$i \begin{pmatrix} D & D^{+-} \\ D^{-+} & D^{--} \end{pmatrix}_{x,y}^{j,k} = \begin{pmatrix} \langle T[\phi_x^j \phi_y^k] \rangle & \langle \phi_y^k \phi_x^j \rangle \\ \langle \phi_x^j \phi_y^k \rangle & \langle T[\phi_y^k \phi_x^j]^* \rangle \end{pmatrix} = i \begin{pmatrix} D & -D^{-+*} \\ D^{-+} & -D^\dagger \end{pmatrix}_{x,y}^{j,k} = i \begin{pmatrix} D^n + i \Im D & -D^f + i \Im D \\ D^f + i \Im D & -D^n + i \Im D \end{pmatrix} \tag{18}$$

$$\begin{aligned}
i D_{x,x'} &= \Theta(t-t') \langle 0 | \phi_x \phi_{x'} |0\rangle + \Theta(t'-t) \langle 0 | \phi_{x'} \phi_x |0\rangle \\
2\Re D_{x,x'} &= -\Theta(t-t') i \langle 0 | [\phi_x, \phi_{x'}] |0\rangle - \Theta(t'-t) i \langle 0 | [\phi_{x'}, \phi_x] |0\rangle = -\epsilon(t-t') i \langle 0 | [\phi_x, \phi_{x'}] |0\rangle = 2D_{x,x'}^n \\
2\Im D_{x,x'} &= -\Theta(t-t') \langle 0 | \{\phi_x, \phi_{x'}\} |0\rangle - \Theta(t'-t) \langle 0 | \{\phi_{x'}, \phi_x\} |0\rangle = -\langle 0 | \{\phi_x, \phi_{x'}\} |0\rangle \\
i D_{x,x'}^{-+} &= \langle 0 | \phi_x \phi_{x'} |0\rangle \\
2\Re D_{x,x'}^{-+} &= -i \langle 0 | \phi_x \phi_{x'} |0\rangle + i \langle 0 | \phi_{x'} \phi_x |0\rangle = -i \langle 0 | [\phi_x, \phi_{x'}] |0\rangle = 2D_{x,x'}^f \\
2\Im D_{x,x'}^{-+} &= -\langle 0 | \phi_x \phi_{x'} |0\rangle - \langle 0 | \phi_{x'} \phi_x |0\rangle = -\langle 0 | \{\phi_x, \phi_{x'}\} |0\rangle \\
i D_{x,x'}^{\bar{a}} &= i(D_{x,x'}^n \pm D_{x,x'}^f) = \frac{1}{2} \epsilon(t-t') \langle 0 | [\phi_x, \phi_{x'}] |0\rangle \pm \frac{1}{2} \langle 0 | \{\phi_x, \phi_{x'}\} |0\rangle = \pm \Theta(\pm(t-t')) \langle 0 | [\phi_x, \phi_{x'}] |0\rangle \\
D^{\bar{a}}(x) &= D(x) - D^{\pm\mp}(x) = \Theta(t) D^{-+}(x) + \Theta(-t) D^{-+}(-x) - D^{-+}(\mp x) \\
&= \Theta(t) D^{-+}(x) + \Theta(-t) D^{-+}(-x) - \Theta(t) D^{-+}(\mp x) - \Theta(-t) D^{-+}(\mp x) \\
&= \Theta(t) [D^{-+}(x) - D^{-+}(\mp x)] + \Theta(-t) [D^{-+}(-x) - D^{-+}(\mp x)] \\
&= \pm \Theta(\pm t) [D^{-+}(x) - D^{-+}(-x)]
\end{aligned} \tag{19}$$

