

Unification of fundamental forces

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Abstract

In this paper, I found a mathematical method to unify fundamental forces. In this method, without using any quantum gravity theories such as string theory and loop quantum gravity I created a field equation for the unification between fundamental forces. Also, I found out the acceleration equation of our universe.

1 Introduction

In mathematical methods in physics textbooks you can see the equation below:

$$\frac{1}{\sqrt{g}}\partial_\alpha(\sqrt{g})U^{\alpha\beta} = U^{\alpha\beta}\Gamma^\gamma_{\gamma\alpha} \quad (1)$$

$U^{\alpha\beta}$ is a tensor field. For unifying the forces I assume that the tensor field has four elements:

$$U^{\alpha\beta} = \eta G^{\alpha\beta} + \lambda F^{\alpha\beta} + \mu G^{\alpha\beta} + \gamma Z^\alpha W_\beta \quad (2)$$

As you see, the fundamental forces are tensor rank 2. For unifying them we have to put a coefficient for each force to have the same degree of freedom. $Z^\alpha W_\beta$ is electroweak, $G^{\alpha\beta}$ is the gluon field, is the electromagnetism field, and is the gravitation field. This research explains the effects of fundamental forces on the curvature of our universe. Finally, I achieved an equation that describes all fundamental forces in one equation. Another point I would like to mention is that the acceleration of our universe affects the final equation which means fundamental forces have a great impact on the acceleration of the universe.

Final field equation:

$$(8\kappa^2(A_\alpha - \frac{q}{m}F_{\lambda\alpha}\frac{dx^\lambda}{d\tau}) + 2\frac{\partial\mathcal{R}}{\partial x^\alpha}g\mathcal{R} + \frac{\partial g}{\partial x^\alpha}\mathcal{R}^2)U^{\alpha\beta} = 0 \quad (3)$$

The equation above is the modern version of $F=ma$. In fact, my model is generalized for gravity and quantum and classical physics.

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2 Proving Field Equation

Let's get started to do the math:

$$\mathcal{L} = \frac{1}{2\kappa} \sqrt{-g} \mathcal{R} = \sqrt{g_{\alpha\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}} \quad (4)$$

$$\frac{1}{\sqrt{g}} \partial_\alpha (\sqrt{g}) U^{\alpha\beta} = U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \quad (5)$$

$$\begin{aligned} \frac{1}{\sqrt{g}} \partial_\alpha \left(\frac{2\kappa \mathcal{L}}{i\mathcal{R}} \right) U^{\alpha\beta} &= U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha}, \\ \frac{\partial \mathcal{L}}{\partial \mathcal{X}^\alpha} &= \frac{1}{2\mathcal{L}} \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}, \\ \frac{2\kappa}{\sqrt{g}} \left(\frac{\frac{1}{2\mathcal{L}} \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} i\mathcal{R} - i \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L} \right) U^{\alpha\beta} &= U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha}, \\ 2\kappa \left(\frac{\frac{1}{2\mathcal{L}} \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} i\mathcal{R} - i \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L} \right) U^{\alpha\beta} &= -\sqrt{-g} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha}, \\ 2\kappa \left(\frac{1}{2} \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \sqrt{-g} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{L} \mathcal{R}^2, \\ \mathcal{L} &= \frac{1}{2\kappa} \sqrt{-g} \mathcal{R} \rightarrow \sqrt{-g} = \frac{2\kappa \mathcal{L}}{\mathcal{R}}, \\ 2\kappa \left(\frac{1}{2} \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \frac{2\kappa \mathcal{L}^2}{\mathcal{R}} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{R}^2 \\ \nabla_\alpha g_{\alpha\beta} &= 0 \\ \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} - \Gamma_{\alpha\alpha}^d g_{d\beta} - \Gamma_{\alpha\beta}^d g_{\alpha d} &= 0 \\ \frac{\partial g_{\alpha\beta}}{\partial x^\alpha} &= \Gamma_{\alpha\alpha}^d g_{d\beta} + \Gamma_{\alpha\beta}^d g_{\alpha d} \end{aligned} \quad (6)$$

$$\begin{aligned} \left(\frac{1}{2} (\Gamma_{\alpha\alpha}^d g_{d\beta} + \Gamma_{\alpha\beta}^d g_{\alpha d}) \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \frac{\mathcal{L}^2}{\mathcal{R}} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{R}^2 \\ \left(\frac{1}{2} \Gamma_{\alpha\alpha}^d g_{d\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} + \frac{1}{2} \Gamma_{\alpha\beta}^d g_{\alpha d} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \frac{\mathcal{L}^2}{\mathcal{R}} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{R}^2 \end{aligned} \quad (7)$$

$$\begin{aligned} \Gamma_{\alpha\beta}^d \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} &= \frac{d^2 x^d}{d\tau^2} + X^d \\ \left(\frac{1}{2} \Gamma_{\alpha\alpha}^d g_{d\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} \mathcal{R} + \frac{1}{2} \frac{d^2 x^d}{d\tau^2} g_{\alpha d} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \frac{\mathcal{L}^2}{\mathcal{R}} U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{R}^2 \\ \left(\left(\frac{d^2 x^d}{d\tau^2} + X^d \right) g_{\alpha d} \mathcal{R} - \frac{\partial \mathcal{R}}{\partial x^\alpha} \mathcal{L}^2 \right) U^{\alpha\beta} &= \mathcal{L}^2 U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \mathcal{R} \\ \mathcal{L}^2 &= -\frac{1}{4\kappa^2} g \mathcal{R}^2 \end{aligned} \quad (8)$$

$$\begin{aligned} (4\kappa^2 \left(\frac{d^2 x^d}{d\tau^2} + X^d \right) g_{\alpha d} \mathcal{R} + \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R}^2) U^{\alpha\beta} &= -g \mathcal{R}^3 U^{\alpha\beta} \Gamma^\gamma_{\gamma\alpha} \\ \Gamma^\gamma_{\gamma\alpha} &= \frac{1}{2g} \frac{\partial g}{\partial x^\alpha} \\ (4\kappa^2 \left(\frac{d^2 x^d}{d\tau^2} + X^d \right) g_{\alpha d} \mathcal{R} + \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R}^2) U^{\alpha\beta} &= -g \mathcal{R}^3 U^{\alpha\beta} \frac{1}{2g} \frac{\partial g}{\partial x^\alpha} \\ (8\kappa^2 \left(\frac{d^2 x^d}{d\tau^2} + X^d \right) g_{\alpha d} + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R}) U^{\alpha\beta} + \frac{\partial g}{\partial x^\alpha} U^{\alpha\beta} \mathcal{R}^2 &= 0 \\ U^{\alpha\beta} &= \eta G^{\alpha\beta} + \lambda F^{\alpha\beta} + \mu G^{\alpha\beta} + \gamma Z^\alpha W_\beta \\ A^d &= \frac{d^2 x^d}{d\tau^2} \\ (8\kappa^2 (A_\alpha + X^d g_{\alpha d}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R}) U^{\alpha\beta} + \frac{\partial g}{\partial x^\alpha} U^{\alpha\beta} \mathcal{R}^2 &= 0 \end{aligned} \quad (9)$$

$$\begin{aligned} X^d &= -\frac{g}{m} F^{d\gamma} \frac{dx^\lambda}{d\tau} g_{\lambda\gamma} \\ (8\kappa^2 (A_\alpha - \frac{g}{m} F^{d\gamma} \frac{dx^\lambda}{d\tau} g_{\lambda\gamma} g_{\alpha d}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R}) U^{\alpha\beta} + \frac{\partial g}{\partial x^\alpha} U^{\alpha\beta} \mathcal{R}^2 &= 0 \end{aligned} \quad (10)$$

$$(8\kappa^2 (A_\alpha - \frac{g}{m} F_{\alpha\lambda} \frac{dx^\lambda}{d\tau}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2) U^{\alpha\beta} = 0$$

$$(8\kappa^2 (A_\alpha - \frac{g}{m} F_{\alpha\lambda} \frac{dx^\lambda}{d\tau}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2) (\eta G^{\alpha\beta} + \lambda F^{\alpha\beta} + \mu G^{\alpha\beta} + \gamma Z^\alpha W_\beta) = 0 \quad (11)$$

For finding the coefficients we have to use numerical methods or experiments.

3 One Example of The Equation Field

For example, I want to consider only the electromagnetic field in the equation above:

$$U^{\alpha\beta} = F^{\alpha\beta} \quad (12)$$

We will see a familiar equation. I have to mention that in this example I consider the flat space-time; however, we can calculate for the curved space-time, but for simplicity, I consider the flat space-time. We chose Beta as zero for finding an equation.