D. Explicit expressions

$$\begin{aligned}
iD^{-+}(x, x') &= \frac{\langle 0|\phi(x)\phi(x')|0\rangle}{\langle 0|0\rangle} \\
&= \int_{m, \mathbf{k}, \mathbf{k}'} \frac{\langle 0|[a(\mathbf{k})e^{-ikx} + a^\dagger(\mathbf{k})e^{ikx}][a(\mathbf{k}')e^{-ik'x'} + a^\dagger(\mathbf{k}')e^{ik'x'}]|0\rangle}{\langle 0|0\rangle} \\
&\quad + \int_{M, \mathbf{k}, \mathbf{k}'} \frac{\langle 0|[b(\mathbf{k})e^{-ikx} + b^\dagger(\mathbf{k})e^{ikx}][b(\mathbf{k}')e^{-ik'x'} + b^\dagger(\mathbf{k}')e^{ik'x'}]|0\rangle}{\langle 0|0\rangle} \\
&= \int_{m, \mathbf{k}, \mathbf{k}'} \frac{\langle 0|a(\mathbf{k})a^\dagger(\mathbf{k}')e^{i(-kx+k'x')}|0\rangle}{\langle 0|0\rangle} + \int_{M, \mathbf{k}, \mathbf{k}'} \frac{\langle 0|b^\dagger(\mathbf{k})b(\mathbf{k}')e^{i(kx-k'x')}|0\rangle}{\langle 0|0\rangle} \\
&= \int_{m, \mathbf{k}} e^{-i\omega_{\mathbf{k}}(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')} + \int_{M, \mathbf{k}} e^{i\Omega_{\mathbf{k}}(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')} \\
iD^{+-}(x, x') &= \frac{\langle 0|\phi(x')\phi(x)|0\rangle}{\langle 0|0\rangle} = iD^{-+}(x', x) \tag{20}
\end{aligned}$$

$$\begin{aligned}
\Theta(t-t')D^{-+}(x, x') &= \int_{\omega} \frac{e^{-i\omega(t-t')}}{\omega + i\epsilon} \int_{m, \mathbf{k}} e^{-i\omega_{\mathbf{k}}(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')} + \int_{\omega} \frac{e^{i\omega(t-t')}}{-\omega + i\epsilon} \int_{M, \mathbf{k}} e^{i\Omega_{\mathbf{k}}(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')} \\
&= \int_{m, \mathbf{k}, \omega} \frac{e^{-i(\omega_{\mathbf{k}}+\omega)(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega + i\epsilon} - \int_{M, \mathbf{k}, \omega} \frac{e^{i(\omega+\Omega_{\mathbf{k}})(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - i\epsilon} \\
&= \int_{m, \mathbf{k}, \omega} \frac{e^{-i\omega(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - \omega_{\mathbf{k}} + i\epsilon} - \int_{M, \mathbf{k}, \omega} \frac{e^{i\omega(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - \Omega_{\mathbf{k}} - i\epsilon} \tag{21}
\end{aligned}$$

$$\begin{aligned}
\Theta(t'-t)D^{+-}(x, x') &= \int_{\omega} \frac{e^{i\omega(t-t')}}{\omega + i\epsilon} \int_{m, \mathbf{k}} e^{i\omega_{\mathbf{k}}(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')} + \int_{\omega} \frac{e^{-i\omega(t-t')}}{-\omega + i\epsilon} \int_{M, \mathbf{k}} e^{-i\Omega_{\mathbf{k}}(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')} \\
&= \int_{m, \mathbf{k}, \omega} \frac{e^{i(\omega_{\mathbf{k}}+\omega)(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega + i\epsilon} - \int_{M, \mathbf{k}, \omega} \frac{e^{-i(\omega+\Omega_{\mathbf{k}})(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - i\epsilon} \\
&= \int_{m, \mathbf{k}, \omega} \frac{e^{i\omega(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - \omega_{\mathbf{k}} + i\epsilon} - \int_{M, \mathbf{k}, \omega} \frac{e^{-i\omega(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega - \Omega_{\mathbf{k}} - i\epsilon} \\
&= - \int_{m, \mathbf{k}, \omega} \frac{e^{-i\omega(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega + \omega_{\mathbf{k}} - i\epsilon} + \int_{M, \mathbf{k}, \omega} \frac{e^{i\omega(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')}}{\omega + \Omega_{\mathbf{k}} + i\epsilon} \tag{22}
\end{aligned}$$

$$\begin{aligned}
D(x, x') &= \Theta(t-t')D^{-+}(x, x') + \Theta(t'-t)D^{+-}(x, x') \\
&= \int_{m, \mathbf{k}, \omega} e^{-i\omega(t-t') + i\mathbf{k}(\mathbf{x}-\mathbf{x}')} \left(\frac{1}{\omega - \omega_{\mathbf{k}} + i\epsilon} - \frac{1}{\omega + \omega_{\mathbf{k}} - i\epsilon} \right) \\
&\quad - \int_{M, \mathbf{k}, \omega} e^{i\omega(t-t') - i\mathbf{k}(\mathbf{x}-\mathbf{x}')} \left(\frac{1}{\omega - \Omega_{\mathbf{k}} - i\epsilon} - \frac{1}{\omega + \Omega_{\mathbf{k}} + i\epsilon} \right) \\
&= \int_{\mathbf{k}} \left[\frac{e^{-ikx}}{k^2 - m^2 + i\epsilon} - \frac{\Theta(\mathbf{k}^2 - M^2)e^{-ikx}}{k^2 + M^2 - i\epsilon} \right] \tag{23}
\end{aligned}$$

$$\begin{aligned}
D^{\bar{a}}(x) &= \pm\Theta(\pm t)[D^{-+}(x) - D^{-+}(-x)] \\
&= \mp\Theta(\pm t)i \left[\int_{m, \mathbf{k}} (e^{-i\omega_{\mathbf{k}}t + i\mathbf{k}\mathbf{x}} - e^{i\omega_{\mathbf{k}}t - i\mathbf{k}\mathbf{x}}) - \int_{M, \mathbf{k}} (e^{i\Omega_{\mathbf{k}}t - i\mathbf{k}\mathbf{x}} - e^{-i\Omega_{\mathbf{k}}t + i\mathbf{k}\mathbf{x}}) \right] \tag{24}
\end{aligned}$$

$$\begin{aligned}
D^{\bar{a}}(x) &= \pm\Theta(\pm t) \int_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} \left[\frac{1}{(k^0 \pm i\epsilon)^2 - \mathbf{k}^2 - m^2} - \frac{\Theta(\mathbf{k}^2 - M^2)}{(k^0 \pm i\epsilon)^2 - \mathbf{k}^2 + M^2} \right] \\
&= \pm\Theta(\pm t) \int_{\mathbf{k}} e^{-i\mathbf{k}\cdot\mathbf{x}} \left[\frac{1}{(k^0 \pm i\epsilon - \omega_{\mathbf{k}})(k^0 \pm i\epsilon + \omega_{\mathbf{k}})} - \frac{\Theta(\mathbf{k}^2 - M^2)}{(k^0 \pm i\epsilon - \Omega_{\mathbf{k}})(k^0 \pm i\epsilon + \Omega_{\mathbf{k}})} \right] \\
&= \pm\Theta(t) i \int_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x}} \left[\frac{e^{i\omega_{\mathbf{k}}t} - e^{-i\omega_{\mathbf{k}}t}}{2\omega_{\mathbf{k}}} - \Theta(\mathbf{k}^2 - M^2) \frac{e^{i\omega_{\mathbf{k}}t} - e^{-i\omega_{\mathbf{k}}t}}{2\Omega_{\mathbf{k}}} \right]
\end{aligned} \tag{25}$$

E. Real space expressions for $m = 0$

the contribution at large k are negligible in a cutoff theory:

$$\begin{aligned}
D(x) &= D_{x0} \rightarrow \int_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x} - \epsilon'k} \int_{\omega} \frac{e^{-i\omega t}}{\omega^2 - k^2 + i\epsilon} \\
&= \int_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{x} - \epsilon'k} \int_{\omega} \frac{e^{-i\omega t}}{(\omega - k + i\epsilon)(\omega + k - i\epsilon)} \\
&= -\frac{i}{4\pi^2} \int dk k^2 dc e^{ik(rc + i\epsilon')} \left(\frac{\Theta(t)e^{-i(k-i\epsilon)t}}{2k - i\epsilon} + \frac{\Theta(-t)e^{i(k-i\epsilon)t}}{2k - i\epsilon} \right) \\
&= -\frac{i}{4\pi^2} \int_0^\infty dk k^2 \frac{e^{ik(r+i\epsilon')} - e^{-ik(r-i\epsilon')}}{ikr} \frac{1}{2k} e^{-i(k-i\epsilon)|t|} \\
&= -\frac{1}{8\pi^2 r} \int_0^\infty dk (e^{ik(r-|t|+i\epsilon')} - e^{-ik(r+|t|-i\epsilon')}) \\
&= \frac{1}{8\pi^2 r} \left(\frac{1}{i(r-|t|+i\epsilon')} - \frac{1}{-i(r+|t|-i\epsilon')} \right) \\
&= -\frac{i}{8\pi^2 r} \left(\frac{1}{r-|t|+i\epsilon'} + \frac{1}{r+|t|-i\epsilon'} \right) \\
&= -\frac{i}{4\pi^2(r^2 - (|t| - i\epsilon')^2)} \\
&= \frac{i}{4\pi^2 x^2 - i\epsilon'} \\
&= P \frac{i}{4\pi^2 x^2} - \frac{1}{4\pi} \delta(x^2)
\end{aligned} \tag{26}$$