$$\begin{aligned}
U^{\alpha\beta} &= F^{\alpha\beta} \\
A^d &= \frac{d^2 x^d}{d\tau^2} \\
\beta &= 0 \\
8\kappa^2(A_\alpha F^{\alpha\beta} - \frac{q}{m} F_{\alpha\lambda} \frac{dx^\lambda}{d\tau} F^{\alpha\beta}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} F^{\alpha\beta} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2 F^{\alpha\beta} &= 0 \\
8\kappa^2(A_1 F^{10} + A_2 F^{20} + A_3 F^{30} - \frac{q}{m} (F_{\alpha 0} \frac{dx^0}{d\tau} + F_{\alpha 1} \frac{dx^1}{d\tau} + F_{\alpha 2} \frac{dx^2}{d\tau} + F_{\alpha 3} \frac{dx^3}{d\tau}) F^{\alpha 0}) \\
+ 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} F^{\alpha 0} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2 F^{\alpha 0} &= 0 \\
\text{For flatspacetime} \\
2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} F^{\alpha 0} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2 F^{\alpha 0} &= 0 \\
& \quad (13) \\
8\kappa^2(A_1 F^{10} + A_2 F^{20} + A_3 F^{30} - \frac{q}{m} (F_{\alpha 0} F^{\alpha 0} \frac{dx^0}{d\tau} + F_{\alpha 1} F^{\alpha 0} \frac{dx^1}{d\tau} + F_{\alpha 2} F^{\alpha 0} \frac{dx^2}{d\tau} + F_{\alpha 3} F^{\alpha 0} \frac{dx^3}{d\tau})) &= 0 \\
F_{\alpha 0} F^{\alpha 0} = F_{10} F^{10} + F_{20} F^{20} + F_{30} F^{30} = -(E_x^2 + E_y^2 + E_z^2) = -E^2 \\
F_{\alpha 1} F^{\alpha 0} = F_{21} F^{20} + F_{31} F^{30} = B_z E_y - B_y E_z \\
F_{\alpha 2} F^{\alpha 0} = F_{12} F^{10} + F_{32} F^{30} = -B_z E_x + B_x E_z \\
F_{\alpha 3} F^{\alpha 0} = F_{13} F^{10} + F_{23} F^{20} = B_y E_x - B_x E_y \\
A_1 F^{10} + A_2 F^{20} + A_3 F^{30} - \frac{q}{m} (F_{\alpha 0} F^{\alpha 0} \frac{dx^0}{d\tau} + F_{\alpha 1} F^{\alpha 0} \frac{dx^1}{d\tau} + F_{\alpha 2} F^{\alpha 0} \frac{dx^2}{d\tau} + F_{\alpha 3} F^{\alpha 0} \frac{dx^3}{d\tau}) &= 0 \\
\vec{A} \cdot \vec{E} - \frac{q}{m} (\gamma E^2 + (B_z E_y - B_y E_z) \frac{dx^1}{d\tau} + (-B_z E_x + B_x E_z) \frac{dx^2}{d\tau} + (B_y E_x - B_x E_y) \frac{dx^3}{d\tau}) &= 0 \\
\vec{A} \cdot \vec{E} - \frac{q}{m} (\gamma E^2 + (\vec{B} \times \vec{E}) \cdot \vec{v}) &= 0 \\
& \quad (14)
\end{aligned}$$

This is familiar $\vec{A} \cdot \vec{E} - \frac{q}{m} (\gamma E^2 + (\vec{B} \times \vec{E}) \cdot \vec{v}) = 0$ for us (Lorentz force equation).

4 Acceleration of Our Universe

For every diagonal metric, such as the Robertson-walker metric :

$$\begin{aligned}
(8\kappa^2(A_\alpha - \frac{q}{m} F_{\alpha\lambda} \frac{dx^\lambda}{d\tau}) + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2) 8\pi G T_{\alpha\beta} &= 0 \\
\text{Assume :} \\
F_{\alpha\lambda} &= 0 \\
(8\kappa^2 A_\alpha + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2) T^{\alpha\beta} &= 0 \\
8\kappa^2 A_\alpha + 2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2 &= 0 \\
A_\alpha &= -\frac{1}{8\kappa^2} (2 \frac{\partial \mathcal{R}}{\partial x^\alpha} g \mathcal{R} + \frac{\partial g}{\partial x^\alpha} \mathcal{R}^2)
\end{aligned} \quad (15)$$

The acceleration is positive because R in the terms is squared, and g is negative; hence, The acceleration is positive. Therefore, the acceleration of the universe is positive. this result is compatible with other assumptions.

5 Conclusion

I do not feel like calculating field equations for all fundamental forces, and all kinds of metrics such as black holes and the universe, but you can calculate with this equation.

$$(8\kappa^2(A_\alpha - \frac{q}{m}F_{\lambda\alpha} \frac{dx^\lambda}{d\tau}) + 2\frac{\partial\mathcal{R}}{\partial x^\alpha}g\mathcal{R} + \frac{\partial g}{\partial x^\alpha}\mathcal{R}^2)U^{\alpha\beta} = 0 \quad (16)$$