$$\begin{aligned}
D^{\bar{a}}(x) &= D_{x0}^{\bar{a}} = \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega \pm i\epsilon)^2 - \mathbf{k}^2} \\
&= \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega \pm i\epsilon - |\mathbf{k}|)(\omega \pm i\epsilon + |\mathbf{k}|)} \\
&= \underbrace{-\Theta(t)i \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \left(\frac{e^{-i(|\mathbf{k}|-i\epsilon)t}}{2|\mathbf{k}|} - \frac{e^{i(|\mathbf{k}+i\epsilon)t}}{2|\mathbf{k}|} \right)}_R + \underbrace{\Theta(-t)i \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \left(\frac{e^{-i(|\mathbf{k}+i\epsilon)t}}{2|\mathbf{k}|} - \frac{e^{i(|\mathbf{k}-i\epsilon)t}}{2|\mathbf{k}|} \right)}_A \\
&= -\left(\frac{\Theta(t)i}{(2\pi)^2} \int dk k^2 d\epsilon e^{ikrc} \frac{e^{-ikt} - e^{ikt}}{2k} - \frac{\Theta(-t)i}{(2\pi)^2} \int dk k^2 d\epsilon e^{ikrc} \frac{e^{-ikt} - e^{ikt}}{2k} \right) e^{-\epsilon|t|} \\
&= -\left(\frac{\Theta(t)i}{(2\pi)^2} \int dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-ikt} - e^{ikt}}{2k} - \frac{\Theta(-t)i}{(2\pi)^2} \int dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-ikt} - e^{ikt}}{2k} \right) e^{-\epsilon|t|} \\
&= -\left(\frac{\Theta(t)}{2(2\pi)^2 r} \int_0^\infty dk (e^{ikr} - e^{-ikr})(e^{-ikt} - e^{ikt}) - \frac{\Theta(-t)}{2(2\pi)^2 r} \int_0^\infty dk (e^{ikr} - e^{-ikr})(e^{-ikt} - e^{ikt}) \right) e^{-\epsilon|t|} \\
&= -\frac{1}{8\pi r} \int_{-\infty}^\infty \frac{dk}{2\pi} (e^{ik(r-t)} + e^{ik(-r+t)} - e^{ik(-r-t)} - e^{ik(r+t)})(\Theta(t) - \Theta(-t))e^{-\epsilon|t|} \\
&= -\frac{\Theta(t)\delta(r-t) + \Theta(-t)\delta(r+t)}{4\pi r} e^{-\epsilon|t|} \\
&= -\frac{\Theta(\pm t)\delta(t \mp r)}{4\pi r} e^{-\epsilon|t|} \\
&= -\Theta(\pm t) \frac{\delta(t^2 - r^2)}{2\pi} e^{-\epsilon|t|} \tag{27}
\end{aligned}$$

$$\begin{aligned}
D^f(x) &= \frac{1}{2}(D^r(x) - D^a(x)) = -\frac{1}{4\pi} \delta(x^2) \epsilon(x^0) \\
D^n(x) &= \frac{1}{2}(D^r(x) + D^a(x)) = -\frac{1}{4\pi} \delta(x^2) \tag{28}
\end{aligned}$$

F. Real space expressions for $m \neq 0$

$$D(k^2) = -i \text{sign}(\epsilon) \int_0^\infty ds e^{i \text{sign}(\epsilon) s(k^2 - m^2 + i\epsilon)}, \tag{29}$$

Analytic for $\Im e^{i\theta} > 0$:

$$\int dx e^{\frac{i}{2} a e^{i\theta} x^2} = \sqrt{\frac{2\pi}{a}} e^{i(\frac{\pi}{4} - \frac{\theta}{2})} \tag{30}$$

For real a :

$$\int \frac{d^4 k}{(2\pi)^4} e^{(ia-\epsilon)k^2 - ikx} = e^{-\frac{ix^2}{4a}} \int \frac{d^4 k}{(2\pi)^4} e^{(ia-\epsilon)(k - \frac{x}{2a})^2} = \frac{e^{i(\pi - \frac{3}{2}\pi)}}{32\pi^2 a^2} e^{-\frac{ix^2}{4a}} = -\frac{i \text{sign}(a)}{32\pi^2 a^2} e^{-\frac{ix^2}{4a}} \tag{31}$$

$$\begin{aligned}
D(x) &= -i\text{sign}(\epsilon) \int \frac{d^4k}{(2\pi)^4} \int_0^\infty ds e^{i\text{sign}(\epsilon')s(k^2 - m^2 + i\epsilon) - ikx} \\
&= -\frac{1}{32\pi^2} \int \frac{da}{a^2} e^{-i\text{sign}(\epsilon)a(m^2 - i\epsilon) - \frac{ix^2}{4a}} \quad \alpha = \frac{1}{4a} \\
&= \frac{1}{8\pi^2} \int_{-\infty}^\infty d\alpha e^{-\frac{im^2 - i\epsilon}{4\alpha} - ix^2\alpha} \\
&= \frac{1}{4\pi^2} \int_0^\infty d\alpha \cos\left(\frac{m^2}{4\alpha} + x^2\alpha\right) \\
&= \frac{1}{4\pi^2} \partial_{x^2} \int_0^\infty \frac{d\alpha}{\alpha} \sin\left(\frac{m^2}{4\alpha} + x^2\alpha\right) \quad \alpha = \frac{\sqrt{m^2 - i\epsilon}}{2\sqrt{|x^2|}} e^\theta \\
&= \frac{1}{4\pi^2} \partial_{x^2} \int_{-\infty}^\infty d\theta \sin\left[\frac{\sqrt{m^2 - i\epsilon}\sqrt{|x^2|}}{2}\right] (e^{-\theta} + \text{sign}(x^2)e^\theta) \\
&= \frac{\Theta(x^2)}{4\pi^2} \partial_{x^2} \int_{-\infty}^\infty d\theta \sin(\sqrt{(m^2 - i\epsilon)x^2} \cosh\theta) \\
&= \frac{\Theta(x^2)}{4\pi} \partial_{x^2} J_0(\sqrt{(m^2 - i\epsilon)x^2})
\end{aligned} \tag{32}$$

$$J'_0 = -J_1 \tag{33}$$

$$\begin{aligned}
D(x) &= \frac{1}{4\pi} \partial_{x^2} \Re H_0^{(1)}(m\sqrt{x^2}) \quad (\text{for } \sqrt{-\epsilon} = i\sqrt{\epsilon}) \\
&= -\frac{\Theta(x^2)}{8\pi} \sqrt{\frac{m^2 - i\epsilon}{x^2}} J_1(\sqrt{(m^2 - i\epsilon)x^2}) \\
&= \frac{1}{4\pi} \partial_{x^2} \Re H_0^{(1)}(m\sqrt{x^2}) \quad (\text{for } \sqrt{-\epsilon} = i\sqrt{\epsilon}) \\
&= \frac{1}{4\pi} \delta(x^2) - \frac{m^2}{8\pi} \Re \frac{H_1^{(1)}(m\sqrt{x^2})}{m\sqrt{x^2}}
\end{aligned} \tag{34}$$

Positive mass square:

$$\begin{aligned}
D^{\bar{a}}(x) &= D_{x0}^{\bar{a}} = \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega \pm i\epsilon)^2 - \mathbf{k}^2 - m^2} \\
&= \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega \pm i\epsilon - \omega_{\mathbf{k}})(\omega \pm i\epsilon + \omega_{\mathbf{k}})} \\
&= \underbrace{-\Theta(t)i \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \left(\frac{e^{-i(\omega_{\mathbf{k}} - i\epsilon)t}}{2\omega_{\mathbf{k}}} - \frac{e^{i(\omega_{\mathbf{k}} + i\epsilon)t}}{2\omega_{\mathbf{k}}} \right)}_R + \underbrace{\Theta(-t)i \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \left(\frac{e^{-i(\omega_{\mathbf{k}} + i\epsilon)t}}{2\omega_{\mathbf{k}}} - \frac{e^{i(\omega_{\mathbf{k}} - i\epsilon)t}}{2\omega_{\mathbf{k}}} \right)}_A \\
&= -\left(\frac{\Theta(t)i}{(2\pi)^2} \int dk k^2 d\epsilon e^{ikrc} \frac{e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}}{2\omega_{\mathbf{k}}} - \frac{\Theta(-t)i}{(2\pi)^2} \int dk k^2 d\epsilon e^{ikrc} \frac{e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}}{2\omega_{\mathbf{k}}} \right) e^{-\epsilon|t|} \\
&= -\left(\frac{\Theta(t)i}{(2\pi)^2} \int dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}}{2\omega_{\mathbf{k}}} - \frac{\Theta(-t)i}{(2\pi)^2} \int dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}}{2\omega_{\mathbf{k}}} \right) e^{-\epsilon|t|} \\
&= -\left(\frac{\Theta(t)}{2(2\pi)^2 r} \int_0^\infty \frac{dkk}{\omega_{\mathbf{k}}} (e^{ikr} - e^{-ikr})(e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}) - \frac{\Theta(-t)}{2(2\pi)^2 r} \int_0^\infty \frac{dkk}{\omega_{\mathbf{k}}} (e^{ikr} - e^{-ikr})(e^{-i\omega_{\mathbf{k}}t} - e^{i\omega_{\mathbf{k}}t}) \right) e^{-\epsilon|t|} \\
&= \mp \frac{1}{8\pi r} \int_{-\infty}^\infty \frac{dkk}{2\pi\omega_{\mathbf{k}}} (e^{i(kr - \omega_{\mathbf{k}}t)} + e^{i(-kr + \omega_{\mathbf{k}}t)} - e^{i(-kr - \omega_{\mathbf{k}}t)} - e^{i(kr + \omega_{\mathbf{k}}t)}) \Theta(\pm t) e^{-\epsilon|t|}
\end{aligned} \tag{35}$$

Negative mass square: $\tilde{\omega}_{\mathbf{k}} = +\sqrt{m^2 - \mathbf{k}^2} > 0$

$$\begin{aligned}
D^r(x) &= \int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{\omega^2 - (\mathbf{k}^2 - m^2)} + \text{homogeneous solutions (vanishing denominator)} \\
&= \int_{\mathbf{k} > m} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega + i\epsilon - \tilde{\omega}_{\mathbf{k}})(\omega + i\epsilon + \tilde{\omega}_{\mathbf{k}})} + \int_{\mathbf{k} < m} e^{i\mathbf{k}\mathbf{x}} \int_{\omega} \frac{e^{-i\omega t}}{(\omega - i\tilde{\omega}_{\mathbf{k}})(\omega + i\tilde{\omega}_{\mathbf{k}})} + \dots \\
&= -\Theta(t)i \left[\int_{\mathbf{k}} e^{i\mathbf{k}\mathbf{x}} \left(\frac{e^{-i(\tilde{\omega}_{\mathbf{k}} - i\epsilon)t}}{2\tilde{\omega}_{\mathbf{k}}} - \frac{e^{i(\tilde{\omega}_{\mathbf{k}} + i\epsilon)t}}{2\tilde{\omega}_{\mathbf{k}}} \right) - i \int_{\mathbf{k} < m} e^{i\mathbf{k}\mathbf{x}} \frac{e^{-\tilde{\omega}_{\mathbf{k}}t}}{2\tilde{\omega}_{\mathbf{k}}} \right] \\
&= -\frac{\Theta(t)}{(2\pi)^2} \left[i \int_m^\infty dk k^2 dce^{ikrc} \frac{e^{-i\tilde{\omega}_k t} - e^{i\tilde{\omega}_k t}}{2\tilde{\omega}_k} + \int_0^m dk k^2 dce^{ikrc} \frac{e^{-\tilde{\omega}_k t}}{2\tilde{\omega}_k} \right] \\
&= -\frac{\Theta(t)}{(2\pi)^2} \left[i \int_m^\infty dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-i\tilde{\omega}_k t} - e^{i\tilde{\omega}_k t}}{2\tilde{\omega}_k} + \int_0^m dk k^2 \frac{e^{ikr} - e^{-ikr}}{ikr} \frac{e^{-\tilde{\omega}_k t}}{2\tilde{\omega}_k} \right] \\
&= -\frac{\Theta(t)}{2(2\pi)^2 r} \left[\int_m^\infty \frac{dkk}{\tilde{\omega}_k} (e^{ikr} - e^{-ikr})(e^{-i\tilde{\omega}_k t} - e^{i\tilde{\omega}_k t}) - i \int_0^m \frac{dkk}{\tilde{\omega}_k} (e^{ikr} - e^{-ikr})e^{-\tilde{\omega}_k t} \right] \tag{36}
\end{aligned}$$